Outline	Introd.	Obj.	BVP	Primal Prob.	Multi-Field Princ.	Dual Prob.	Dual FEF	Dual Anal.	Num. Appl. Clos.
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Duality in the Geometrically Exact Analysis of Three-Dimensional Framed Structures

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- Closure





Displacement-based finite element formulations

- are on the basis of most of the finite element models used in computer analysis of structures;
- assume the configuration variables as primary unknowns;
- lead to approximate *kinematically admissible solutions* in which *stress discontinuities may occur across element boundaries*; stress 'averaging' procedures are required;
- are well developed for both linear and nonlinear analyses.





Equilibrium-based finite element formulations

- are less common than the displacement-based finite element formulations;
- lead to approximate *statically admissible solutions*;
- have a special appeal for practical design engineers due to the exact transmission of stresses across interelement boundaries, thus avoiding the need for 'averaging' procedures;
- are not well studied in the context of the geometrically nonlinear analysis of framed structures.





Objectives and Scope

To present, in the framework of the quasi-static linear elastic analysis of geometrically exact framed structures modeled using the three-dimensional Reissner-Simo beam theory:

- Two dual energy-based variational formulations: one (Primal) derived from the well known Principle of Stationary Total Potential Energy, and the other (Dual) resulting from the Principle of Stationary Total Complementary Energy;
- An equilibrium-based (hybrid-mixed) finite element formulation relying on a modified Principle of Complementary Energy;
- A duality based method in which both primal and dual variational problems are studied in conjunction.





Kinematical Considerations

- The deformed geometry of a beam is described by the centroidal axis and the set of orientations of cross-sections;
- Only initially straight beam configurations and initially undistorted cross-sections are assumed;
- The geometric shape of the cross-sections is assumed to be arbitrary and constant along the beam;
- The cross-sections are assumed to suffer only rigid body motions during deformation;
- The beam theory is valid for arbitrarily large displacements and rotations **Geometrically Exact Beam Theory**.





Kinematical Considerations (Cont.)

- The deformed configuration of a beam is described by the position of the line of centroids of the cross-sections and also the rotations of the cross-sections;
- The rotations of the cross-sections are described using the Euler-Rodrigues formula, which is assumed to be parameterized through the total rotation vector as follows:

$$\mathbf{Q} = \mathbf{I} + rac{\sin heta}{ heta} \mathbf{\Theta} + rac{1 - \cos heta}{ heta^2} \mathbf{\Theta}^2$$

where $\boldsymbol{\Theta} = \textit{Skew}(\boldsymbol{ heta})$ and $heta = \| \boldsymbol{ heta} \|$





Reissner-Simo Beam BVP (Material Form)

• Differential Equations

Equilibrium $\{ \mathcal{T}_{e}^{r}(\mathbf{d})\sigma^{r} + \mathbf{q} = \mathbf{0}, \text{ in } \Omega \}$ Elasticity $\{ \sigma^{r} = \frac{\partial W(\varepsilon^{r}(\mathbf{d}))}{\partial \varepsilon^{r}}, \text{ in } \Omega \}$ Compatibility $\{ \varepsilon^{r} = \varepsilon^{r}(\mathbf{d}), \text{ in } \Omega \}$

- Neumann Boundary Conditions: $nH\sigma^r = \bar{q}$, on Γ_N
- Dirichlet Boundary Conditions: $\mathbf{d} = \bar{\mathbf{d}}$, on Γ_D

Remark: If the strain energy $W(\varepsilon^r)$ is differentiable and convex, by means of the Legendre transformation, the constitutive relations can be alternatively established using the format

$$arepsilon^r = rac{\partial W_c(oldsymbol{\sigma}^r)}{\partial oldsymbol{\sigma}^r}, ext{ in } \Omega$$



Principle of Stationary Total Potential Energy

• Let \mathcal{U}_k and \mathcal{V}_k be the kinematically admissible function spaces

$$\mathcal{U}_k = \{ \mathbf{d} \in \mathcal{H}^1(\Omega) | \ \mathbf{d} = \bar{\mathbf{d}} \text{ on } \Gamma_D \}$$
$$\mathcal{V}_k = \{ \delta \mathbf{d} \in \mathcal{H}^1(\Omega) | \ \delta \mathbf{d} = \mathbf{0} \text{ on } \Gamma_D \}$$

The total potential energy associated with vector **d** is the one-field functional Π_p(**d**) : U_k(Ω) → R given by

$$\Pi_{\rho}(\mathbf{d}) = \int_{\Omega} [W(\varepsilon^{r}(\mathbf{d})) - \mathbf{q} \cdot \mathbf{d}] \ dS - [\mathbf{\bar{q}} \cdot \mathbf{d}]_{\Gamma_{N}}$$

 Principle of Stationary Total Potential Energy (PSTPE): vector d ∈ U_k is a solution of the BVP iff δΠ_p = 0 ∀δd ∈ V_k, i.e., a beam is in equilibrium iff its total potential energy takes a stationary value for all kinematically admissible displacement fields.





Hybrid-Multi-Field Variational Principles

- The PSTPE can be generalized by means of the Lagrangian multiplier method leading to a **Generalized Variational Principle (GVP)**;
- The GVP can afterwards be particularized into different **Hybrid-Multi-Field Principles**, *e.g.*:
 - Principles of Hu-Washizu;
 - Principles of Hellinger-Reissner;
 - Principle of Total Complementary Energy, etc.





Principle of Stationary Total Complementary Energy

 $\bullet~$ Let \mathcal{U}_s and \mathcal{V}_s be the statically admissible function spaces

$$\mathcal{U}_{s} = \{(\boldsymbol{\sigma}^{r}, \mathbf{d}) \in (\mathcal{H}^{1}(\Omega) \times \mathcal{H}^{1}(\Omega)) | \mathcal{T}_{e}^{r}(\mathbf{d})\boldsymbol{\sigma}^{r} + \mathbf{q} = \mathbf{0} \text{ in } \Omega \text{ and} \\ n\mathbf{H}\boldsymbol{\sigma}^{r} - \bar{\mathbf{q}} = \mathbf{0} \text{ on } \Gamma_{N}\}$$

$$\mathcal{V}_{s} = \{ (\delta \boldsymbol{\sigma}^{r}, \mathbf{d}) \in \mathcal{H}^{1}(\Omega) \times \mathcal{H}^{1}(\Omega) | \mathcal{T}_{e}^{r}(\mathbf{d}) \delta \boldsymbol{\sigma}^{r} = \mathbf{0} \text{ in } \Omega \text{ and} \\ n \mathbf{H} \delta \boldsymbol{\sigma}^{r} = \mathbf{0} \text{ on } \Gamma_{N} \}$$

 The complementary energy associated with (σ^r, d) is the 2-field functional Π_c : U_s(Ω) → R given by

$$\Pi_{c}(\boldsymbol{\sigma}^{r},\mathbf{d}) = \int_{0}^{L} [W_{c}(\boldsymbol{\sigma}^{r}) - \boldsymbol{\sigma}^{r} \cdot \boldsymbol{\varepsilon}^{r}(\mathbf{d}) + \boldsymbol{\sigma}^{r} \cdot \mathcal{T}_{c}^{r}(\mathbf{d})\mathbf{d}] dS - [n\mathbf{H}\boldsymbol{\sigma}^{r} \cdot \bar{\mathbf{d}}]_{\Gamma_{D}}$$

• Principle of Stationary Total Complementary Energy (PSTCE): the pair $(\sigma^r, \mathbf{d}) \in \mathcal{U}_s$ is a solution of the BVP iff $\delta \Pi_c = 0 \ \forall (\delta \sigma^r, \mathbf{d}) \in \mathcal{V}_s$.





Hybrid-Mixed Complementary Energy

• If the equilibrium equations are assumed to be relaxed within the framework of the PSTCE, the following hybrid-mixed complementary energy $\Pi_c^g : \chi(\Omega) \to \mathcal{R}$ can be obtained

$$\Pi_{c}^{g}(\boldsymbol{\sigma}^{r}, \mathbf{d}, \mathbf{d}^{\Gamma}) = \sum_{b=1}^{B} \int_{\Omega_{b}} [W_{c}(\boldsymbol{\sigma}_{b}^{r}) - \boldsymbol{\sigma}_{b}^{r} \cdot \boldsymbol{\varepsilon}_{b}^{r}(\mathbf{d}_{b}) + \mathbf{q}_{b} \cdot \mathbf{d}_{b}] \ d\Omega_{b}$$
$$+ [\mathbf{\bar{q}} \cdot \mathbf{d}^{\Gamma}]_{\Gamma_{N} \cup \Gamma_{int}} + [n\mathbf{H}\boldsymbol{\sigma}^{r} \cdot (\mathbf{d} - \mathbf{J}_{N}\mathbf{d}^{\Gamma})]_{\Gamma_{N} \cup \Gamma_{int}} + [n\mathbf{H}\boldsymbol{\sigma}^{r} \cdot (\mathbf{d} - \mathbf{J}_{D}\mathbf{\bar{d}})]_{\Gamma_{D}}$$

- J_N and J_D represent transformation matrices mapping global vectors (matrices) onto local element vectors (matrices) defined on $\Gamma_N \cup \Gamma_{int}$ and Γ_D , respectively;
- The functions in class $\chi(\Omega)$ consist of pairs $(\sigma_b^r, \mathbf{d}_b) \in \mathcal{H}^0(\Omega_b) \times \mathcal{H}^1(\Omega_b)$, with $1 \le b \le B$, and a real-valued vector \mathbf{d}^{Γ} defined on $\Gamma_N \cup \Gamma_{int}$.





Hybrid-Mixed Complementary Variational Principle

• The variational (weak) problem $\delta \Pi_c^g = 0, \ \forall (\delta \sigma^r, \delta \mathbf{d}, \delta \mathbf{d}^{\Gamma}) \in \chi(\Omega)$ is formally equivalent to the following system of Euler-Lagrange equations

$$\mathcal{T}_{e}^{r}(\mathbf{d}_{b})\boldsymbol{\sigma}_{b}^{r}+\mathbf{q}_{b}=\mathbf{0} \text{ in } \Omega_{b}$$

$$\boldsymbol{\varepsilon}_{b}^{r}(\boldsymbol{\sigma}_{b}^{r})-\boldsymbol{\varepsilon}_{b}^{r}(\mathbf{d}_{b})=\mathbf{0} \text{ in } \Omega_{b}$$

$$\bar{\mathbf{q}}-n\mathbf{J}_{N}^{T}\mathbf{H}\boldsymbol{\sigma}^{r}=\mathbf{0} \text{ in } \Gamma_{N}\cup\Gamma_{int}$$

$$\mathbf{d}-\mathbf{J}_{N}\mathbf{d}^{\Gamma}=\mathbf{0} \text{ in } \Gamma_{N}\cup\Gamma_{int}$$

$$\mathbf{d}-\mathbf{J}_{D}\bar{\mathbf{d}}=\mathbf{0} \text{ on } \Gamma_{D}$$

with $1 \leq b \leq B$.





• Element variables:

$$\sigma^{r^h} = \begin{bmatrix} \mathbf{n}^r \\ \mathbf{m}^r_i + (\mathbf{m}^r_j - \mathbf{m}^r_i)\frac{s}{L} \end{bmatrix}, \quad \mathbf{d}^h = \begin{bmatrix} \mathbf{u}_i + (\mathbf{u}_j - \mathbf{u}_i)\frac{s}{L} \\ \boldsymbol{\theta} \end{bmatrix}$$

• Nodal variables: \mathbf{d}^{Γ} (generalized displacements)

Remarks:

- As the approximate displacements are one degree greater than the approximate rotations, this formulation is capable of representing zero shear solutions and is, thus, completely free from shear locking;
- Using these approximations, the formulation can provide solutions that satisfy the equilibrium differential equations in strong form, as well as the stress continuity conditions (when assuming zero distributed loads);
- Furthermore, the necessary and sufficient condition for solvability of the discrete linearized system of equations is fulfilled either for a single element or a patch of elements with appropriate boundary conditions $(n_{\sigma r} \ge n_d n_r)$.





Linearized Global System of Equations

• The linearized global system of equations can be stated as

$$\begin{split} \mathbf{r}(\mathbf{p}) + \mathbf{T}(\mathbf{p})\Delta\mathbf{p} &= \mathbf{0} \ , \\ \mathbf{p} &= \left[\begin{array}{c} \mathbf{p}_{\sigma^{r}} \\ \mathbf{p}_{\mathbf{d}} \end{array} \right], \ \mathbf{T} &= \left[\begin{array}{c} \mathbf{F} & \mathbf{A}^{T} \\ \mathbf{A} & \mathbf{K}_{c} \end{array} \right] \\ \mathbf{K}_{eg} &= \mathbf{A}\mathbf{F}^{-1}\mathbf{A}^{T} - \mathbf{K}_{c} \end{split}$$

(for the classification of the stability of the equilibrium)

•
$$\mathbf{F} = \frac{\partial^2 \Pi_c^g}{\partial \mathbf{p}_{\sigma^r} \partial \mathbf{p}_{\sigma^r}}$$
 - flexibility matrix
• $\mathbf{A} = \frac{\partial^2 \Pi_c^g}{\partial \mathbf{p}_d \partial \mathbf{p}_{\sigma^r}}$ - equilibrium matrix; \mathbf{A}^T - compatibility matrix
• $\mathbf{K}_c = \frac{\partial^2 \Pi_c^g}{\partial \mathbf{p}_d \partial \mathbf{p}_d}$ - stiffness matrix



Fully Linear Case (FLC) vs Geom. Nonlinear Case (GNC)

FLC	GNC
$\Pi_{p}(\mathbf{d})$ is convex	$\Pi_{\rho}(\mathbf{d})$ is nonconvex
$\Pi_c(\sigma^r)$ is concave	$\Pi_c(\pmb{\sigma}^r, \mathbf{d})$ is a saddle functional
$\epsilon \geq (\epsilon_k, \epsilon_s)$	$\epsilon \stackrel{\geq}{<} ?(\epsilon_k, \epsilon_s)^*$

* Extremum conditions of Π_p and Π_c are required (Nobel and Sewell 1972, Gao and Strang 1989)

•
$$\epsilon = \left|\overline{\Pi}_{p} - \overline{\Pi}_{c}\right|, \ \epsilon_{k} = \left|\overline{\Pi}_{p} - \Pi_{p}\right|, \ \epsilon_{s} = \left|\Pi_{c} - \overline{\Pi}_{c}\right|$$

•
$$\overline{\Pi}_{p} = \inf_{\mathbf{d} \in \mathcal{U}_{k}} \Pi_{p}(\mathbf{d}), \ \overline{\Pi}_{c} = \sup_{\boldsymbol{\sigma}^{r} \in \mathcal{U}_{s}} \Pi_{c}(\boldsymbol{\sigma}^{r})$$











Deformed Configurations







Diagrams of moments for $P = 2.5 \times 10^{-5}$ (16FE)





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Cantileve	Cantilever beam subject to an end force												
Ene	Energies												













Deformed Configurations (5FE per leg)







Diagrams of moments for P = 10000 (5FE per leg)







Energies for P = 15000







Problem Definition







Deformed Configurations

Lateral Torsion Buckling





























Problem Definition















Energies for T = 210







- The present hybrid-mixed FE formulation, established within the framework of the geometrically exact (Reissner-Simo) analysis of 3D framed structures, is:
 - variationally consistent;
 - completely free from shear locking;
 - capable of producing statically admissible approximate solutions;
- The present duality based method opens a new way on a posteriori error estimation and on possible bounding aspects within the framework of geometrically nonlinear analysis of framed structures.





- Consider higher-order polynomial sets of approximate functions within the dual FE formulation (*p*-type refinement schemes);
- Consider initially curved beam elements within the framework of the dual formulation;
- Incorporate general cross-sectional in-plane changes and out-of-plane warping phenomena within the framework of the dual formulation;
- Include physical nonlinearities within the framework of the dual formulation;





- Extend the dual formulation to shells and membranes;
- Derive (hybrid-) mixed FE formulations from other (hybrid-) multi-field variational principles;
- Investigate, from a mathematical point of view, the numerical stability of the present dual FE formulation;
- Investigate alternative error estimation methods which can provide guaranteed upper bounds of the exact error of the approximate solutions (considering both global and local quantities of interest).



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