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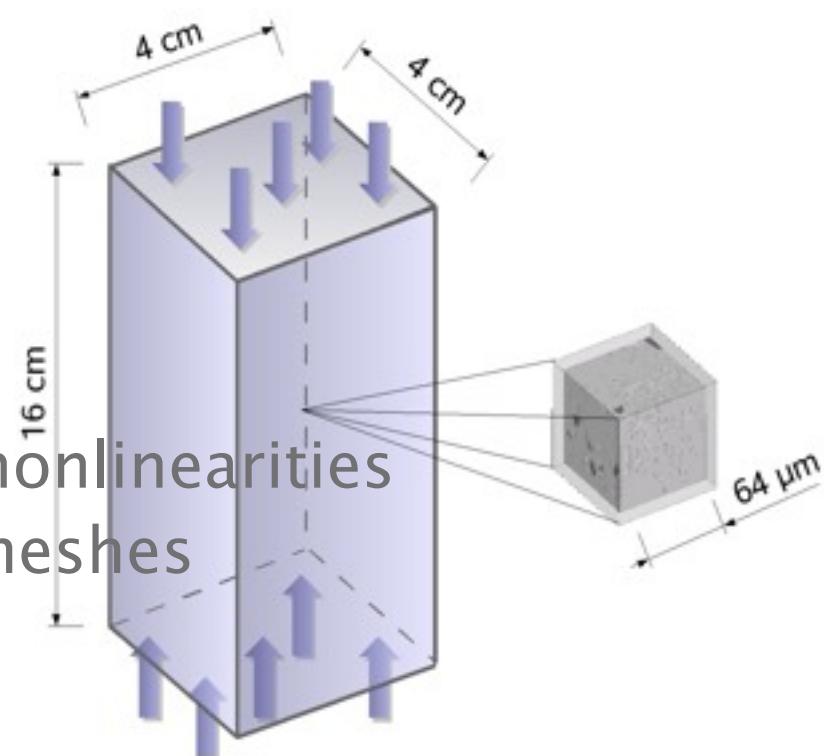
Finite Elements for Large Strains – A double mixed (M^2) Formulation

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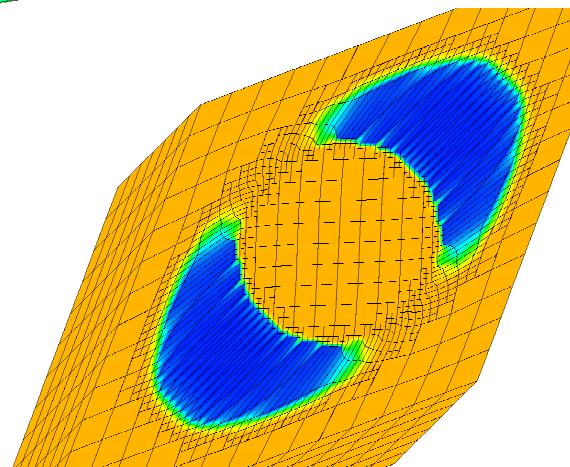
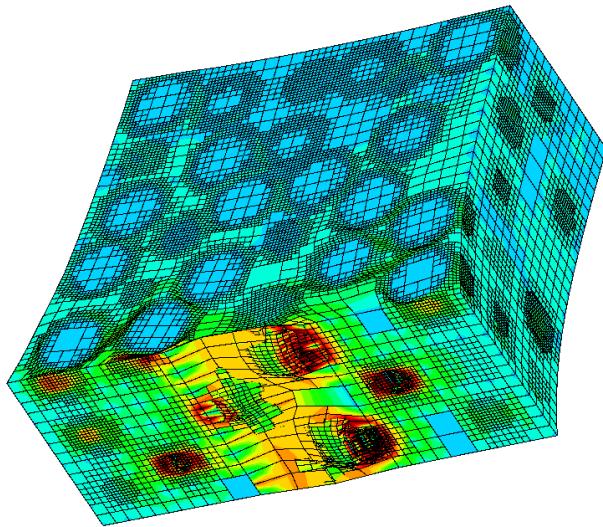
Motivation

Development of user friendly elements

- robustness
- simple treatment of
 - incompressible materials
 - complex geometries
 - geometrical and material nonlinearities
 - coarse and non-uniform meshes
- efficiency
- simple implementation

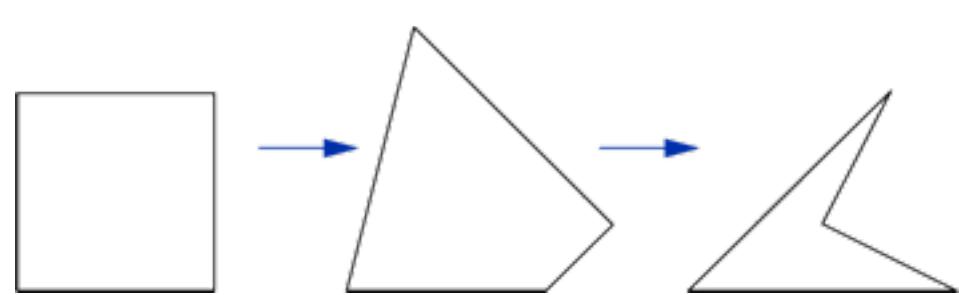


Numerical simulation of incompressible materials is still challenging within the context of large deformations



Challenges:

- Locking in incompressibility
- Large deformations
- Non-convex element shapes
- Robustness





Deformation Measures

$$\boldsymbol{F}(\boldsymbol{X}) = \text{Grad } \varphi(\boldsymbol{X})$$

$$(\boldsymbol{F}, \text{cof } \boldsymbol{F}, \det \boldsymbol{F})$$

$$d\boldsymbol{x} = \boldsymbol{F} d\boldsymbol{X}, \quad \boldsymbol{n} da = \text{cof}[\boldsymbol{F}] \boldsymbol{N} dA \quad \text{and} \quad dv = \det[\boldsymbol{F}] dV$$

$$\text{cof} [\boldsymbol{F}] = \begin{bmatrix} F_{22}F_{33} - F_{32}F_{23} & F_{31}F_{23} - F_{21}F_{33} & F_{21}F_{32} - F_{31}F_{22} \\ F_{32}F_{13} - F_{12}F_{33} & F_{11}F_{33} - F_{31}F_{13} & F_{31}F_{12} - F_{11}F_{32} \\ F_{12}F_{23} - F_{22}F_{13} & F_{21}F_{13} - F_{11}F_{23} & F_{11}F_{22} - F_{21}F_{12} \end{bmatrix}$$

$$\boldsymbol{C} = \boldsymbol{F}^T \boldsymbol{F}$$



Polyconvex Material

[**Definition of Polyconvexity**] $\mathbf{F} \mapsto W(\mathbf{F})$ is polyconvex if and only if there exists a function $P : \mathbb{R}^{3 \times 3} \times \mathbb{R}^{3 \times 3} \times \mathbb{R} \mapsto \mathbb{R}$ (in general non-unique) such that

$$W(\mathbf{F}) = P(\mathbf{F}, \text{cof}\mathbf{F}, |\det\mathbf{F}|) \quad (10)$$

and $(\mathbf{F}, \text{cof}\mathbf{F}, \det\mathbf{F}) \in \mathbb{R}^{19} \mapsto P(\mathbf{F}, \text{cof}\mathbf{F}, \det\mathbf{F}) \in \mathbb{R}$ is convex for all points $\mathbf{X} \in \mathbb{R}^3$. \square

Potential

$$\Pi(\mathbf{x}) = \int_{\mathcal{B}} W(\mathbf{F}(\mathbf{x})) dV + \Pi^{ext}(\mathbf{x})$$

Hu–Washizu Functional

$$\Pi(\mathbf{x}, \mathbf{H}^c, \theta, p, \mathbf{B}) = \int_{\mathcal{B}} \mathcal{P}(\mathbf{F}, \mathbf{H}^c, \theta) dV + \int_{\mathcal{B}} p(J - \theta) dV + \int_{\mathcal{B}} \mathbf{B}(\text{cof}\mathbf{F} - \mathbf{H}^c) dV$$



Hu-Washizu Functional

$$\Pi(\boldsymbol{x}, \boldsymbol{H}^e, \theta, p, \boldsymbol{B}) = \int_{\mathcal{B}} \mathcal{P}(\boldsymbol{C}, \boldsymbol{H}^e, \theta) dV + \int_{\mathcal{B}} p(J - \theta) dV + \int_{\mathcal{B}} \boldsymbol{B} \cdot (\text{cof} \boldsymbol{C} - \boldsymbol{H}^e) dV$$



Hu-Washizu Functional

$$\Pi(\boldsymbol{x}, \boldsymbol{H}^c, \theta, p, \boldsymbol{B}) = \int_{\mathcal{B}} \mathcal{P}(\boldsymbol{C}, \boldsymbol{H}^c, \theta) dV + \int_{\mathcal{B}} p(J - \theta) dV + \int_{\mathcal{B}} \boldsymbol{B} \cdot (\text{cof}\boldsymbol{C} - \boldsymbol{H}^c) dV$$

Variation

$$\delta_C \Pi = \int_{\mathcal{B}} \frac{1}{2} \delta \boldsymbol{C} \cdot \underbrace{2(\partial_C \mathcal{P} + \frac{p}{2J} \text{cof}\boldsymbol{C} + \boldsymbol{B} : \partial_C \text{cof}\boldsymbol{C})}_{S} dV,$$

$$\delta_\theta \Pi = \int_{\mathcal{B}} \delta \theta (\partial_\theta \mathcal{P} - p) dV \quad \rightarrow \quad p = \partial_\theta \mathcal{P},$$

$$\delta_p \Pi = \int_{\mathcal{B}} \delta p (J - \theta) dV \quad \rightarrow \quad J = \theta,$$

$$\delta_H \Pi = \int_{\mathcal{B}} \delta \boldsymbol{H}^c \cdot (\partial_H \mathcal{P} - \boldsymbol{B}) dV \quad \rightarrow \quad \boldsymbol{B} = \partial_H \mathcal{P},$$

$$\delta_B \Pi = \int_{\mathcal{B}} \delta \boldsymbol{B} \cdot (\text{cof}\boldsymbol{C} - \boldsymbol{H}^c) dV \quad \rightarrow \quad \boldsymbol{H}^c = \text{cof}\boldsymbol{C}.$$



Hu-Washizu Functional

Choice of Strain Energy Function

$$\mathcal{P}(\mathbf{C}, \mathbf{H}^c, \theta) = \mathcal{P}_1(\mathbf{C}) + \mathcal{P}_2(\mathbf{H}^c) + \mathcal{P}_3(\theta)$$



Hu-Washizu Functional

Choice of Strain Energy Function

$$\mathcal{P}(\mathbf{C}, \mathbf{H}^c, \theta) = \mathcal{P}_1(\mathbf{C}) + \mathcal{P}_2(\mathbf{H}^c) + \mathcal{P}_3(\theta)$$

Linearization, Definitions

$$\mathbb{C}_1 = \partial_{CC}^2 \mathcal{P}, \quad \mathbb{C}_2 = \partial_{H^c H^c}^2 \mathcal{P}, \quad \mathbb{C}_3 = \partial_{\theta\theta}^2 \mathcal{P}$$

$$\mathbb{P}^c := \partial_C \text{cof} \mathbf{C} = (\text{cof} \mathbf{C} \otimes \text{cof} \mathbf{C} - \text{cof} \mathbf{C} \boxtimes \text{cof} \mathbf{C}) / \det \mathbf{C}, \quad \Xi = \partial_{CC}^2 \text{cof} \mathbf{C}$$

$$(\mathbf{A} \boxtimes \mathbf{B}) : \mathbf{a} \otimes \mathbf{b} = \mathbf{A} \mathbf{a} \otimes \mathbf{B} \mathbf{b}$$

$$\overline{\mathbb{C}} = \frac{p}{2J} \mathbb{P}^c - \frac{p}{4J^3} \text{cof} \mathbf{C} \otimes \text{cof} \mathbf{C} + \mathbf{B} : \Xi$$



Linearization of Hu-Washizu Functional

$$\begin{aligned}\Delta\delta_C\Pi &= \int_{\mathcal{B}} \frac{1}{2} \Delta\delta\mathbf{C} \cdot \mathbf{S} dV + \int_{\mathcal{B}} \frac{1}{2} \delta\mathbf{C} \cdot [4(\mathbb{C}_1 + \bar{\mathbb{C}}) : \frac{1}{2} \Delta\mathbf{C}] dV \\ &\quad + \int_{\mathcal{B}} \frac{1}{2} \delta\mathbf{C} \cdot (\frac{1}{J} \text{cof}\mathbf{C} \Delta p) dV + \int_{\mathcal{B}} \frac{1}{2} \delta\mathbf{C} \cdot 2(\mathbb{P}^c : \Delta\mathbf{B}) dV\end{aligned}$$

$$\Delta\delta_\theta\Pi = \int_{\mathcal{B}} \delta\theta \mathbb{C}_3 \Delta\theta dV - \int_{\mathcal{B}} \delta\theta \Delta p dV,$$

$$\Delta\delta_p\Pi = \int_{\mathcal{B}} \delta p \frac{1}{J} \text{cof}\mathbf{C} \cdot \frac{1}{2} \Delta\mathbf{C} dV - \int_{\mathcal{B}} \delta p \Delta\theta dV,$$

$$\Delta\delta_H\Pi = \int_{\mathcal{B}} \delta\mathbf{H}^c \cdot (\mathbb{C}_2 : \Delta\mathbf{H}^c) dV - \int_{\mathcal{B}} \delta\mathbf{H}^c \cdot \Delta\mathbf{B} dV.$$

$$\Delta\delta_B\Pi = \int_{\mathcal{B}} \delta\mathbf{B} \cdot (2\mathbb{P}^c : \frac{1}{2} \Delta\mathbf{C}) dV - \int_{\mathcal{B}} \delta\mathbf{B} \cdot \Delta\mathbf{H}^c dV$$



Finite Element Approximation

- a quadratic interpolation for x , $\boldsymbol{u} = \mathbf{N} \boldsymbol{d}_u$
- a linear interpolation for \mathbf{H}^c and \mathbf{B} $\mathbf{H}^c = \mathbf{N}_c \boldsymbol{d}_c, \quad \mathbf{B} = \mathbf{N}_b \boldsymbol{d}_b$
- a constant interpolation for θ and p . $\theta = \mathbf{N}_\theta \boldsymbol{d}_\theta \quad p = \mathbf{N}_p \boldsymbol{d}_p$

$$\begin{aligned}\Delta \delta_C \Pi &= \delta \boldsymbol{d}_u^T \left[\mathbf{k}^{geo} + \int_{\mathcal{B}} \mathbf{B}^T 4 (\mathbf{C}_1 + \overline{\mathbf{C}}) \mathbf{B} dV \right] \Delta \boldsymbol{d}_u \\ &+ \delta \boldsymbol{d}_u^T \left\{ \int_{\mathcal{B}} \mathbf{B}^T \frac{1}{J} \text{cof} \mathbf{C} \mathbf{N}_p dV \Delta \boldsymbol{d}_p + \int_{\mathcal{B}} \mathbf{B}^T 2 \mathbb{P}^c \mathbf{N}_b dV \Delta \boldsymbol{d}_b \right\} \\ &= \delta \boldsymbol{d}_u^T \left\{ \mathbf{k}_{uu} \Delta \boldsymbol{d}_u + \mathbf{k}_{up} \Delta \boldsymbol{d}_p + \mathbf{k}_{ub} \Delta \boldsymbol{d}_b \right\},\end{aligned}$$
$$\frac{1}{2} \delta \mathbf{C} = \mathbf{B} \delta \boldsymbol{d}_u \quad \text{and} \quad \frac{1}{2} \Delta \mathbf{C} = \mathbf{B} \Delta \boldsymbol{d}_u$$



Finite Element Approximation

$$\Delta\delta_\theta\Pi = \delta\mathbf{d}_\theta^T \int_{\mathcal{B}} \mathbf{N}_\theta^T \mathbf{C}_3 \mathbf{N}_\theta dV \Delta\mathbf{d}_\theta - \delta\mathbf{d}_\theta^T \int_{\mathcal{B}} \mathbf{N}_\theta^T \mathbf{N}_p dV \Delta\mathbf{d}_p$$

$$= \delta\mathbf{d}_\theta^T \mathbf{k}_{\theta\theta} \Delta\mathbf{d}_\theta + \delta\mathbf{d}_\theta^T \mathbf{k}_{\theta p} \Delta\mathbf{d}_p,$$

$$\Delta\delta_p\Pi = \delta\mathbf{d}_p^T \int_{\mathcal{B}} \mathbf{N}_p^T \frac{1}{J} \text{cof} \mathbf{C} \mathbf{B} dV \Delta\mathbf{d}_u - \delta\mathbf{d}_p^T \int_{\mathcal{B}} \mathbf{N}_p^T \mathbf{N}_\theta dV \Delta\mathbf{d}_\theta$$

$$= \delta\mathbf{d}_p^T \mathbf{k}_{pu} \Delta\mathbf{d}_u + \delta\mathbf{d}_p^T \mathbf{k}_{p\theta} \Delta\mathbf{d}_\theta,$$

$$\Delta\delta_c\Pi = \delta\mathbf{d}_c^T \int_{\mathcal{B}} \mathbf{N}_c^T \mathbf{C}_2 \mathbf{N}_c dV \Delta\mathbf{d}_c - \delta\mathbf{d}_c^T \int_{\mathcal{B}} \mathbf{N}_c^T \mathbf{N}_b dV \Delta\mathbf{d}_b$$

$$= \delta\mathbf{d}_c^T \mathbf{k}_{cc} \Delta\mathbf{d}_c + \delta\mathbf{d}_c^T \mathbf{k}_{cb} \Delta\mathbf{d}_b,$$

$$\Delta\delta_b\Pi = \delta\mathbf{d}_b^T \int_{\mathcal{B}} \mathbf{N}_b^T 2 \mathbb{P}^c \mathbf{B} dV \Delta\mathbf{d}_u - \delta\mathbf{d}_b^T \int_{\mathcal{B}} \mathbf{N}_b^T \mathbf{N}_c dV \Delta\mathbf{d}_c$$

$$= \delta\mathbf{d}_b^T \mathbf{k}_{bu} \Delta\mathbf{d}_u + \delta\mathbf{d}_b^T \mathbf{k}_{bc} \Delta\mathbf{d}_c$$



Finite Element Approximation

$$\begin{bmatrix} \Delta\delta_C\Pi \\ \Delta\delta_\theta\Pi \\ \Delta\delta_p\Pi \\ \Delta\delta_H\Pi \\ \Delta\delta_B\Pi \end{bmatrix} = \begin{bmatrix} \delta\mathbf{d}_u \\ \delta\mathbf{d}_\theta \\ \delta\mathbf{d}_p \\ \delta\mathbf{d}_c \\ \delta\mathbf{d}_b \end{bmatrix}^T \begin{bmatrix} k_{uu} & 0 & k_{up} & 0 & k_{ub} \\ 0 & k_{\theta\theta} & k_{\theta p} & 0 & 0 \\ k_{pu} & k_{p\theta} & 0 & 0 & 0 \\ 0 & 0 & 0 & k_{cc} & k_{cb} \\ k_{bu} & 0 & 0 & k_{bc} & 0 \end{bmatrix} \begin{bmatrix} \Delta\mathbf{d}_u \\ \Delta\mathbf{d}_\theta \\ \Delta\mathbf{d}_p \\ \Delta\mathbf{d}_c \\ \Delta\mathbf{d}_b \end{bmatrix}$$



Newton Solution Procedure

$$k_{uu}\Delta d_u + k_{up}\Delta d_p + k_{ub}\Delta d_b = -r_u$$

$$k_{\theta\theta}\Delta d_\theta + k_{\theta p}\Delta d_p = -r_\theta$$

$$k_{pu}\Delta d_u + k_{p\theta}\Delta d_\theta = -r_p$$

$$k_{cc}\Delta d_c + k_{cb}\Delta d_b = -r_c$$

$$k_{bu}\Delta d_u + k_{bc}\Delta d_c = -r_b$$

$$\left[k_{uu} + k_{up}\tilde{k}_{pp}^{-1}k_{pu} + k_{ub}\tilde{k}_{cc}k_{bu} \right] \Delta d_u = - \left[r_u + k_{up}\tilde{k}_{pp}^{-1}\tilde{r}_p + k_{ub}k_{cb}^{-1}\tilde{r}_c \right]$$

$$\tilde{k}_{cc} = k_{cb}^{-1}k_{cc}k_{bc}^{-1} \quad \text{and} \quad \tilde{r}_c = -r_c + k_{cc}k_{bc}^{-1}r_b.$$

Mixed finite element formulation

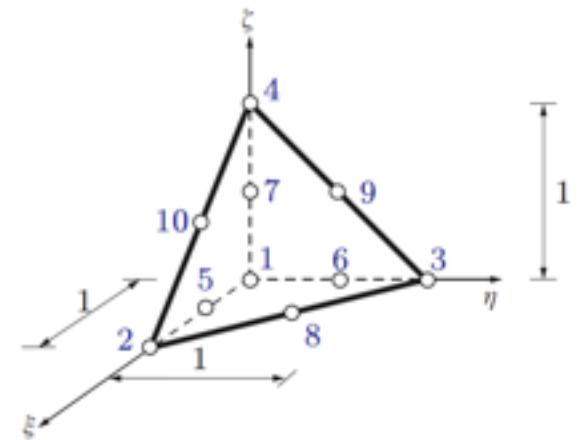
Displacement field

$$\boldsymbol{u} = \sum_{I=1}^{nel} N_I \boldsymbol{d}_u^I = \boldsymbol{N} \boldsymbol{d}_u$$

$$\boldsymbol{N} = \left[\begin{array}{cccc|cccc|cccc} N_1^u & N_2^u & \dots & N_{nel}^u & 0 & 0 & \dots & 0 & 0 & 0 & \dots & 0 \\ 0 & 0 & \dots & 0 & N_1^u & N_2^u & \dots & N_{nel}^u & 0 & 0 & \dots & 0 \\ 0 & 0 & \dots & 0 & 0 & 0 & \dots & 0 & N_1^u & N_2^u & \dots & N_{nel}^u \end{array} \right]_{3 \times 3 \times nel}$$

$$\boldsymbol{d}_u = [d_{ux}^1 \ d_{ux}^2 \ \dots \ d_{ux}^{nel} \ | \ d_{uy}^1 \ d_{uy}^2 \ \dots \ d_{uy}^{nel} \ | \ d_{uz}^1 \ \dots \ d_{uz}^{nel}]^T$$

$$nel = 10$$



Mixed finite element formulation

Displacement field

$$\mathbf{u} = \sum_{I=1}^{nel} N_I \mathbf{d}_u^I = \mathbf{N} \mathbf{d}_u$$

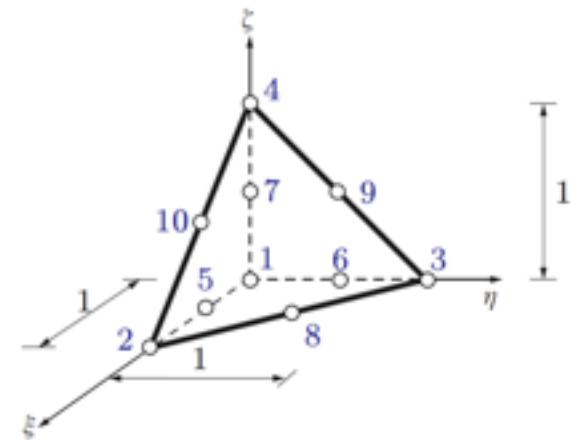
$$\mathbf{N} = \left[\begin{array}{cccc|cccc|cccc} N_1^u & N_2^u & \dots & N_{nel}^u & 0 & 0 & \dots & 0 & 0 & 0 & \dots & 0 \\ 0 & 0 & \dots & 0 & N_1^u & N_2^u & \dots & N_{nel}^u & 0 & 0 & \dots & 0 \\ 0 & 0 & \dots & 0 & 0 & 0 & \dots & 0 & N_1^u & N_2^u & \dots & N_{nel}^u \end{array} \right]_{3 \times 3 \times nel}$$

$$\mathbf{d}_u = [d_{ux}^1 \ d_{ux}^2 \ \dots \ d_{ux}^{nel} \ | \ d_{uy}^1 \ d_{uy}^2 \ \dots \ d_{uy}^{nel} \ | \ d_{uz}^1 \ \dots \ d_{uz}^{nel}]^T$$

$$nel = 10$$

Pressure Terms

$$p_e = \bar{p} \quad \text{and} \quad \theta_e = \bar{\theta} \quad \text{constant}$$



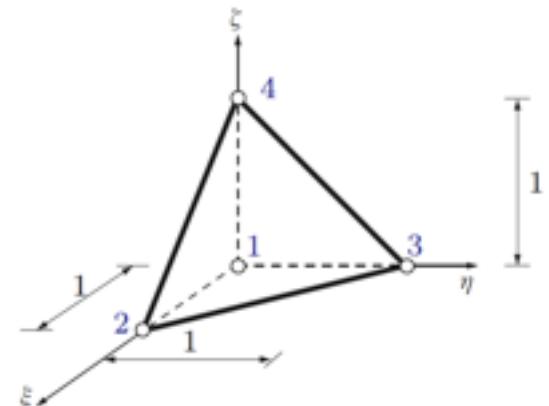
Mixed finite element formulation

Approximation of co factor

$$\mathbf{H}^c = \mathbf{N}_c \mathbf{d}_c$$

$$\begin{bmatrix} H_{11}^c \\ H_{22}^c \\ H_{33}^c \\ H_{12}^c \\ H_{23}^c \\ H_{13}^c \end{bmatrix} = \begin{bmatrix} N_1^c & \dots & N_{nelc}^c \\ 0 & \dots & 0 \end{bmatrix} \begin{bmatrix} 0 & \dots & 0 \\ N_1^c & \dots & N_{nelc}^c \\ 0 & \dots & 0 \end{bmatrix} \dots \begin{bmatrix} 0 & \dots & 0 \\ N_1^c & \dots & N_{nelc}^c \end{bmatrix} \begin{bmatrix} d_{c11}^1 \\ d_{c11}^2 \\ d_{c11}^3 \\ \dots \\ d_{c13}^{(nelc-1)} \\ d_{c13}^{nelc} \end{bmatrix}$$

$$nelc = 4$$



Approximation of Lagrange Multiplier

$$\mathbf{B} = \mathbf{N}_b \mathbf{d}_b$$



Mixed finite element formulation

$$\mathbf{r}_u = \int_{\mathcal{B}_e} \mathbf{B}^T \mathbf{S} dV = \mathbf{0},$$

$$r_\theta = \int_{\mathcal{B}_e} (\partial_\theta \mathcal{P} - \bar{p}) dV = 0,$$

$$r_p = \int_{\mathcal{B}_e} (J - \bar{\theta}) dV = 0,$$

$$\mathbf{r}_c = \int_{\mathcal{B}_e} \mathbf{N}_c^T (\partial_H \mathcal{P} - \mathbf{B}) dV = \mathbf{0},$$

$$\mathbf{r}_b = \int_{\mathcal{B}_e} \mathbf{N}_b^T (\text{cof} \mathbf{C} - \mathbf{H}^c) dV = \mathbf{0}$$



Mixed finite element formulation

$$\mathbf{r}_u = \int_{\mathcal{B}_e} \mathbf{B}^T \mathbf{S} dV = \mathbf{0},$$

$$\mathbf{r}_\theta = \int_{\mathcal{B}_e} (\partial_\theta \mathcal{P} - p) dV = 0,$$

$$\mathbf{r}_p = \int_{\mathcal{B}_e} (J - \bar{\theta}) dV = 0,$$

$$\mathbf{r}_c = \int_{\mathcal{B}_e} \mathbf{N}_c^T (\partial_H \mathcal{P} - \mathbf{B}) dV = \mathbf{0},$$

$$\mathbf{r}_b = \int_{\mathcal{B}_e} \mathbf{N}_b^T (\text{cof} \mathbf{C} - \mathbf{H}^c) dV = \mathbf{0}$$

$$\bar{p} = \frac{1}{V_e} \int_{\mathcal{B}_e} \partial_\theta \mathcal{P} dV \quad \text{and} \quad \bar{\theta} = \frac{1}{V_e} \int_{\mathcal{B}_e} J dV = \frac{v_e}{V_e}$$



Mixed finite element formulation

$$\delta \mathbf{d}_b^T \left(\int_{\mathcal{B}_e} \mathbf{N}_b^T \text{cof} \mathbf{C} dV - \int_{\mathcal{B}_e} \mathbf{N}_b^T \mathbf{N}_c dV \mathbf{d}_c \right) = 0.$$

$$\mathbf{d}_c = (\mathbf{A}_c)^{-1} \mathbf{f}_c$$

$$\mathbf{f}_c := \int_{\mathcal{B}_e} \mathbf{N}_b^T \text{cof} \mathbf{C} dV \quad \text{and} \quad \mathbf{A}_c := \int_{\mathcal{B}_e} \mathbf{N}_b^T \mathbf{N}_b dV$$

Block structure of $\mathbf{H}^c = \mathbf{N}_c \mathbf{d}_c$

$$\begin{bmatrix} \mathbf{M}^e & \mathbf{0} & \dots & \dots \\ \mathbf{0} & \mathbf{M}^e & \dots & \dots \\ \vdots & \vdots & \ddots & \vdots \\ \dots & \dots & \dots & \mathbf{M}^e \end{bmatrix} \begin{Bmatrix} d_{c11} \\ d_{c22} \\ \vdots \\ d_{c13} \end{Bmatrix} = \begin{Bmatrix} f_{c11} \\ f_{c22} \\ \vdots \\ f_{c13} \end{Bmatrix}$$



Mixed finite element formulations

Three-dimensional tetraeder

$$\boldsymbol{M}^e = \frac{V_e}{20} \begin{bmatrix} 2 & 1 & 1 & 1 \\ 1 & 2 & 1 & 1 \\ 1 & 1 & 2 & 1 \\ 1 & 1 & 1 & 2 \end{bmatrix} \quad \Rightarrow \quad (\boldsymbol{M}^e)^{-1} = \frac{4}{V_e} \begin{bmatrix} 4 & -1 & -1 & -1 \\ -1 & 4 & -1 & -1 \\ -1 & -1 & 4 & -1 \\ -1 & -1 & -1 & 4 \end{bmatrix}$$



Local Algorithm at element level

- 1) Update of basic variables: $\boldsymbol{d}_u \Leftarrow \boldsymbol{d}_u + \Delta \boldsymbol{d}_u$

1st Pre-Gauss-loop:

- 2.1) Compute and store at each Gauss point:

$$F_{ij} = \delta_{ij} + \sum_{I=1}^{nel} N_{I,j} d_{u;iI}, \text{ cof}\mathbf{C}, \det\mathbf{F}$$

- 2.2) Calculate volume in the reference and actual configuration:

$$V_e = \int_{\mathcal{B}_e} dV \text{ and } v_e = \int_{\mathcal{S}_e} dv$$

- 2.3) Compute local right hand side regarding \mathbf{H}^c :

$$\mathbf{f}_c = \int_{\mathcal{B}_e} \mathbf{N}_b^T \text{cof}\mathbf{C} dV$$

- 3) Compute coefficient matrix, dofs regarding \mathbf{H}_c and volume dilatation:

\mathbf{M}^e following Equation (59)

$$d_{cij} = (\mathbf{M}^e)^{-1} f_{cij}$$

$$\bar{\theta} = v_e/V_e$$



Local algorithm at element level

- 2nd Pre-Gauss-loop:
- 4)
4.1) Compute and store at each Gausspoint: $\mathbf{H}^c = \mathbf{N}_c \mathbf{d}_c$
4.2) Compute local right hand side regarding \mathbf{B} :
- $$\mathbf{f}_b = \int_{\mathcal{B}_e} \mathbf{N}_c^T (\partial_H \mathcal{P}) dV$$
- 5) Compute pressure $\bar{p} = \partial_\theta \mathcal{P}|_{\theta=\bar{\theta}}$
6) Compute $\mathbf{d}_{bij} = (\mathbf{M}^e)^{-1} \mathbf{f}_{bij}$



Local algorithm at element level

- Final Gauss-loop:
- 7)
 - 7.1) Compute Lagrange multiplier $\mathbf{B} = \mathbf{N}_b \mathbf{d}_b$
 - 7.2) Compute stress $\mathbf{S} = 2(\partial_C \mathcal{P} + \frac{1}{2} \bar{p}/J \text{cof} \mathbf{C} + \mathbf{B} : \partial_C \text{cof} \mathbf{C})$
 - 7.3) Compute right hand side $\mathbf{r}_u = \int_{\mathcal{B}_e} \mathbf{B}^T \mathbf{S} dV$
 - 7.4) Compute individual stiffness matrices:
 $k_{uu}, k_{\theta p}, k_{\theta\theta}, k_{bc}, k_{cc}, k_{bu}, k_{pu}$
 - 8) Calculate static condensation matrices $\tilde{\mathbf{k}}_{pp}, \tilde{\mathbf{k}}_{cc}$
 - 9) Return overall residual vector $\mathbf{r} := \mathbf{r}_u$
and stiffness matrix $\mathbf{k} := \mathbf{k}_{uu} + \mathbf{k}_{up}(\tilde{\mathbf{k}}_{pp})^{-1}\mathbf{k}_{pu} + \mathbf{k}_{ub}\tilde{\mathbf{k}}_{cc}\mathbf{k}_{bu}$



Examples: isotropic strain energy function

$$\psi^{iso}(\mathbf{C}, \text{cof}\mathbf{C}, \det\mathbf{C}) = \frac{\alpha}{2} (\text{tr}\mathbf{C})^2 + \frac{\beta}{2} (\text{tr}[\text{cof}\mathbf{C}])^2 - \gamma \ln(\sqrt{\det \mathbf{C}}) + \epsilon_1(\det\mathbf{C}^{\epsilon_2} + \det\mathbf{C}^{-\epsilon_2} - 2)$$

Mixed form

$$\mathcal{P}^{iso}(\mathbf{C}, \mathbf{H}^c, \theta) = \frac{\alpha}{2} (\text{tr}\mathbf{C})^2 + \frac{\beta}{2} (\text{tr}\mathbf{H}^c)^2 - \gamma \ln(\theta) + \epsilon_1(\theta^{2\epsilon_2} + \theta^{-2\epsilon_2} - 2)$$



Examples: isotropic strain energy function

$$\psi^{iso}(\mathbf{C}, \text{cof}\mathbf{C}, \det\mathbf{C}) = \frac{\alpha}{2} (\text{tr}\mathbf{C})^2 + \frac{\beta}{2} (\text{tr}[\text{cof}\mathbf{C}])^2 - \gamma \ln(\sqrt{\det \mathbf{C}}) + \epsilon_1(\det\mathbf{C}^{\epsilon_2} + \det\mathbf{C}^{-\epsilon_2} - 2)$$

Mixed form

$$\mathcal{P}^{iso}(\mathbf{C}, \mathbf{H}^c, \theta) = \frac{\alpha}{2} (\text{tr}\mathbf{C})^2 + \frac{\beta}{2} (\text{tr}\mathbf{H}^c)^2 - \gamma \ln(\theta) + \epsilon_1(\theta^{2\epsilon_2} + \theta^{-2\epsilon_2} - 2)$$

Stress free initial configuration

$$\gamma = 6\alpha + 12\beta$$



Examples: isotropic strain energy function

$$\psi^{iso}(\mathbf{C}, \text{cof}\mathbf{C}, \det\mathbf{C}) = \frac{\alpha}{2} (\text{tr}\mathbf{C})^2 + \frac{\beta}{2} (\text{tr}[\text{cof}\mathbf{C}])^2 - \gamma \ln(\sqrt{\det \mathbf{C}}) + \epsilon_1 (\det\mathbf{C}^{\epsilon_2} + \det\mathbf{C}^{-\epsilon_2} - 2)$$

Mixed form

$$\mathcal{P}^{iso}(\mathbf{C}, \mathbf{H}^c, \theta) = \frac{\alpha}{2} (\text{tr}\mathbf{C})^2 + \frac{\beta}{2} (\text{tr}\mathbf{H}^c)^2 - \gamma \ln(\theta) + \epsilon_1 (\theta^{2\epsilon_2} + \theta^{-2\epsilon_2} - 2)$$

Stress free initial configuration

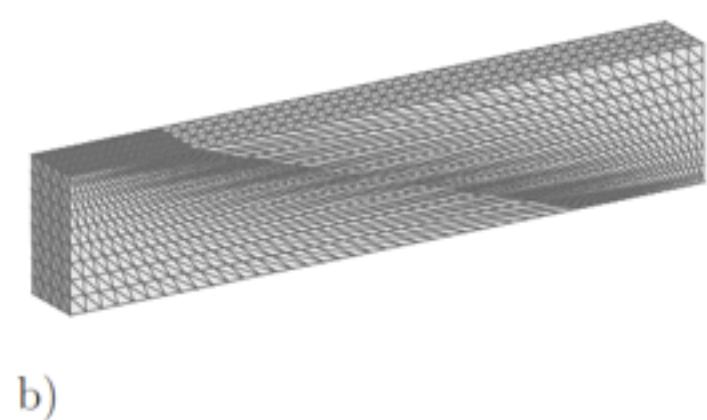
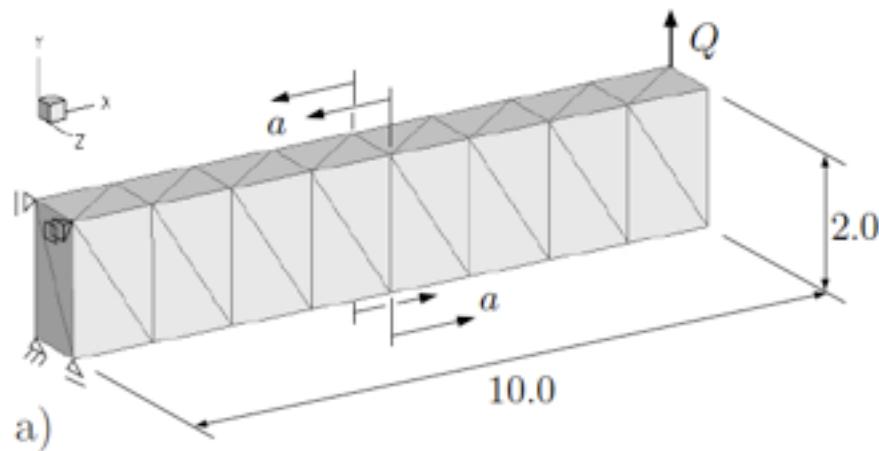
$$\gamma = 6\alpha + 12\beta$$

$$\partial_C \mathcal{P}^{iso} = \alpha (\text{tr}\mathbf{C}) \mathbf{1}, \quad \partial_{H^c} \mathcal{P}^{iso} = \beta (\text{tr}\mathbf{H}^c) \mathbf{1}, \quad \partial_\theta \mathcal{P}^{iso} = -\frac{\gamma}{\theta} + 2\epsilon_1 \epsilon_2 (\theta^{2\epsilon_2-1} - \theta^{-2\epsilon_2-1})$$

$$\partial_{CC}^2 \mathcal{P}^{iso} = \alpha \mathbf{1} \otimes \mathbf{1}, \quad \partial_{H^c H^c}^2 \mathcal{P}^{iso} = \beta \mathbf{1} \otimes \mathbf{1},$$

$$\partial_{\theta\theta}^2 \mathcal{P}^{iso} = \frac{\gamma}{\theta^2} + 2\epsilon_1 \epsilon_2 [2\epsilon_2 (\theta^{2\epsilon_2-2} - \theta^{-2\epsilon_2-2}) - \theta^{2\epsilon_2-2} + \theta^{-2\epsilon_2-2}]$$

Mesh distortion sensitivity: Cantilever beam



$$\alpha = 42.0 \text{ kPa}, \quad \beta = 84.0 \text{ kPa}, \quad \varepsilon_1 = 100 \text{ kPa} \quad \text{and} \quad \varepsilon_2 = 10$$

Mesh distortion sensitivity: Cantilever beam

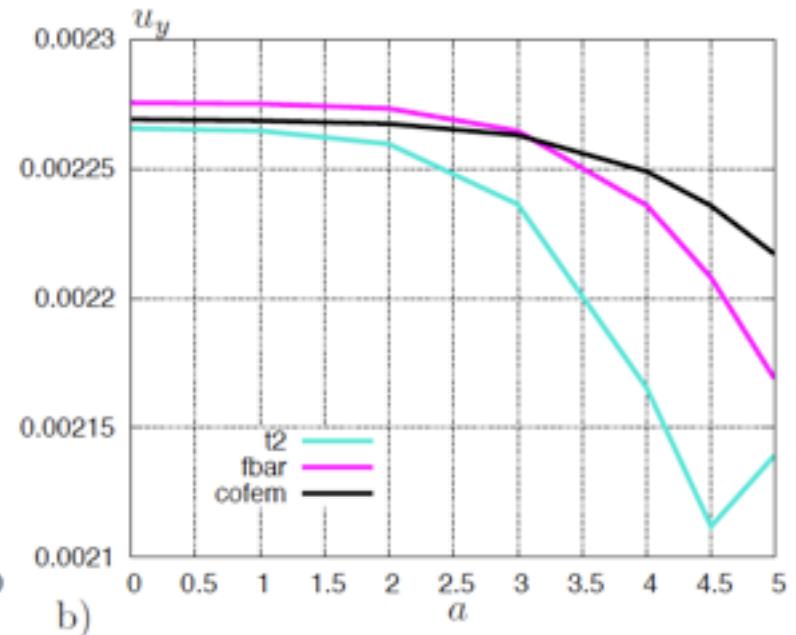
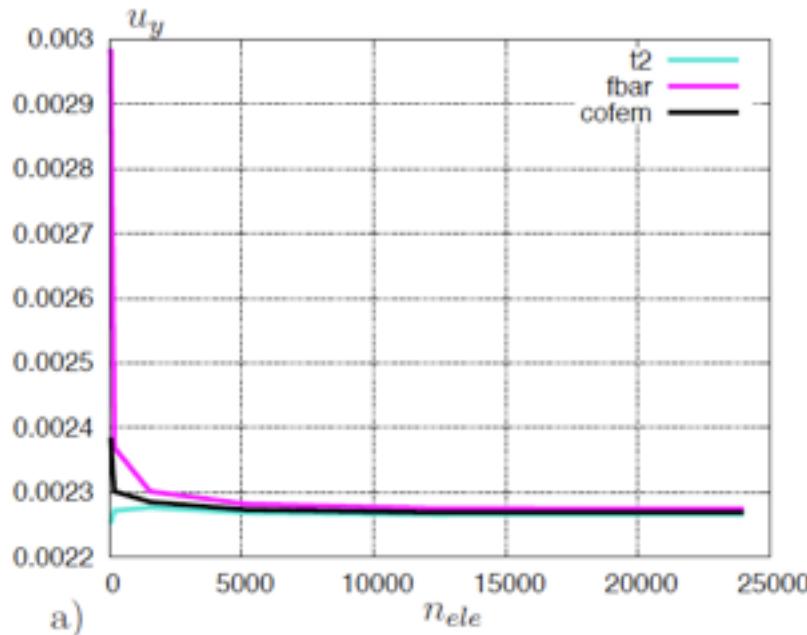


Figure 4: Vertical displacement at the loading point versus a) number of elements n_{ele} (for $a = 0$) and b) distortion a (for $n_{ele} = 12288$).

Mesh distortion sensitivity: Cantilever beam

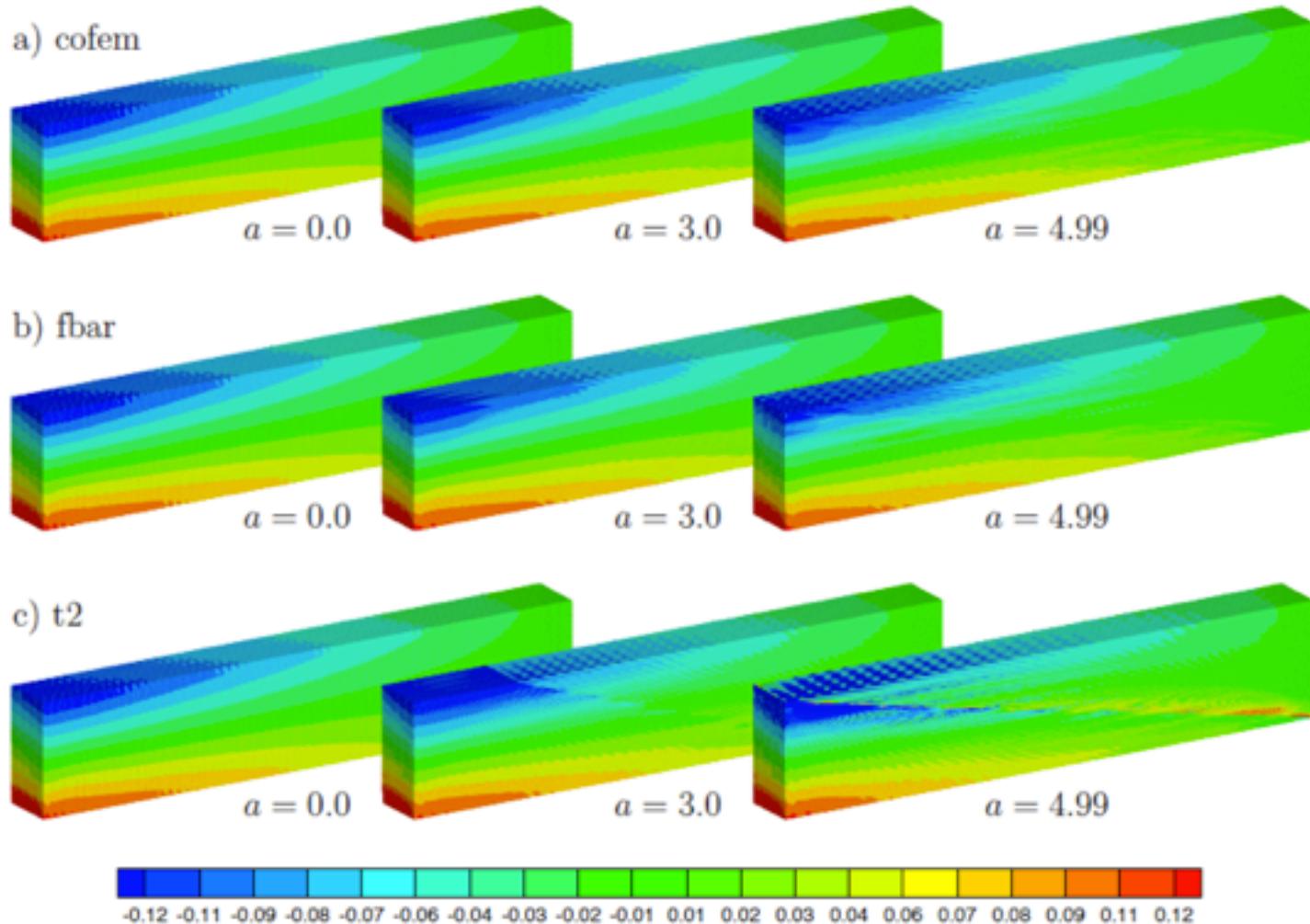


Figure 5: Distribution of normal Cauchy stress in x -direction σ_{11} for the a) cofem-element, b) fbar-element and c) t2-element at different distortion values.

Cook's membran problem

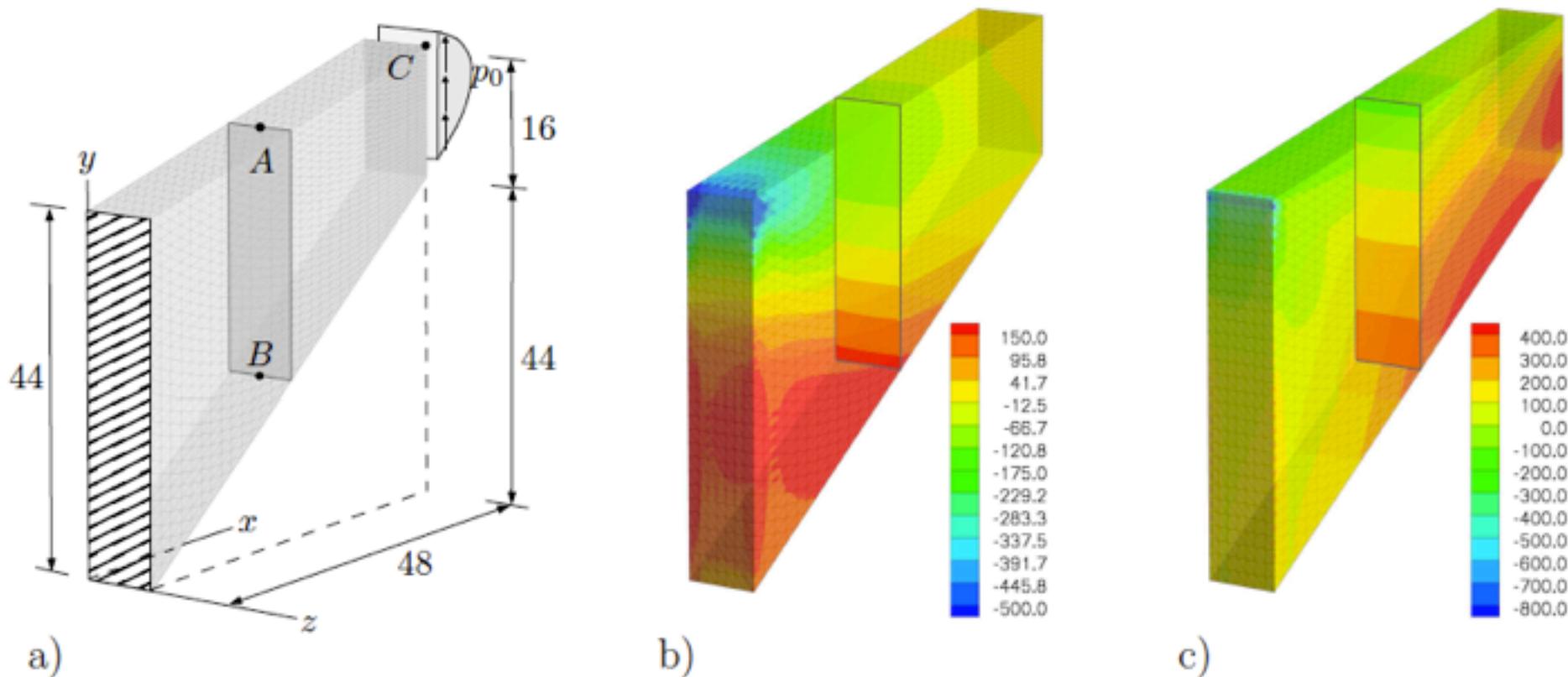


Figure 6: Illustration of the a) boundary value problem and the resulting distribution of normal stresses b) σ_{xx} - and c) σ_{yy} . The considered mesh consists of 32928 tetrahedral elements using the proposed formulation.

Cook's membran problem

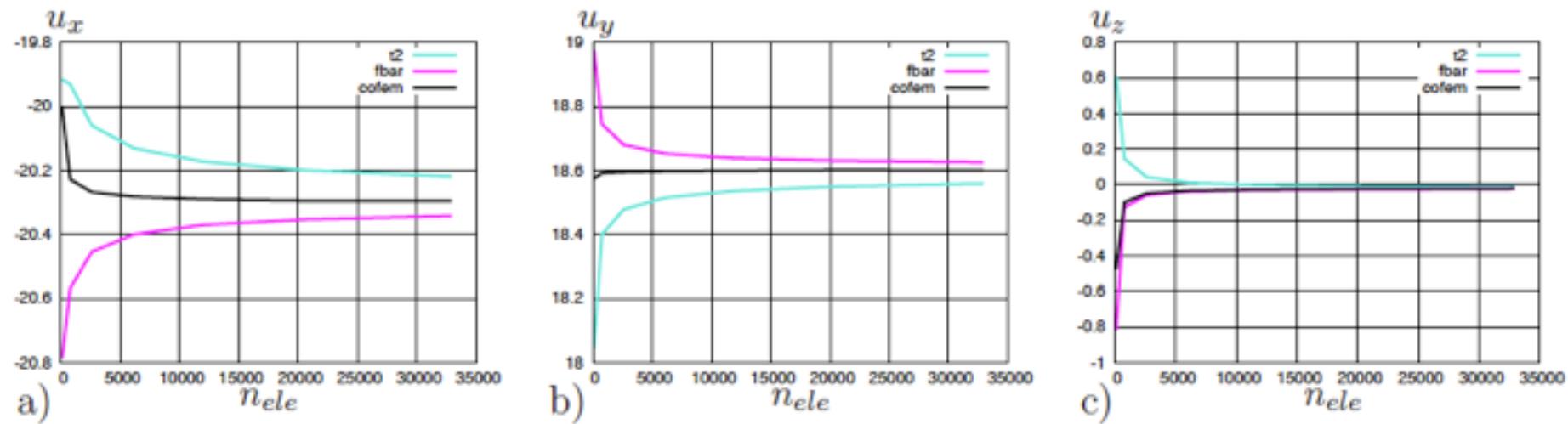


Figure 7: Displacements a) u_x and b) u_y and c) u_z at point C versus number of elements using the isotropic strain energy function \mathcal{P}^{iso} .

Cook's membran problem

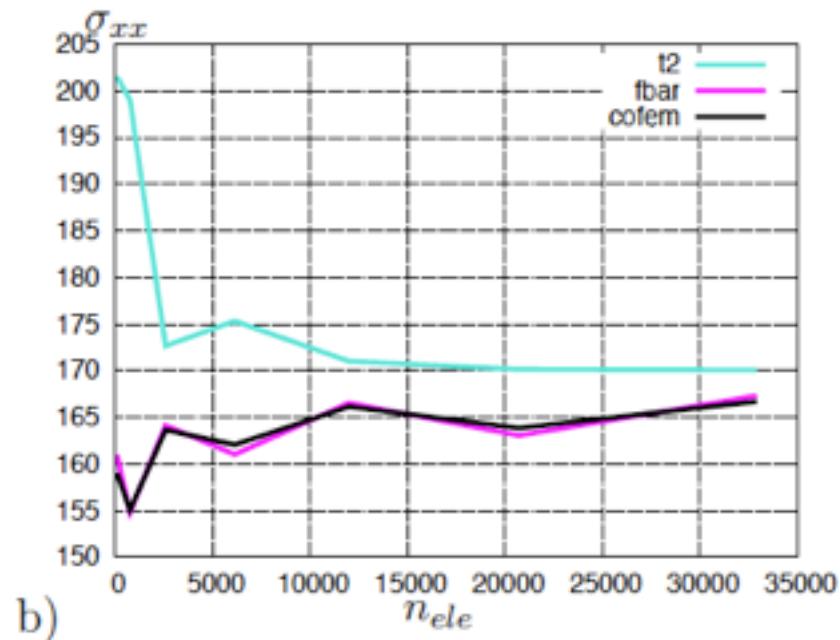
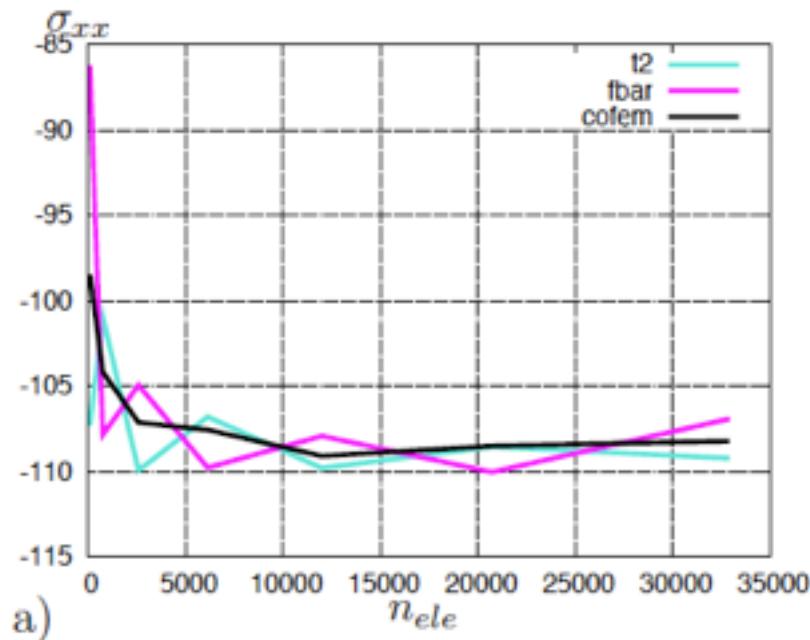


Figure 8: Stress values of σ_{xx} at a) point A and b) point B using the isotropic strain energy function \mathcal{P}^{iso} and three different FE-formulations.

Cook's membran problem

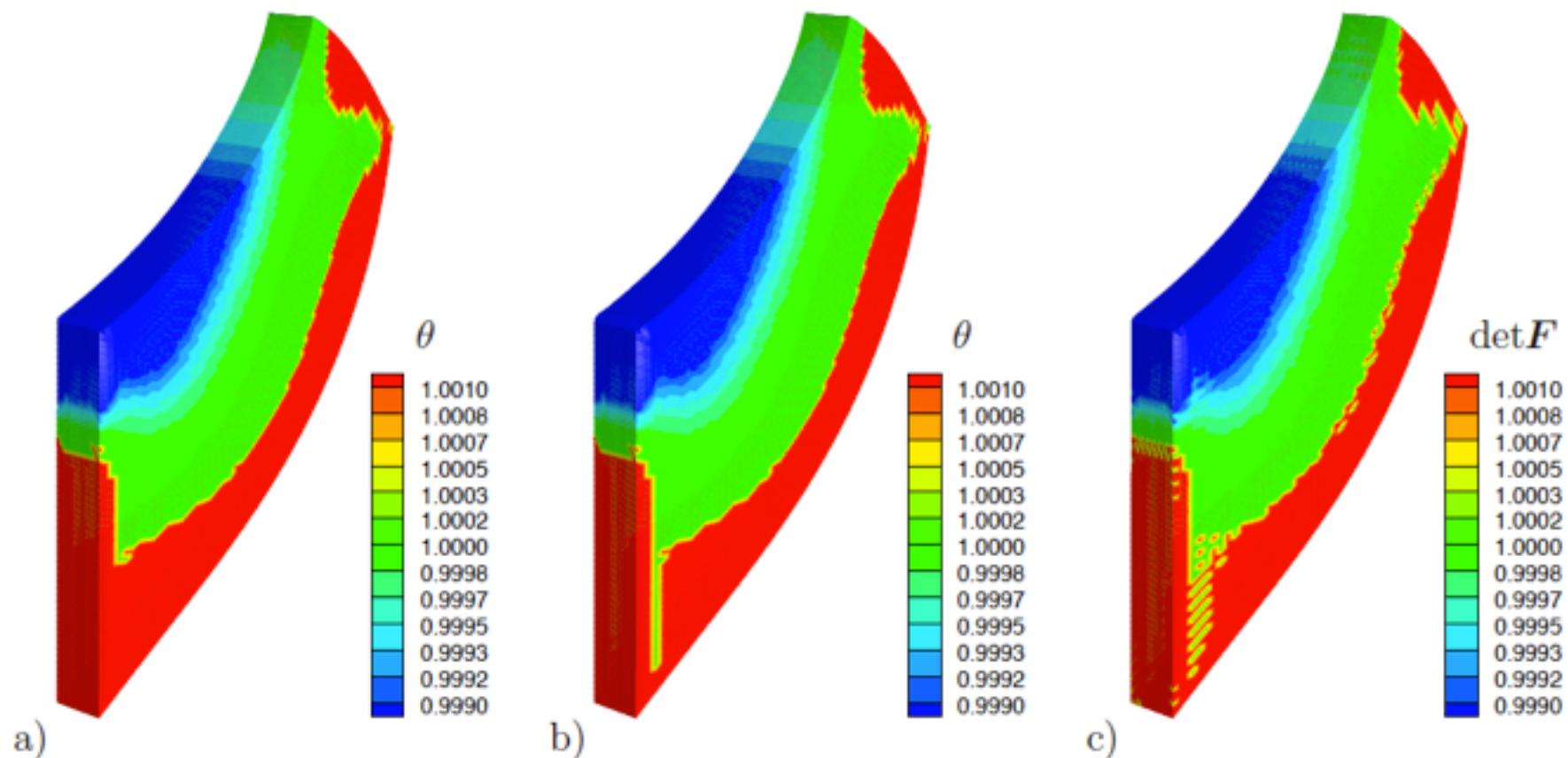


Figure 9: Distribution of $\theta/\det F$ in the deformed configuration for the a) cofem-, b) fbar- and c) t2 formulation; the number of elements is $n_{ele} = 32928$.



Anisotropic strain energy function

$$\begin{aligned}\psi^{ti}(C, \text{cof}C, \det C) = & g_0 \left[\frac{1}{g_C + 1} (\text{tr}[CM])^{g_C+1} \right. \\ & \left. + \frac{1}{g_H + 1} (\text{tr}[\text{cof}[C]M])^{g_H+1} + \frac{1}{g_\theta} (\det C)^{-g_\theta} \right]\end{aligned}$$

$$M = a \otimes a$$

Schröder et al. 2009



Anisotropic strain energy function

$$\begin{aligned}\psi^{ti}(C, \text{cof}C, \det C) = & g_0 \left[\frac{1}{g_C + 1} (\text{tr}[CM])^{g_C+1} \right. \\ & \left. + \frac{1}{g_H + 1} (\text{tr}[\text{cof}[C]M])^{g_H+1} + \frac{1}{g_\theta} (\det C)^{-g_\theta} \right]\end{aligned}$$

$$M = a \otimes a$$

Schröder et al. 2009

Mixed form

$$\mathcal{P}^{ti}(C, H^c, \theta) = g_0 \left[\frac{1}{g_C + 1} (\text{tr}[CM])^{(g_C+1)} + \frac{1}{g_H + 1} (\text{tr}[H^c M])^{(g_H+1)} + \frac{1}{g_\theta} (\theta)^{-2g_\theta} \right]$$



Anisotropic strain energy function

$$\begin{aligned}\psi^{ti}(C, \text{cof}C, \det C) = & g_0 \left[\frac{1}{g_C + 1} (\text{tr}[CM])^{g_C+1} \right. \\ & \left. + \frac{1}{g_H + 1} (\text{tr}[\text{cof}[C]M])^{g_H+1} + \frac{1}{g_\theta} (\det C)^{-g_\theta} \right]\end{aligned}$$

$$M = a \otimes a$$

Schröder et al. 2009

Mixed form

$$\mathcal{P}^{ti}(C, H^c, \theta) = g_0 \left[\frac{1}{g_C + 1} (\text{tr}[CM])^{(g_C+1)} + \frac{1}{g_H + 1} (\text{tr}[H^c M])^{(g_H+1)} + \frac{1}{g_\theta} (\theta)^{-2g_\theta} \right]$$

$$\partial_C \mathcal{P}^{ti} = g_0 (\text{tr}[CM])^{g_C} M, \quad \partial_{H^c} \mathcal{P}^{ti} = g_0 (\text{tr}[H^c M])^{g_H} M.$$

$$\partial_\theta \mathcal{P}^{ti} = -2 g_0 \theta^{(-2g_\theta-1)}$$



Anisotropic strain energy function

$$\begin{aligned}\psi^{ti}(C, \text{cof}C, \det C) = & g_0 \left[\frac{1}{g_C + 1} (\text{tr}[CM])^{g_C+1} \right. \\ & \left. + \frac{1}{g_H + 1} (\text{tr}[\text{cof}[C]M])^{g_H+1} + \frac{1}{g_\theta} (\det C)^{-g_\theta} \right]\end{aligned}$$

$$M = a \otimes a$$

Schröder et al. 2009

Mixed form

$$\mathcal{P}^{ti}(C, H^c, \theta) = g_0 \left[\frac{1}{g_C + 1} (\text{tr}[CM])^{(g_C+1)} + \frac{1}{g_H + 1} (\text{tr}[H^c M])^{(g_H+1)} + \frac{1}{g_\theta} (\theta)^{-2g_\theta} \right]$$

$$\partial_C \mathcal{P}^{ti} = g_0 (\text{tr}[CM])^{g_C} M, \quad \partial_{H^c} \mathcal{P}^{ti} = g_0 (\text{tr}[H^c M])^{g_H} M.$$

$$\partial_\theta \mathcal{P}^{ti} = -2 g_0 \theta^{(-2g_\theta-1)}$$

$$g_0 = 3000.0 \text{ kPa}, \quad g_C = 4.0, \quad g_H = 8 \quad \text{and} \quad g_\theta = 1$$

Cook's membran problem

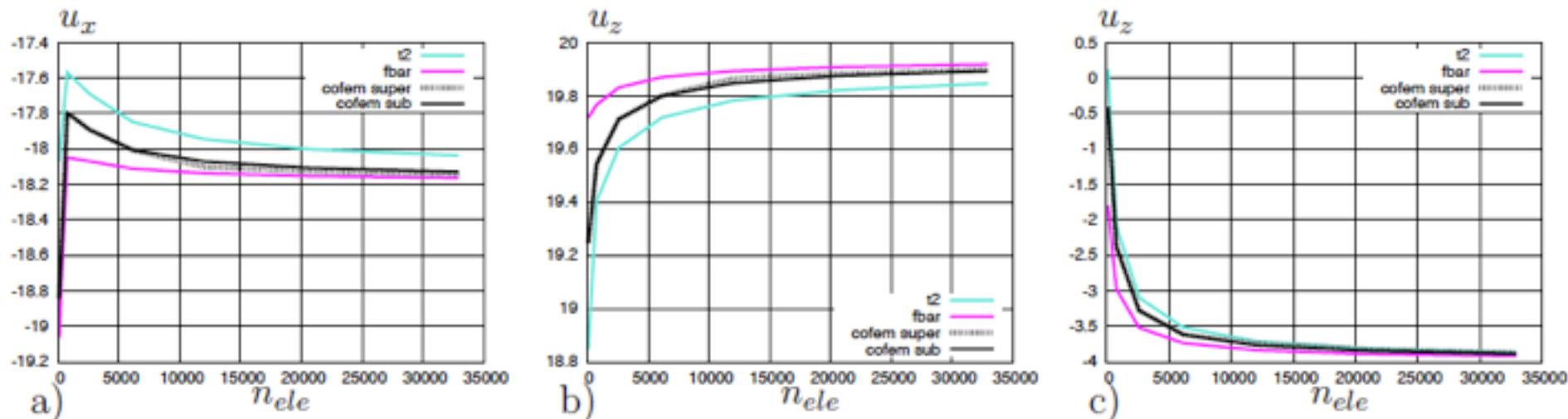


Figure 11: Displacements a) u_x and b) u_y and u_z at point C versus number of elements using the transversely isotropic strain energy function.

Cook's membran problem

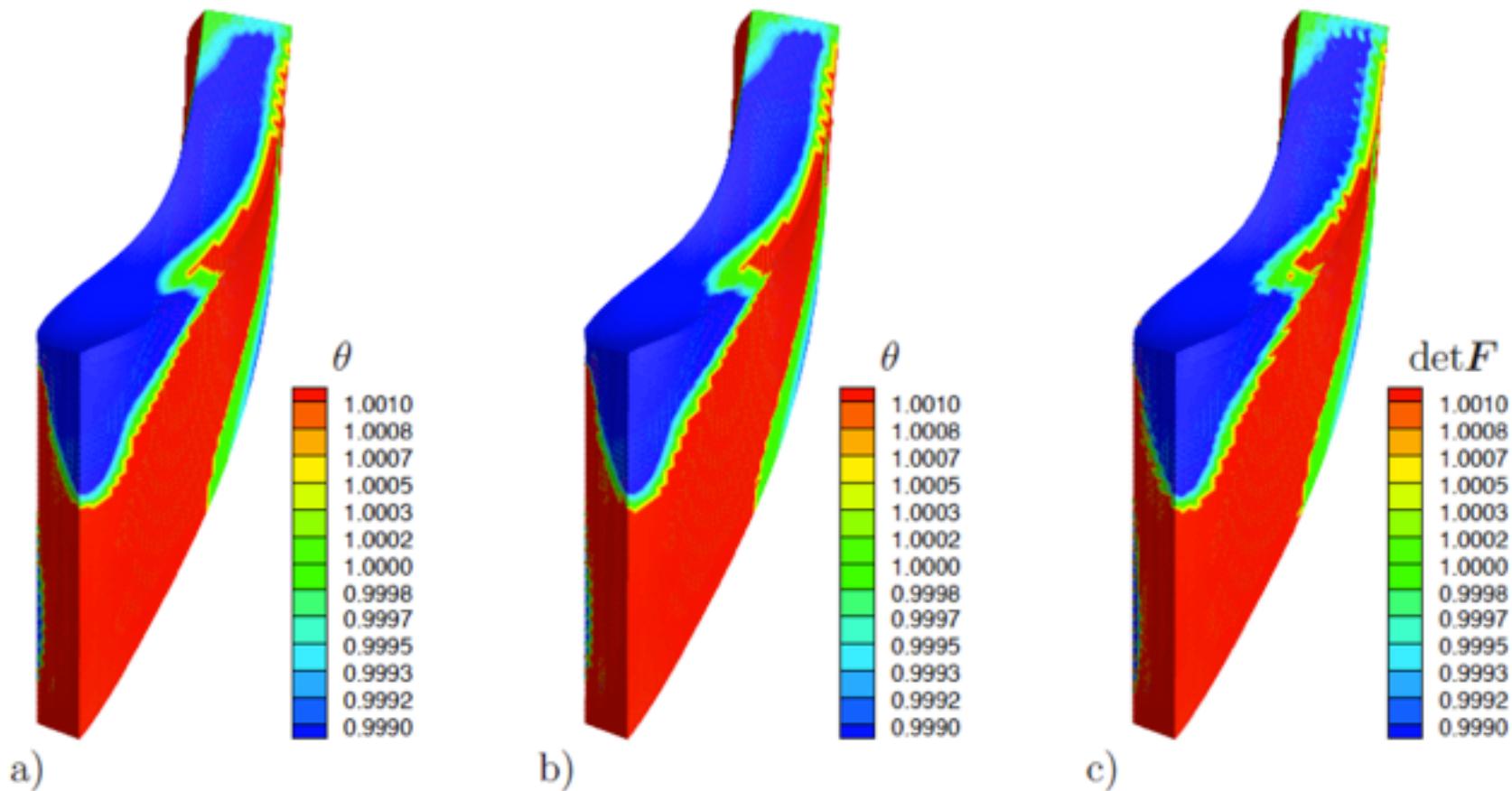


Figure 13: Distribution of $\theta/\det F$ in the deformed configuration for the a) cofem-, b) fbar- and c) t2 formulation using the transversely isotropic strain energy function $\mathcal{P} := \mathcal{P}^{iso} + \mathcal{P}^{ti}$; the number of elements is $n_{ele} = 32928$.

Simulation of arterial walls

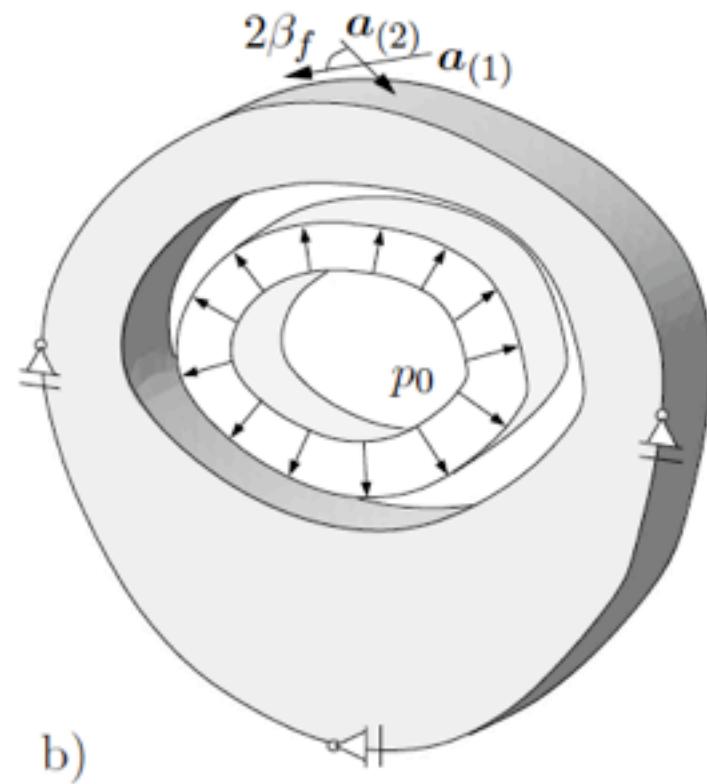
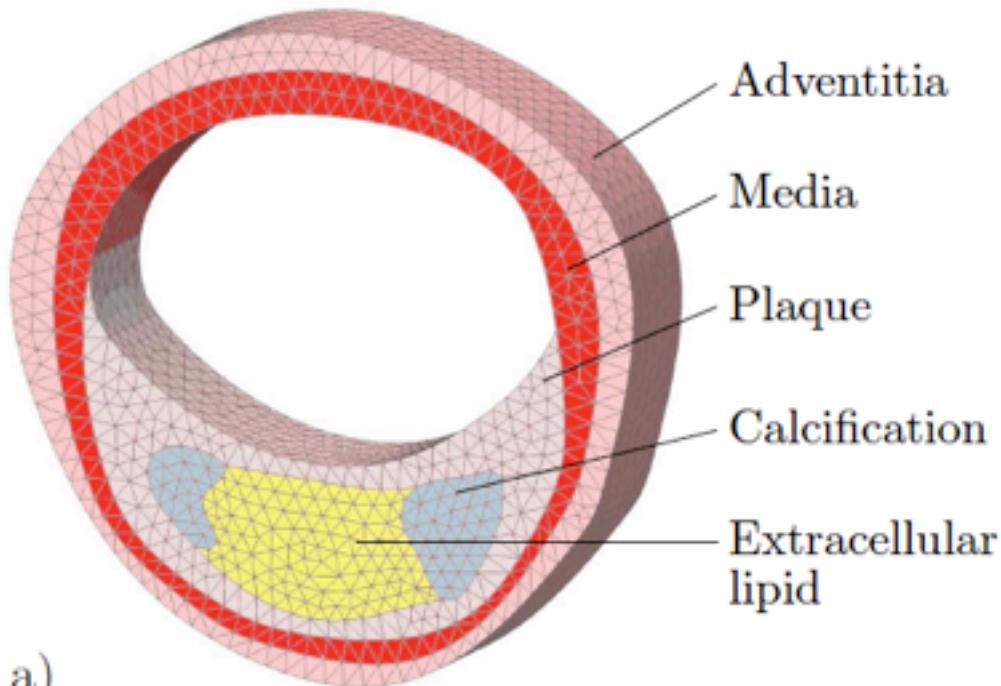


Figure 15: Illustration of the a) individual healthy and deceased tissue components and b) the boundary value problem; the considered mesh consists of 20860 tetrahedral elements.



Simulation of arterial walls

Isotropic material model for calcification and extracellular lipid

$$\mathcal{P}^{iso}(\mathbf{C}, \mathbf{H}^c, \theta) = c_1 \left(\frac{\text{tr} \mathbf{C}}{(\det \mathbf{C})^{1/3}} - 3 \right) + \epsilon_1 \left(\theta^{2\epsilon_2} + \frac{1}{\theta^{2\epsilon_2}} - 2 \right)$$

Transversely isotropic material model for media and adventitia

$$\mathcal{P}_{(a)}^{ti}(\mathbf{C}, \mathbf{H}^c, \theta) = \alpha_1 \langle \text{tr} [(\mathbf{H}^c)(1 - \mathbf{M}_{(a)})] - 2 \rangle^{\alpha_2}$$

Complete strain energy function for media and adventitia

$$\mathcal{P}^{bio} = \mathcal{P}^{iso} + \sum_{a=1}^2 \mathcal{P}_{(a)}^{ti}$$

Isotropic material model for plaque

$$\mathcal{P}^{MR}(\mathbf{C}, \mathbf{H}^c, \theta) = \beta_1 \text{tr} \mathbf{C} + \eta_1 \text{tr} (\mathbf{H}^c) + \delta_1 \theta^2 - \delta_2 \ln(\theta^2)$$

Schröder & Neff 2004, Balzani 2006, Brands et al. 2008



Simulation of arterial walls

Table 1: Material parameters for the individual tissues taken from Brands et al. (2008).

Layer	Model	Set	c_1 [kPa]	ε_1 [kPa]	ε_2 [-]	α_1 [kPa]	α_2 [-]	β_1 [kPa]	η_1 [kPa]	δ_1 [kPa]	β_f [°]
Media	\mathcal{P}^{bio}	1	17.5	100.0	50.0	5.0 E05	7.0	–	–	–	43.39
		2	17.5	499.8	2.4	30001.9	5.1	–	–	–	43.39
Adventitia	\mathcal{P}^{bio}	1	7.5	100.0	20.0	1.5 E10	20.0	–	–	–	49.0
		2	6.6	23.9	10.0	1503.0	6.3	–	–	–	49.0
Calcification	\mathcal{P}^{iso}	1,2	6800.0	50.0	10.0	–	–	–	–	–	–
Lipid	\mathcal{P}^{iso}	1,2	700.0	5250.0	10.0	–	–	–	–	–	–
Plaque	\mathcal{P}^{MR}	1,2	–	–	–	–	–	80.0	250.0	2000.0	–

Simulation of arterial walls

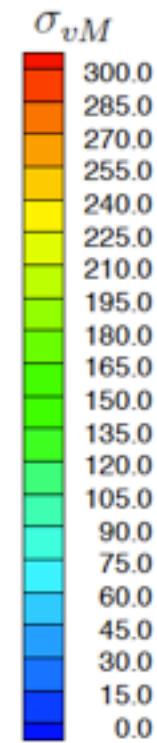
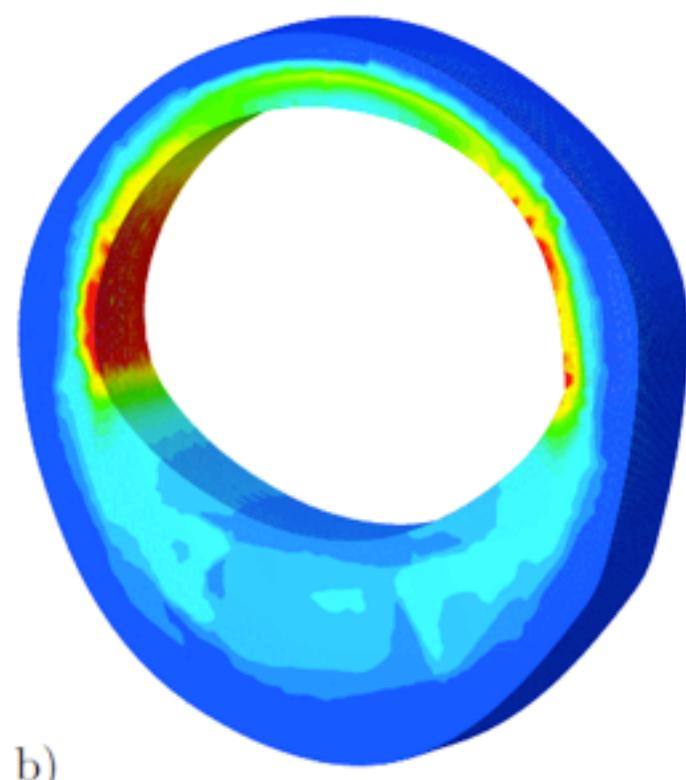
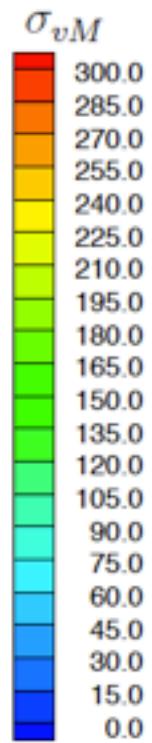
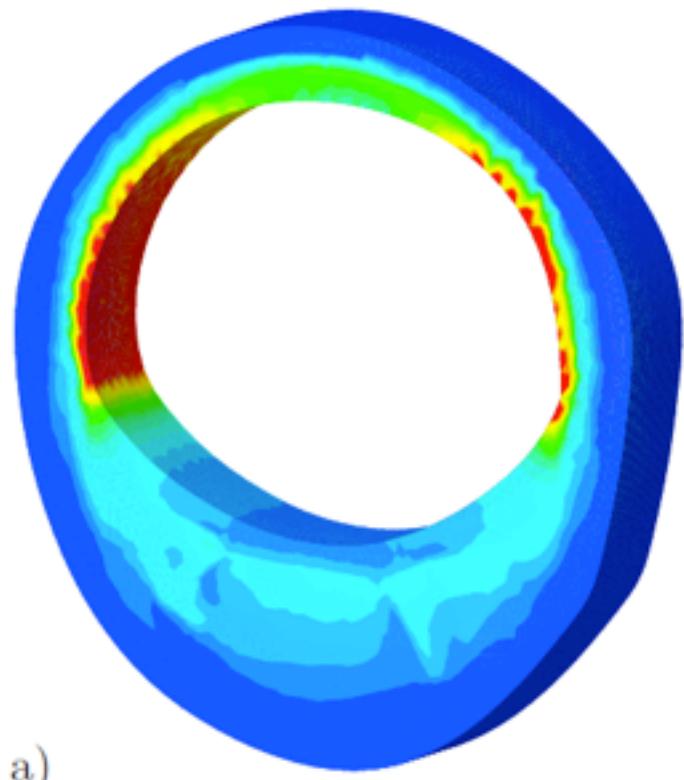


Figure 17: Distribution of the von Mises stresses σ_{vM} in the deformed configuration for the a) cofem- and b) fbar-formulation using parameter set 2.

Development of the EI9 element

- Motivation
- Split of the strain energy function
 - Different Treatment for homogenous and inhomogenous part
- Variational formulation
- Ansatz and Implementation
- Numerical Tests
 - non-uniformly meshed beam
 - incompressible block
 - surface buckling
- Conclusion and outlook



Split of the Strain Energy Function

Additive split into a homogeneous and an inhomogeneous part (Nadler & Rubin, 2003)

$$W(\mathbf{F}) = W_{\text{hom}}(\bar{\mathbf{F}}) + W_{\text{inh}}(\hat{\tilde{\mathbf{F}}})$$

Split of the deformation gradient:

homogeneous part

$$\mathbf{F} = \bar{\mathbf{F}} \cdot \hat{\tilde{\mathbf{F}}} \quad ,$$

$$\bar{\mathbf{F}} = \frac{1}{V} \int_{\Omega} \mathbf{F} d\Omega$$



Homogeneous Part

Compressible Neo-Hooke Material

$$W_{\text{hom}}(\bar{\mathbf{F}}) = \frac{\mu}{2} \left(\bar{J}^{-2/3} \operatorname{tr}(\bar{\mathbf{C}}) - 3 \right) + \frac{K}{\beta^2} (\bar{J}^{-\beta} - 1 + \beta \ln(\bar{J}))$$

$$\begin{aligned}\bar{J} &= \det(\bar{\mathbf{F}}) \\ \bar{\mathbf{C}} &= \bar{\mathbf{F}}^T \bar{\mathbf{F}} \quad .\end{aligned}$$

$$\bar{\mathbf{P}} = \frac{\partial W_{\text{hom}}}{\partial \bar{\mathbf{F}}} = \mu \bar{J}^{-2/3} \left(\bar{\mathbf{F}} - \frac{1}{3} \operatorname{tr}(\bar{\mathbf{C}}) \bar{\mathbf{F}}^{-T} \right) + \frac{K}{\beta} (1 - \bar{J}^{-\beta}) \bar{\mathbf{F}}^{-T}$$



Inhomogeneous Part (1)

Linear Elastic Material:

$$W_{\text{inh}}(\hat{\tilde{\mathbf{F}}}) = \frac{1}{2} (\hat{\tilde{\mathbf{F}}} - \mathbf{1}) : \mathbb{C} : (\hat{\tilde{\mathbf{F}}} - \mathbf{1})$$

$$\hat{\tilde{\mathbf{P}}} = (K - \frac{2}{3}\mu) (\text{tr}(\hat{\tilde{\mathbf{F}}}) - 3) \mathbf{1} + \mu (\hat{\tilde{\mathbf{F}}} + \hat{\tilde{\mathbf{F}}}^T - 2\mathbf{1})$$

with

$$\mathbb{C}_{iklm} = (K - \frac{2}{3}\mu) \delta_{ik}\delta_{lm} + \mu (\delta_{il}\delta_{km} + \delta_{im}\delta_{kl})$$

Inhomogeneous Part (2)

Multiplicative Split of the inhomogeneous part of the displacement gradient

$$\hat{\tilde{F}} = \tilde{F} + \hat{H} \quad \text{with} \quad \tilde{F} = \bar{F}^{-1} \cdot F(x)$$

leading to

$$F = \bar{F} \cdot \left(\bar{F}^{-1} \cdot F(x) + \hat{H} \right) = F(x) + \bar{F} \cdot \hat{H}$$

Ansatz for inhomogeneous part of H leading to \hat{H}

$$M_1 = (1 - \xi^2) \quad M_2 = (1 - \eta^2) \quad M_3 = (1 - \zeta^2)$$

(Wilson, Taylor, Doherty & Ghaboussi, 1973)



Potential and Variation

Hu-Washizu

$$\Pi(x, \hat{H}, P) = \int_{\Omega} [W_{\text{hom}}(\bar{F}) + W_{\text{inh}}(\hat{\tilde{F}}) - P : \bar{F} \cdot \hat{H}] \, d\Omega - P_{ext}$$



Potential and Variation

Hu-Washizu

$$\Pi(x, \hat{\mathbf{H}}, \mathbf{P}) = \int_{\Omega} \left[W_{\text{hom}}(\bar{\mathbf{F}}) + W_{\text{inh}}(\hat{\tilde{\mathbf{F}}}) - \mathbf{P} : \bar{\mathbf{F}} \cdot \hat{\mathbf{H}} \right] d\Omega - P_{ext}$$

$$\int_{\Omega} \delta \bar{\mathbf{F}} : \frac{\partial W_{\text{hom}}}{\partial \bar{\mathbf{F}}} d\Omega + \int_{\Omega} \delta \tilde{\mathbf{F}} : \frac{\partial W_{\text{inh}}}{\partial \hat{\tilde{\mathbf{F}}}} d\Omega - \int_{\Omega} \mathbf{P} : \delta \bar{\mathbf{F}} \cdot \hat{\mathbf{H}} d\Omega - \delta P_{ext} = 0$$

$$\int_{\Omega} \delta \hat{\mathbf{H}} : \frac{\partial W_{\text{inh}}}{\partial \hat{\tilde{\mathbf{F}}}} d\Omega - \int_{\Omega} \mathbf{P} : \bar{\mathbf{F}} \cdot \delta \hat{\mathbf{H}} d\Omega = 0$$

$$\int_{\Omega} \delta \mathbf{P} : \bar{\mathbf{F}} \cdot \hat{\mathbf{H}} d\Omega = 0$$



Potential and Variation

Hu-Washizu

$$\Pi(x, \hat{\mathbf{H}}, \mathbf{P}) = \int_{\Omega} \left[W_{\text{hom}}(\bar{\mathbf{F}}) + W_{\text{inh}}(\hat{\tilde{\mathbf{F}}}) - \mathbf{P} : \bar{\mathbf{F}} \cdot \hat{\mathbf{H}} \right] d\Omega - P_{ext}$$

$$\int_{\Omega} \delta \bar{\mathbf{F}} : \frac{\partial W_{\text{hom}}}{\partial \bar{\mathbf{F}}} d\Omega + \int_{\Omega} \delta \tilde{\mathbf{F}} : \frac{\partial W_{\text{inh}}}{\partial \hat{\tilde{\mathbf{F}}}} d\Omega - \int_{\Omega} \mathbf{P} : \delta \bar{\mathbf{F}} \cdot \hat{\mathbf{H}} d\Omega - \delta P_{ext} = 0$$

$$\int_{\Omega} \delta \hat{\mathbf{H}} : \frac{\partial W_{\text{inh}}}{\partial \hat{\tilde{\mathbf{F}}}} d\Omega - \int_{\Omega} \mathbf{P} : \bar{\mathbf{F}} \cdot \delta \hat{\mathbf{H}} d\Omega = 0$$

$$\int_{\Omega} \delta \mathbf{P} : \bar{\mathbf{F}} \cdot \hat{\mathbf{H}} d\Omega = 0$$

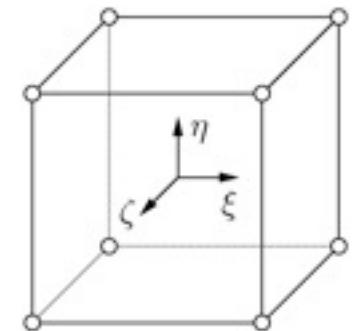
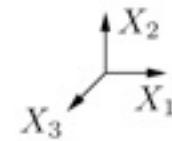
with

$$\delta \tilde{\mathbf{F}} = -\bar{\mathbf{F}}^{-1} \cdot \delta \bar{\mathbf{F}} \cdot \bar{\mathbf{F}}^{-1} \cdot \mathbf{F}(x) + \bar{\mathbf{F}}^{-1} \cdot \delta \mathbf{F}(x)$$

Ansatz and Implementation

$$\boldsymbol{x}^h = \sum_{I=1}^N N_I \boldsymbol{x}_I = \sum_{I=1}^N N_I (\boldsymbol{X}_I + \boldsymbol{u}_I)$$

$$N_I(\xi, \eta, \zeta) = \frac{1}{8}(1 + \xi\xi_I)(1 + \eta\eta_I)(1 + \zeta\zeta_I)$$





$$\boldsymbol{F}^h = \begin{pmatrix} F_{11}^h \\ F_{22}^h \\ F_{33}^h \\ F_{12}^h \\ F_{21}^h \\ F_{23}^h \\ F_{32}^h \\ F_{13}^h \\ F_{31}^h \end{pmatrix} = \sum_{I=1}^8 \boldsymbol{B}_I \boldsymbol{x}_I , \quad \boldsymbol{B}_I = \begin{pmatrix} N_{I,X} & 0 & 0 \\ 0 & N_{I,Y} & 0 \\ 0 & 0 & N_{I,Z} \\ N_{I,Y} & 0 & 0 \\ 0 & N_{I,X} & 0 \\ 0 & N_{I,Z} & 0 \\ 0 & 0 & N_{I,Y} \\ N_{I,Z} & 0 & 0 \\ 0 & 0 & N_{I,X} \end{pmatrix}$$

$$\bar{\boldsymbol{F}}^h = \frac{1}{\Omega_e} \int_{\Omega_e} \boldsymbol{F}^h \, d\Omega = \sum_{I=1}^8 \frac{1}{\Omega_e} \int_{\Omega_e} \boldsymbol{B}_I \, d\Omega \, \boldsymbol{x}_I = \sum_{I=1}^8 \bar{\boldsymbol{B}}_I \, \boldsymbol{x}_I$$



$$\hat{\mathbf{H}}^h = \begin{pmatrix} \hat{H}_{11}^h \\ \hat{H}_{22}^h \\ \hat{H}_{33}^h \\ \hat{H}_{12}^h \\ \hat{H}_{21}^h \\ \hat{H}_{23}^h \\ \hat{H}_{32}^h \\ \hat{H}_{13}^h \\ \hat{H}_{31}^h \end{pmatrix} = \sum_{L=1}^4 \mathbf{G}_L \boldsymbol{\alpha}_L , \quad \mathbf{G}_L = \begin{pmatrix} M_{L,X} & 0 & 0 \\ 0 & M_{L,Y} & 0 \\ 0 & 0 & M_{L,Z} \\ M_{L,Y} & 0 & 0 \\ 0 & M_{L,X} & 0 \\ 0 & M_{L,Z} & 0 \\ 0 & 0 & M_{L,Y} \\ M_{L,Z} & 0 & 0 \\ 0 & 0 & M_{L,X} \end{pmatrix} \quad \mathbf{G}_4 = \begin{pmatrix} M_{4,X} & M_{4,Y} & M_{4,Z} \\ M_{4,X} & M_{4,Y} & M_{4,Z} \\ M_{4,X} & M_{4,Y} & M_{4,Z} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

with

$$M_1 = (1 - \xi^2) \quad M_2 = (1 - \eta^2) \quad M_3 = (1 - \zeta^2) \quad M_4 = \xi\eta\zeta$$

weak form

$$\bigcup_{e=1}^{n_e} \left(\sum_{I=1}^8 \delta \mathbf{u}_I^T \bar{\mathbf{B}}_I^T \bar{\mathbf{P}}^h \Omega_e + \sum_{I=1}^8 \delta \mathbf{u}_I^T \int_{\Omega_e} \tilde{\mathbf{B}}_I^T \hat{\mathbf{P}}^h d\Omega \right) - \delta P_{\text{ext}}^h = 0$$

$$\sum_{K=1}^4 \delta \boldsymbol{\alpha}_K^T \int_{\Omega_e} \mathbf{G}_K^T \hat{\mathbf{P}}^h d\Omega = 0$$

Residual and tangent stiffness matrix

$$\mathbf{R}_I^u = \bar{\mathbf{B}}_I^T \bar{\mathbf{P}}^h \Omega_e + \int_{\Omega_e} \tilde{\mathbf{B}}_I^T \hat{\mathbf{P}}^h d\Omega - \mathbf{P}_I^{ext}$$

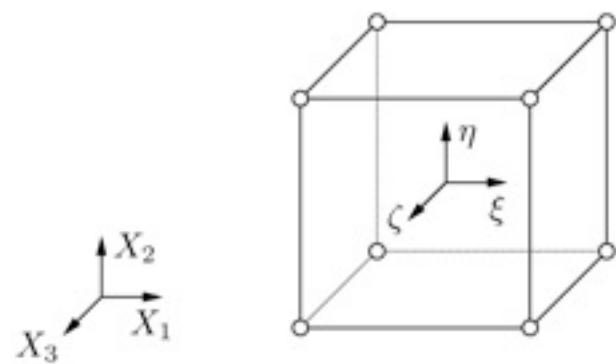
$$\mathbf{R}_K^\alpha = \int_{\Omega_e} \mathbf{G}_K^T \hat{\mathbf{P}}^h d\Omega$$

$$\mathbf{K}_{IJ}^{uu} = \bar{\mathbf{B}}_I^T \bar{\mathbb{D}} \bar{\mathbf{B}}_J \Omega_e + \int_{\Omega_e} \tilde{\mathbf{B}}_I^T \mathbb{C} \tilde{\mathbf{B}}_J d\Omega + \mathbf{K}_{IJ}^{uu\,geo}$$

$$\mathbf{K}_{IL}^{u\alpha} = \int_{\Omega_e} \tilde{\mathbf{B}}_I^T \mathbb{C} \mathbf{G}_L d\Omega$$

$$\mathbf{K}_{KJ}^{\alpha u} = \int_{\Omega_e} \mathbf{G}_K^T \mathbb{C} \tilde{\mathbf{B}}_J d\Omega$$

$$\mathbf{K}_{KL}^{\alpha\alpha} = \int_{\Omega_e} \mathbf{G}_K^T \mathbb{C} \mathbf{G}_L d\Omega$$



$$\delta \mathbf{u}_I \cdot \mathbf{K}_{IJ}^{uu\,geo} \cdot \Delta \mathbf{u}_J = \int_{\Omega_e} \Delta \delta \tilde{\mathbf{F}} : \mathbb{C} : (\hat{\tilde{\mathbf{F}}} - \mathbf{1}) d\Omega$$

Equation system on element level

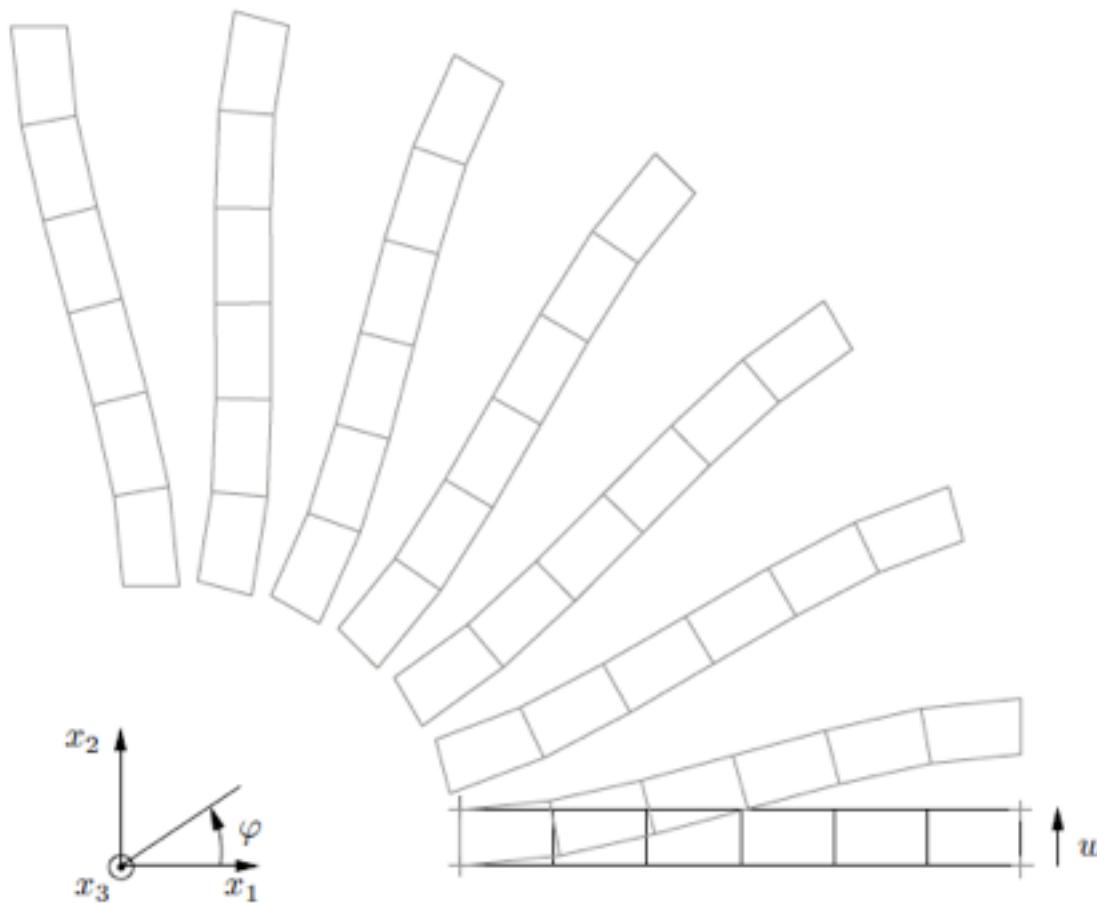
$$\mathbf{K}^{uu} \Delta \mathbf{u} + \mathbf{K}^{u\alpha} \Delta \alpha = -\mathbf{R}^u$$

$$\mathbf{K}^{\alpha u} \Delta \mathbf{u} + \mathbf{K}^{\alpha\alpha} \Delta \alpha = -\mathbf{R}^\alpha$$

(Simo, Armero & Taylor, 1993)

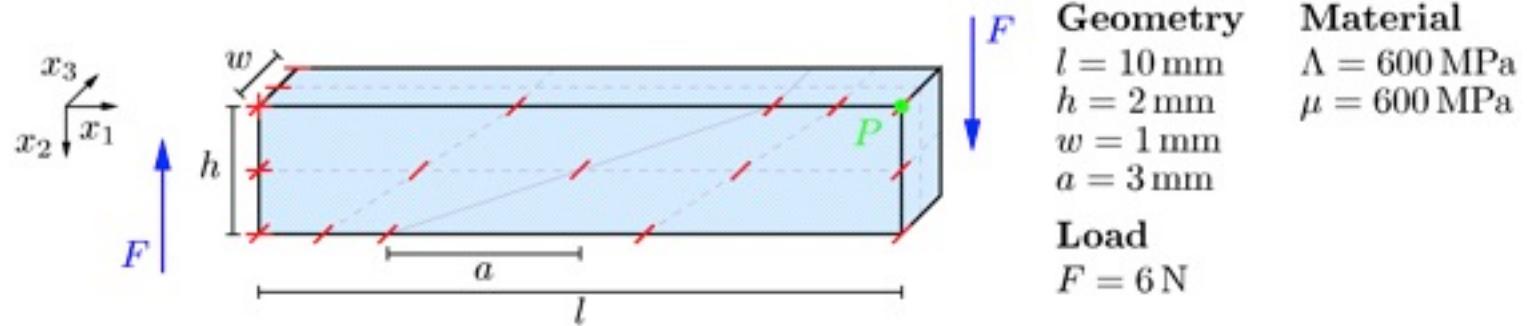


Numerical Example: Objectivity Test

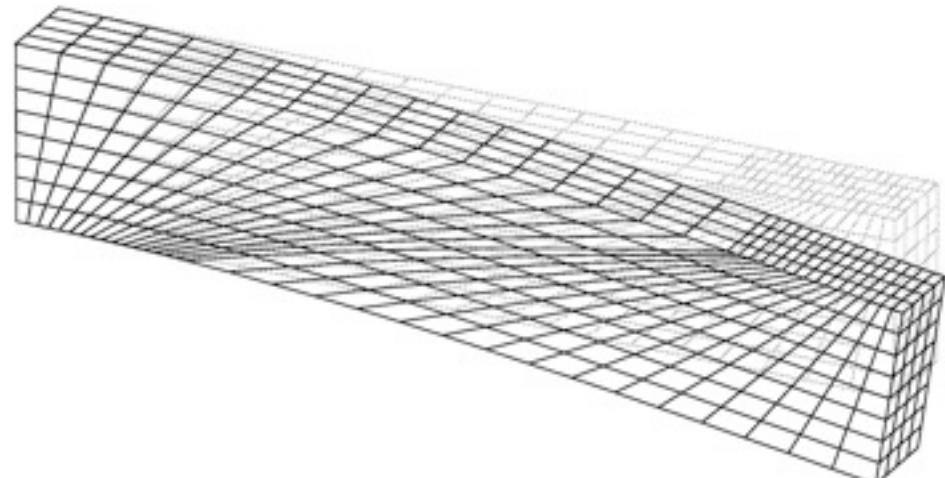


Numerical Example: Beam (1)

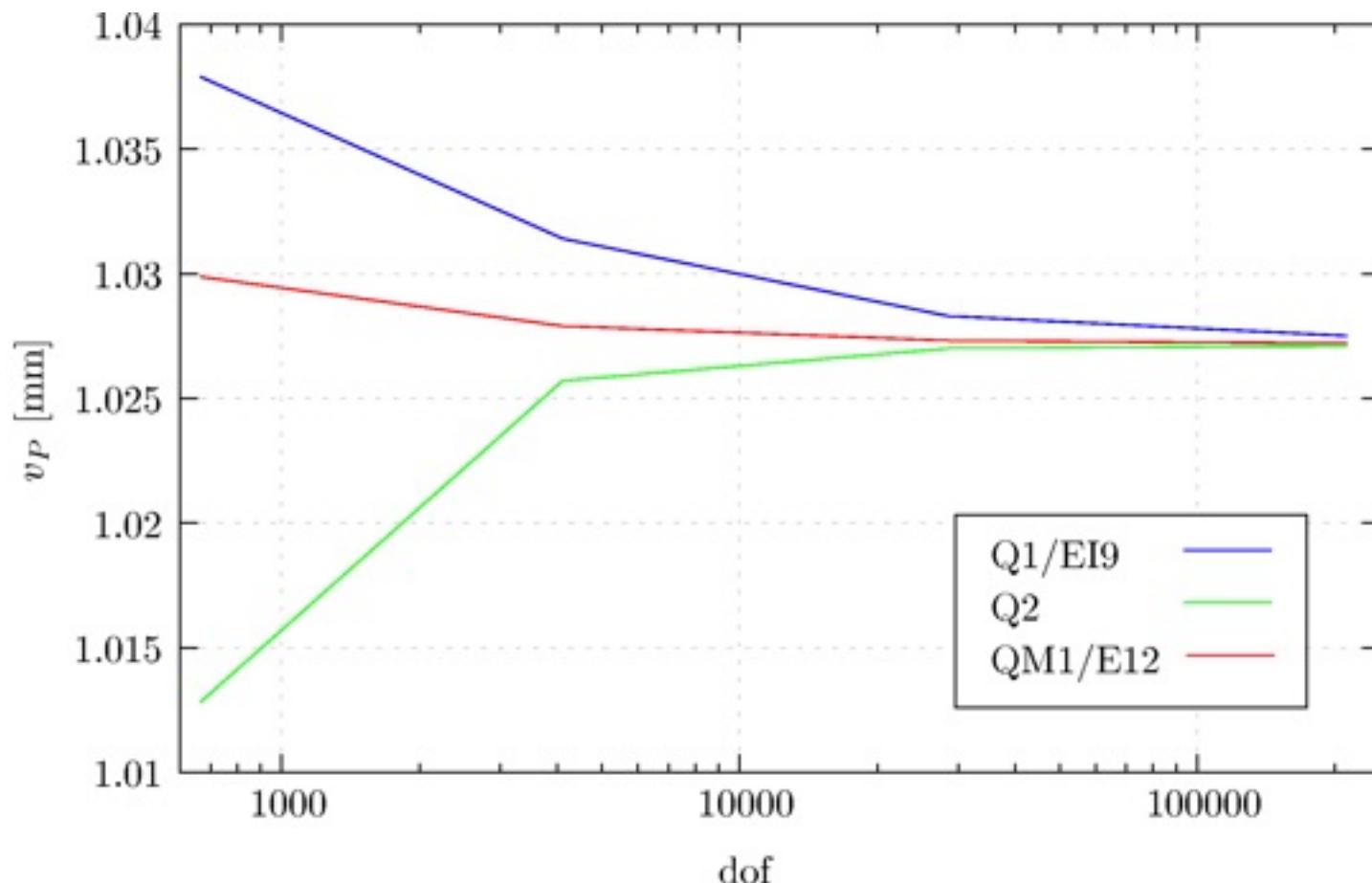
System



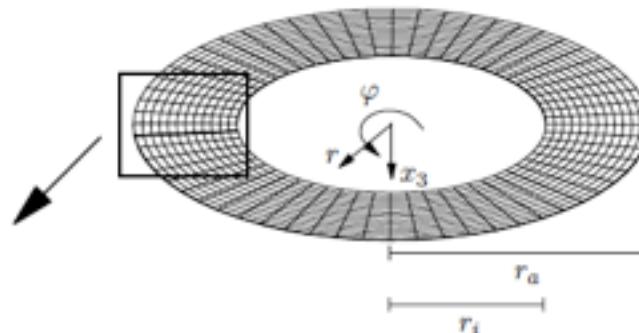
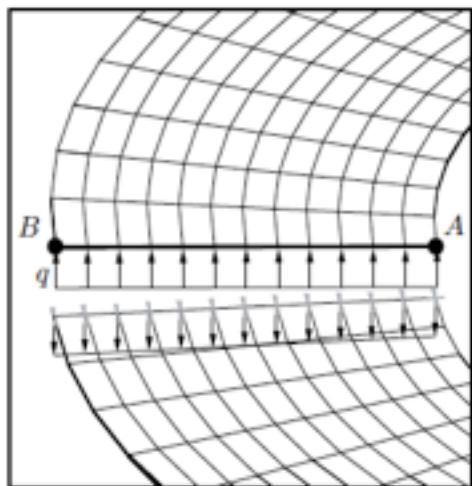
Deformed System



Numerical Example: Beam (2)



Numerical Example: Thin Plate Ring

**Geometry**

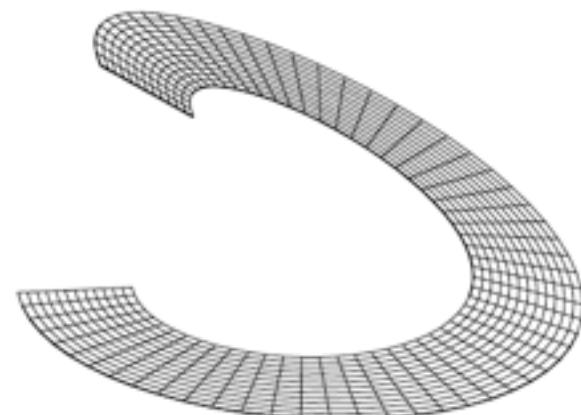
$r_a = 6 \text{ mm}$
 $r_i = 10 \text{ mm}$
 $h = 0.03 \text{ mm}$

Material

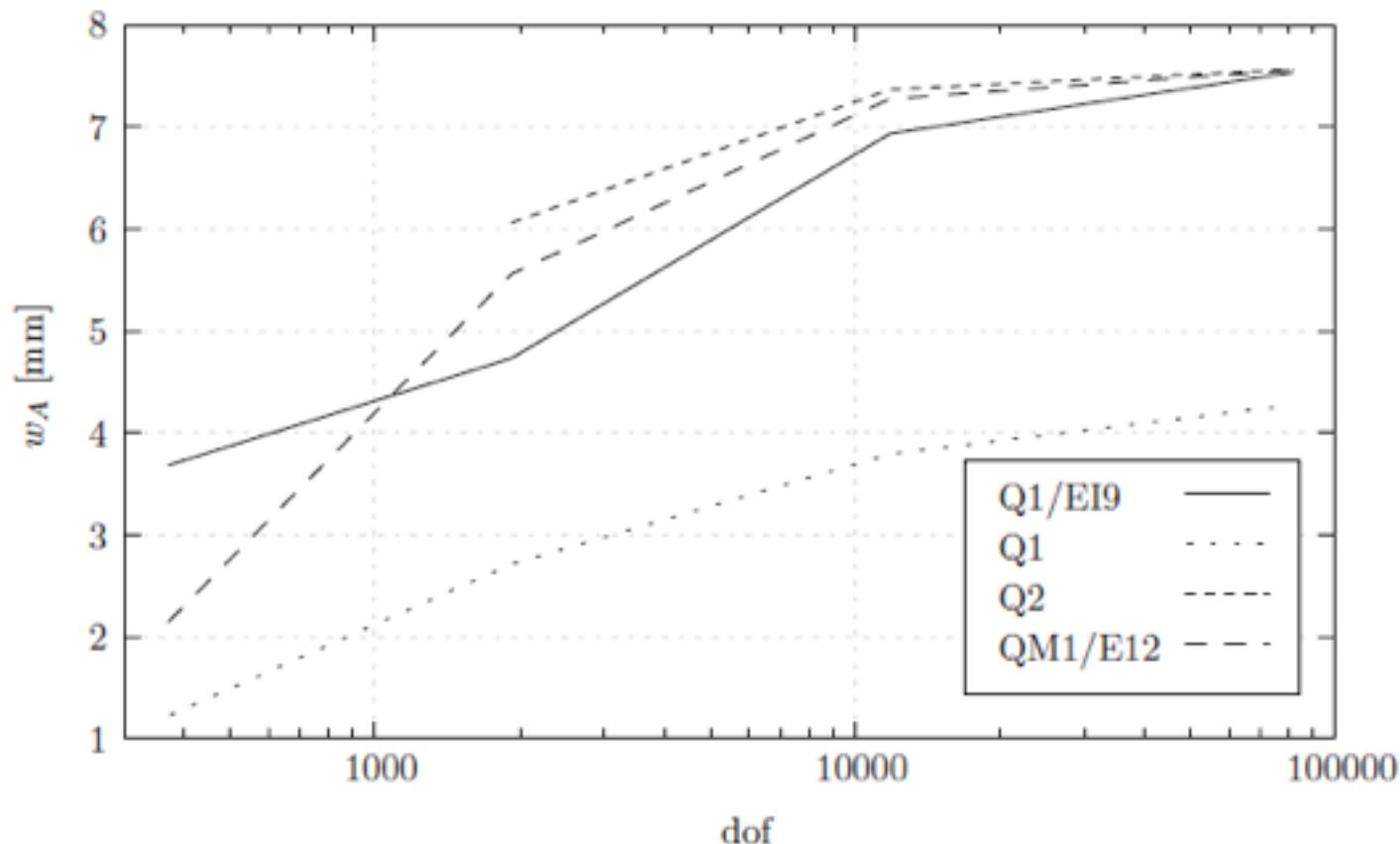
$K = 7000000 \text{ MPa}$
 $\mu = 10500000 \text{ MPa}$
 $\beta = -2$

Load

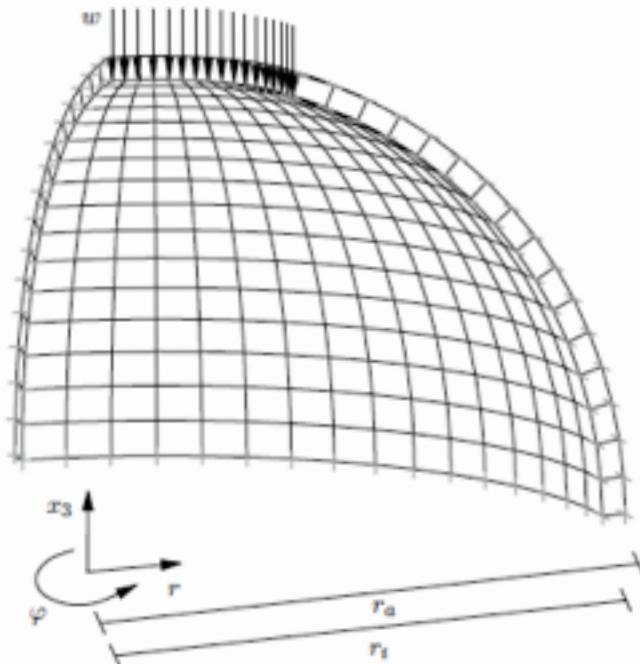
$q = 6.67 \text{ MPa}$



Numerical Example: Thin Plate Ring (2)

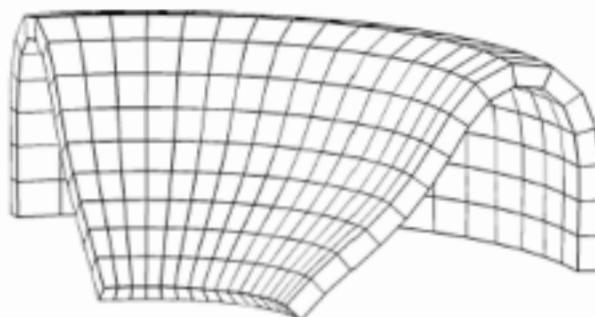


Numerical Example: Spherical Shell

**Geometry** $r_o = 9.75 \text{ mm}$
 $r_i = 10.25 \text{ mm}$ **Load** $w = 10 \text{ mm}$ **Material** $K = 175000 \text{ MPa}$
 $\mu = 80769.23 \text{ MPa}$
 $\beta = -2$

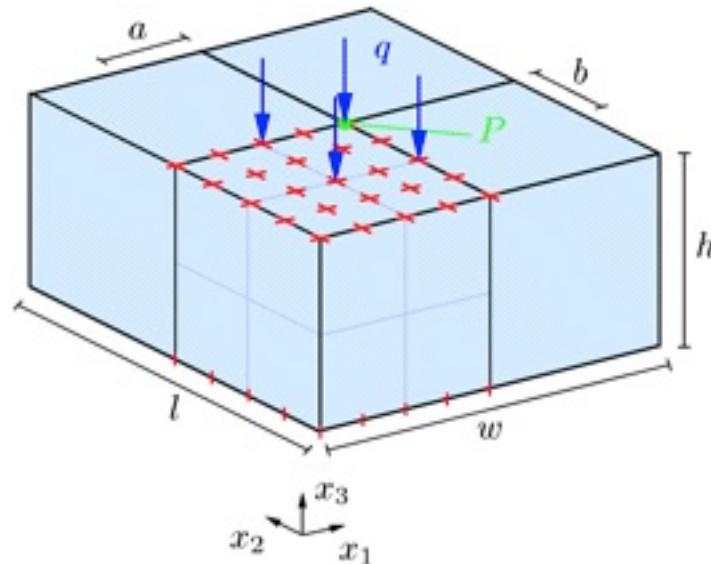
Degrees of freedom	Q1/EI9	Q1	Q2	QM1/EI12
1615	5767.66	9319.42	5485.56	5529.04
9471	5479.85	6508.51	5350.54	5379.84
62335	5377.19	5646.24	5348.75	
445695	5347.95	5416.33	5338.94	5340.71

Figure 13. Spherical shell: System, load and material data



Numerical Example: Block (1)

System

**Geometry**

$h = 50 \text{ mm}$
 $w = 100 \text{ mm}$
 $l = 100 \text{ mm}$
 $a = 25 \text{ mm}$
 $b = 25 \text{ mm}$

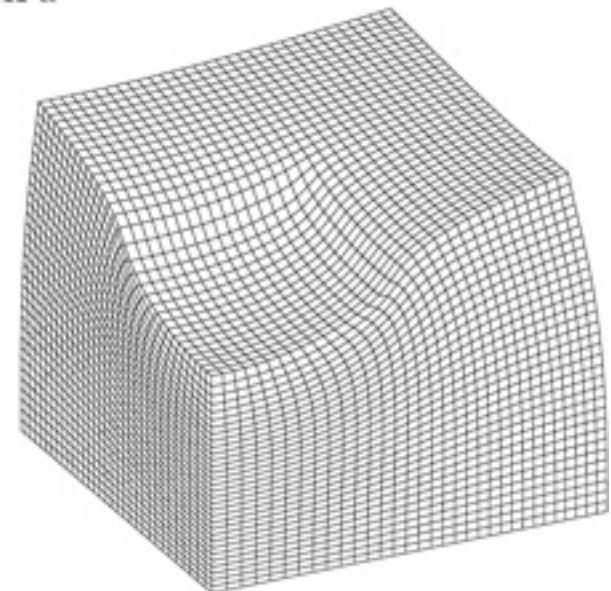
Load

$q = 3 \text{ MPa}$

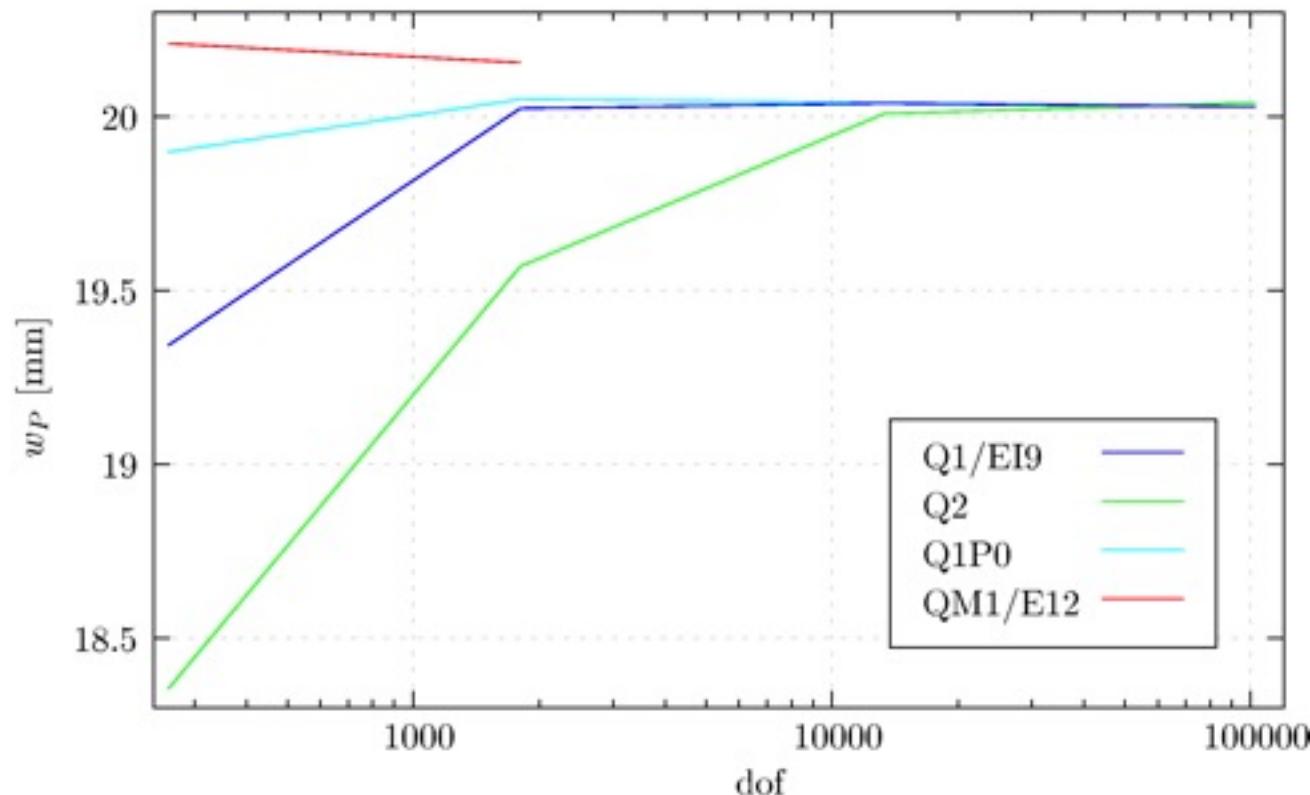
Material

$\Lambda = 499.92568 \text{ MPa}$
 $\mu = 1.61148 \text{ MPa}$

Deformed System



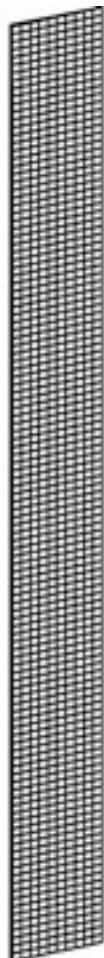
Numerical Example: Block (2)



Degrees of freedom	Q1/EI9	Q1/EI12	Q2	Q1P0	QM1/EI12
1800	4	4	4	4	5
102432	6	6	6	6	7

Numerical Example: Surface Buckling

(a)

**Geometry**

$$\alpha = 2^\circ$$

$$r_i = 0.01 \text{ mm}$$

$$r_a = 1 \text{ mm}$$

$$h = 10 \text{ mm}$$

Material

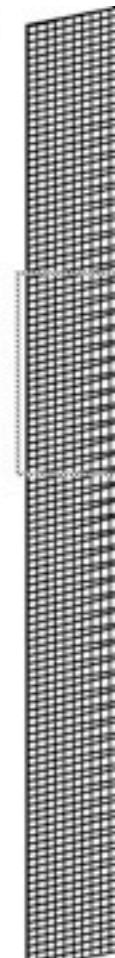
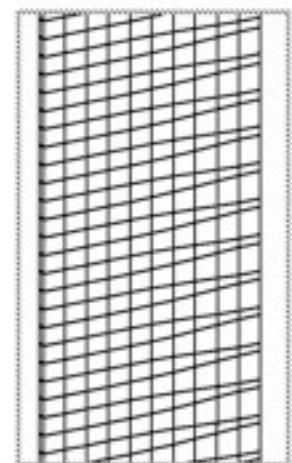
$$\Lambda = 216.04938 \text{ MPa}$$

$$\mu = 92.59259 \text{ MPa}$$

(b)



(c)

**Q1****Q2****Q1P0****Q1/EI9 (b)****QM1/E12 (c)**



Conclusion and Outlook

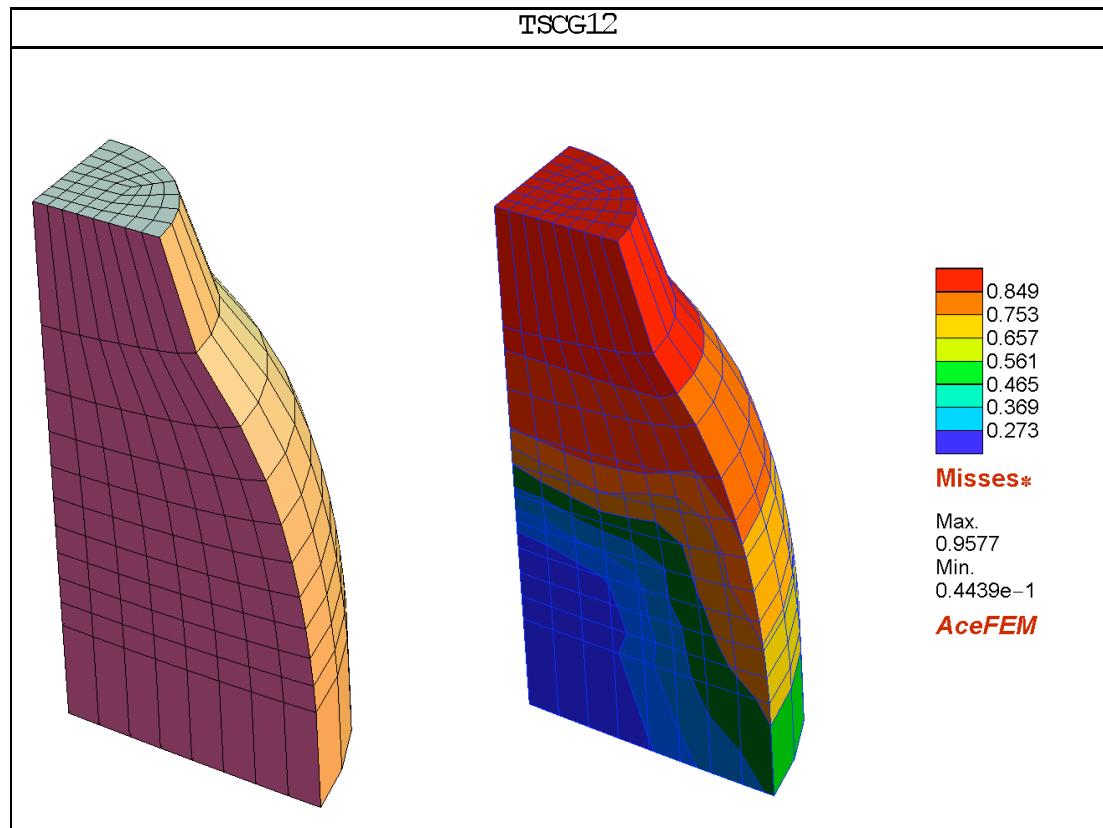
Advantages of Q1/EI9

- robust, no hourgassing
- no adjustment of any parameters
- works for coarse and unaligned meshes
- handles incompressible material
- has good bending properties

Outlook

- Implementation of inelastic constitutive models
(more complex than for standard elements)
- Solution by TSCG12 element

TSCG12 Element for plasticity



using 12 enhanced modes and an expansion of shape function derivatives using Taylor series

Korelc, Wriggers, Soric (2010)