

National Academy of Sciences of Ukraine

Institute of Magnetism



Thin Films of Ni-Mn-Ga

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- L'vov V.A., Taras Shevchenko University, Kiev, Ukraine
- Ohtsuka M., IMRAM, Tohoku University, Sendai, Japan

PLAN

1

Introduction

2

Submicron Ni-Mn-Ga/substrate thin films ferromagnetic martensites:

- Fabrication by magnetron sputtering
- Transformation behavior
- Structure and substructure
- Magnetic properties

Generally, all properties are found to be dependent regularly on film thickness and substrate nature

3

Actuation ability
Current work

Development of thin film technologies of FSMA's

(1) Free-standing NiMnGa films

They are implemented and characterized as actuators:
optical microscanner

(2) *submicron* Ni-Mn-Ga/substrate thin film composites

Possible applications: magnetic MEMS, microsensors,
memory storage etc.

Fabrication methods used: molecular beam epitaxy, magnetron sputtering,
evaporation technique, laser ablation and electron beam deposition

Fabrication of Ni-Mn-Ga/substrate Thin Film Composites

RF magnetron sputtering:

- Target: Hot pressing of Ni, Mn and $\text{Mn}_{30}\text{Ga}_{70}$ mixed powders
→ nominal compositions $\text{Ni}_{49.5}\text{Mn}_{28.0}\text{Ga}_{22.5}$, $\text{Ni}_{52}\text{Mn}_{24}\text{Ga}_{24}$
- Substrates: Al_2O_3 ceramic, Mo foil (5 and 10 μm), $\text{MgO}(100)$,
[Si(100)+SiN_x 500nm] wafers

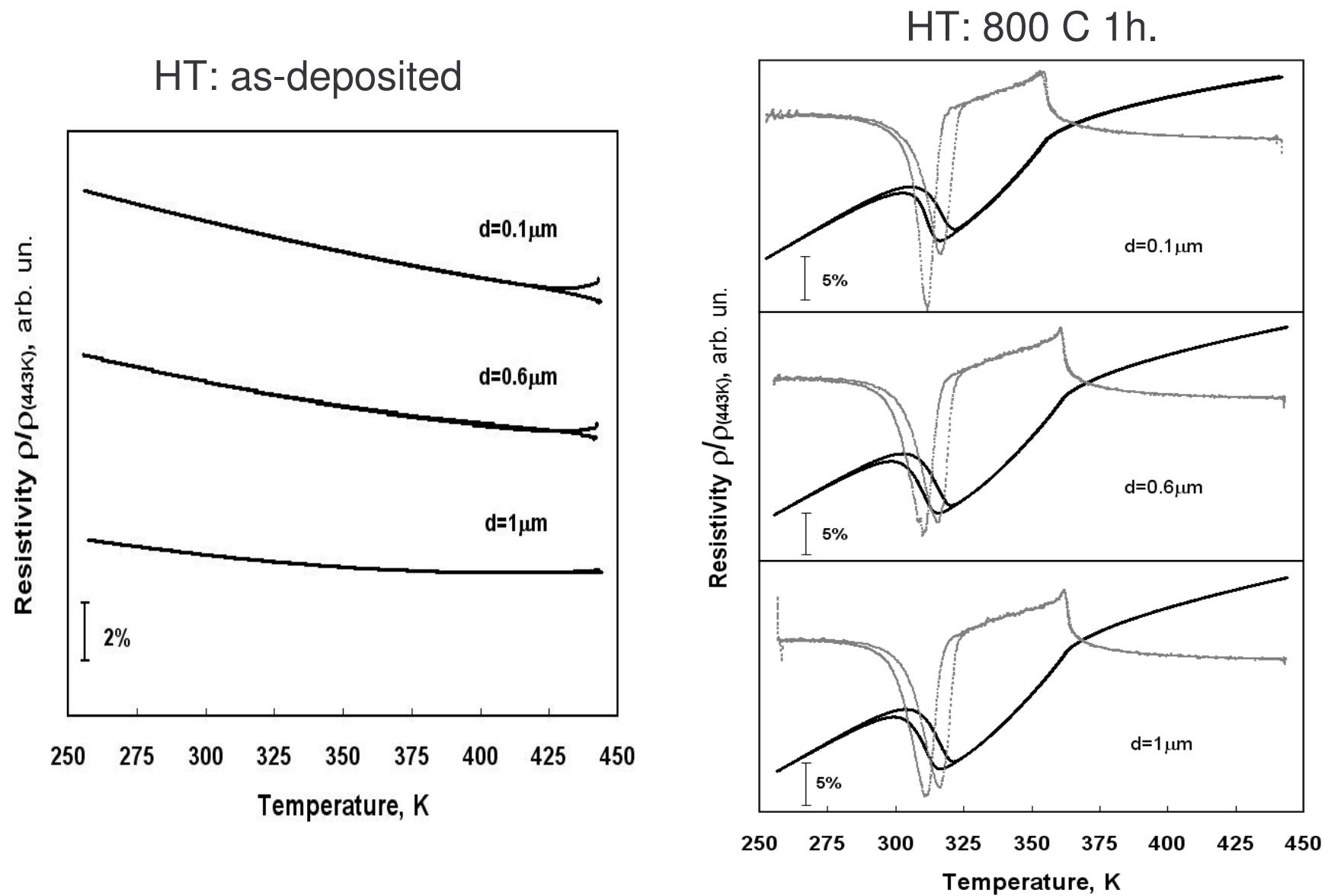
Sputtering conditions:

- RF power: 200 W
- Substrate temperature: 323 K
- Ar gas flow rate: 230 mm^3s^{-1}
- **Film thicknesses: 0.1 - 1 μm , 5 μm (reference)**

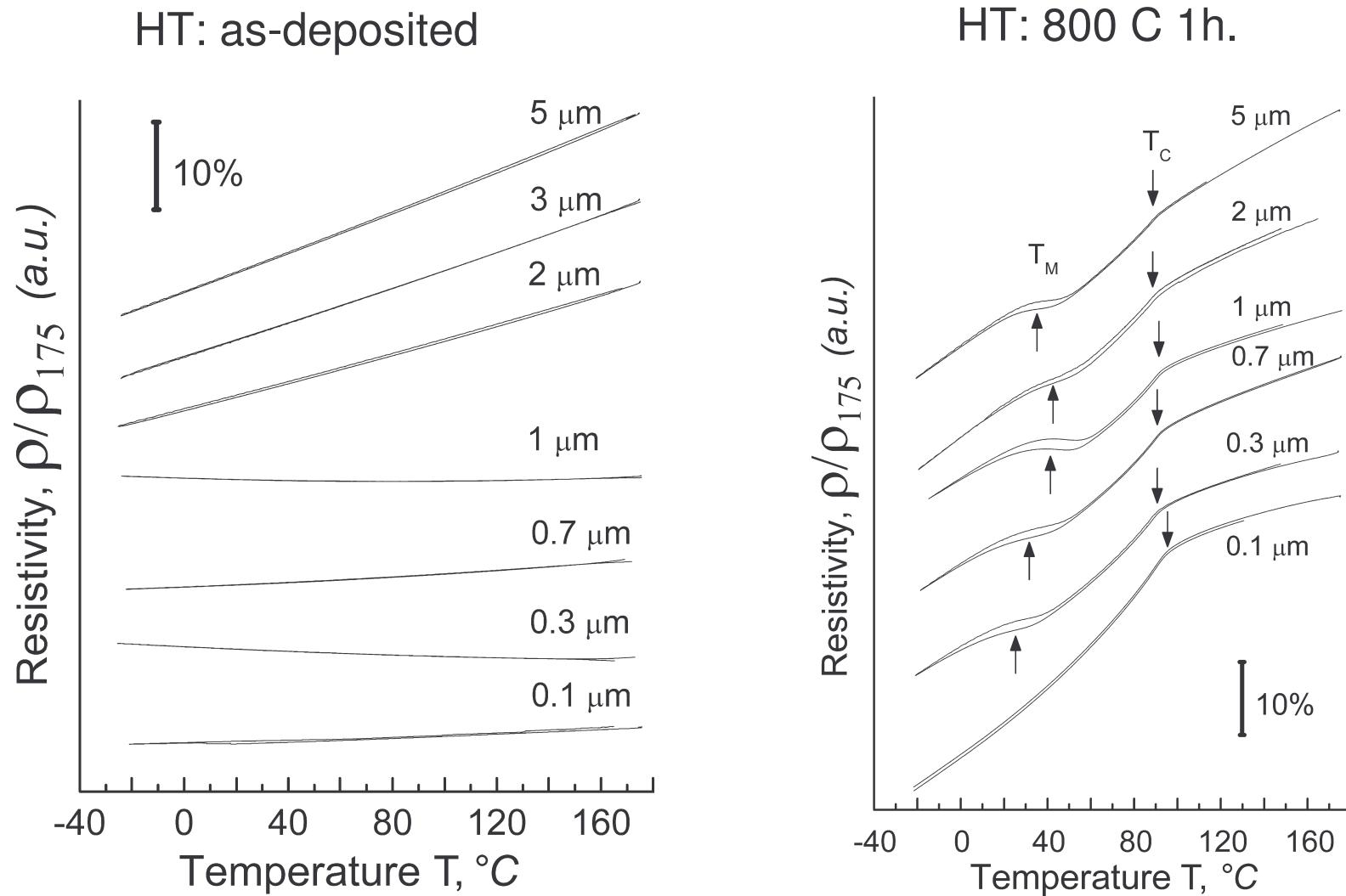
Heat treatment:

- Temperature: **1073 K**
- Time: **1h.**
- Pressure: $\sim 2 \cdot 10^{-4}$ Pa

Transformation behavior of Ni49/MgO(100)

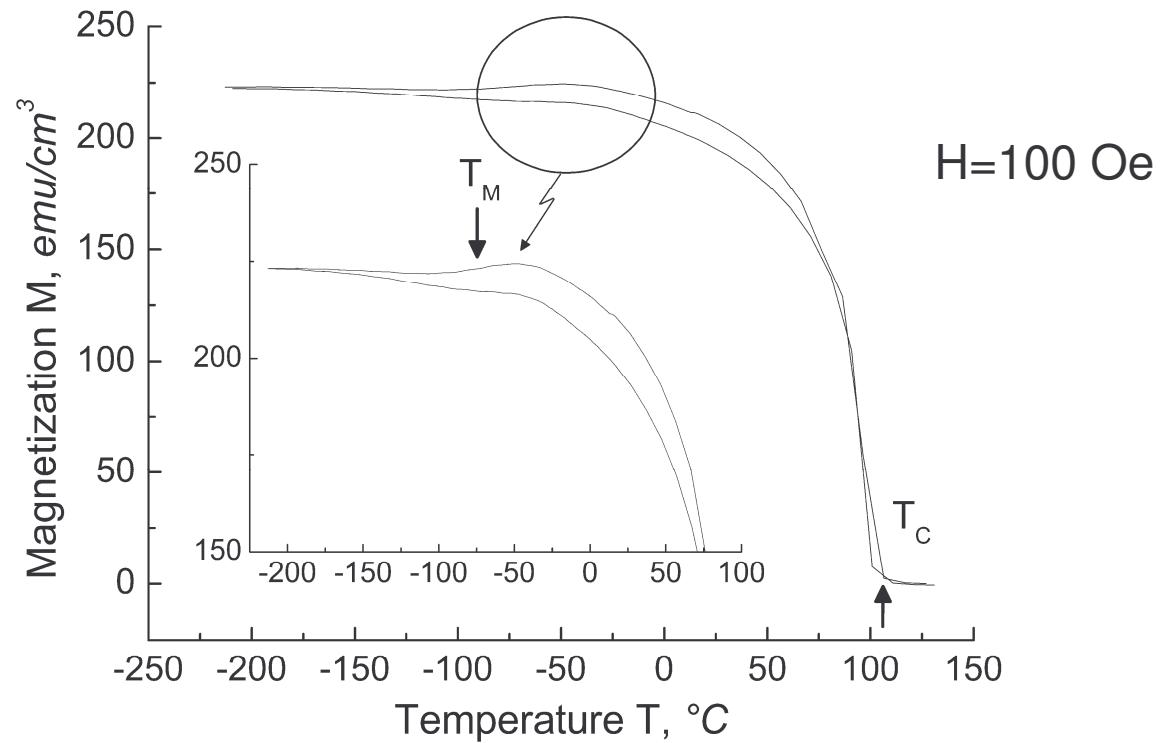


Transformation behavior of Ni49/Si(100)

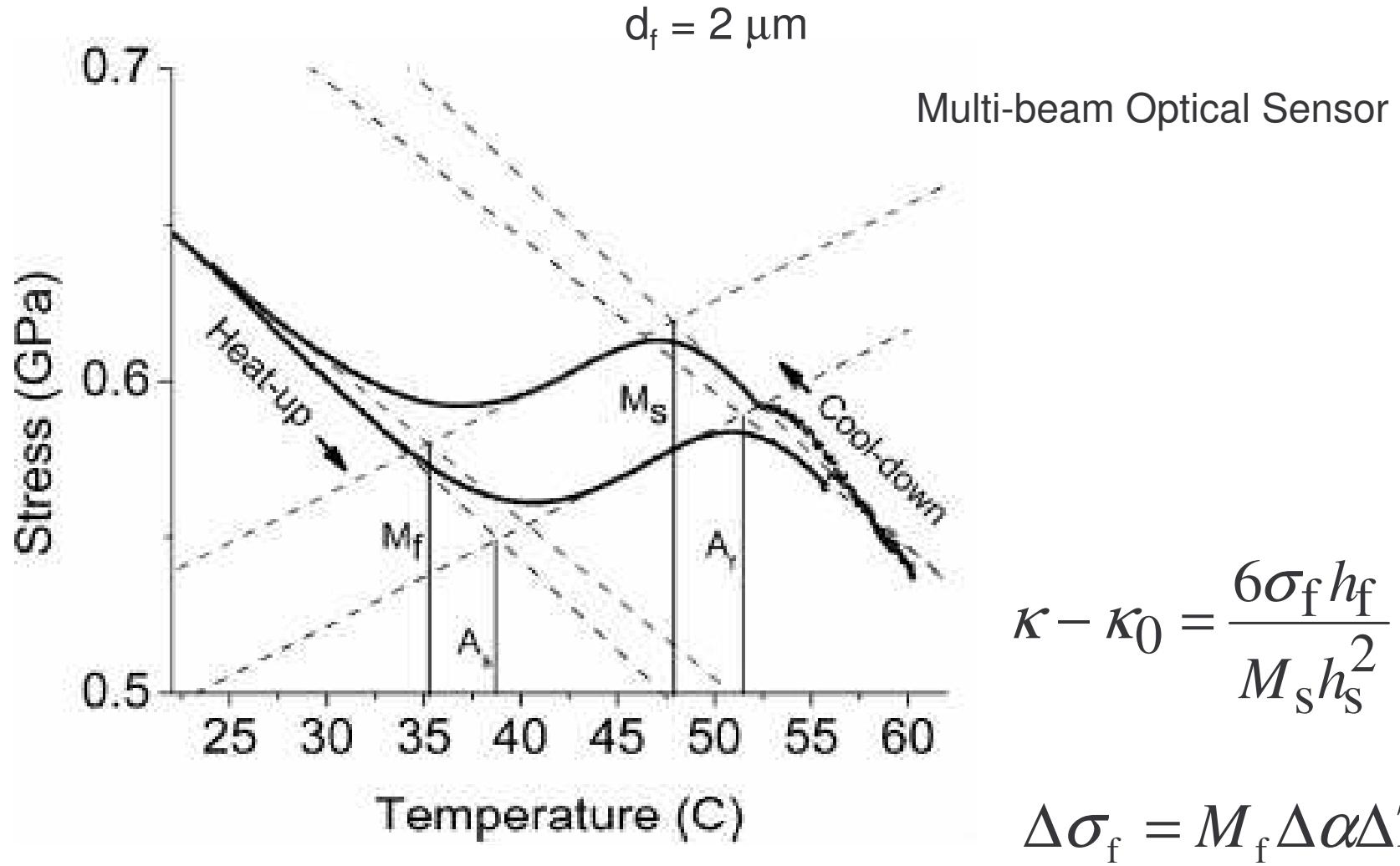


Transformation behavior of Ni49/Si(100)

$$d_f = 0.1 \text{ } \mu\text{m}$$



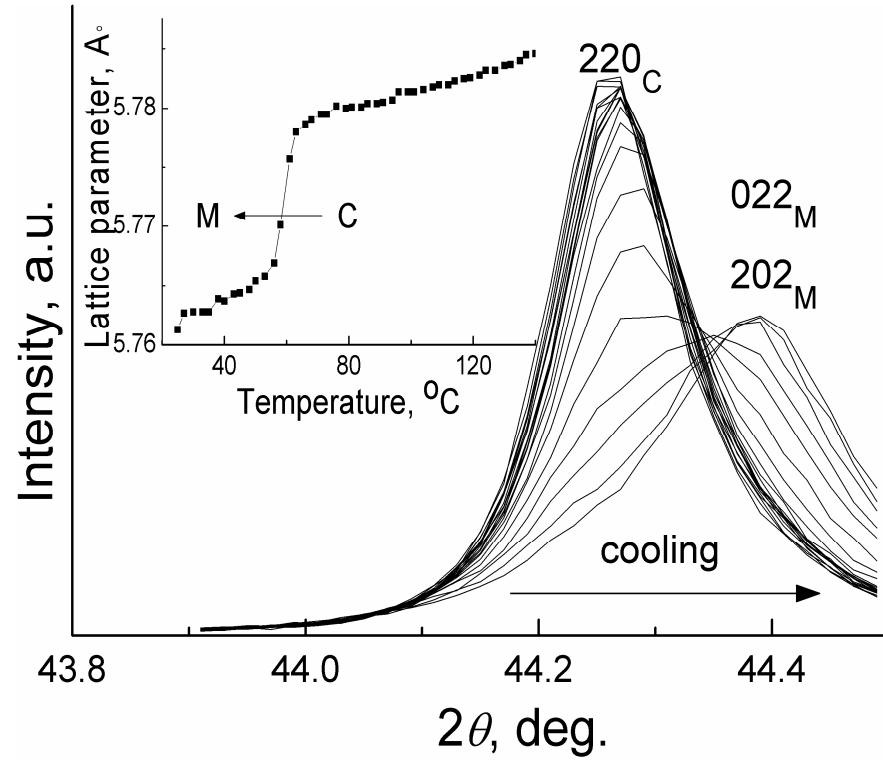
Substrate curvature analysis for Ni49/Si



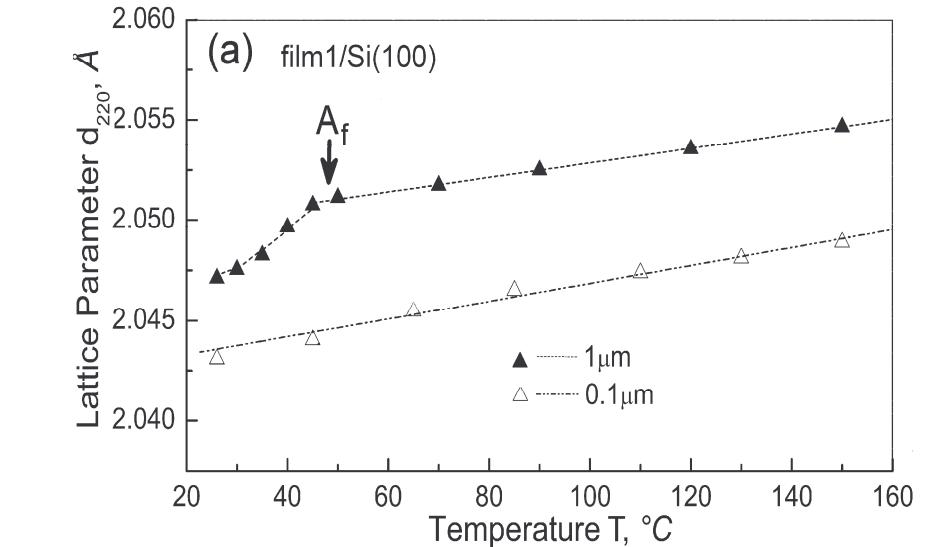
Transformation behavior in XRPD

Ni49/Si(100)

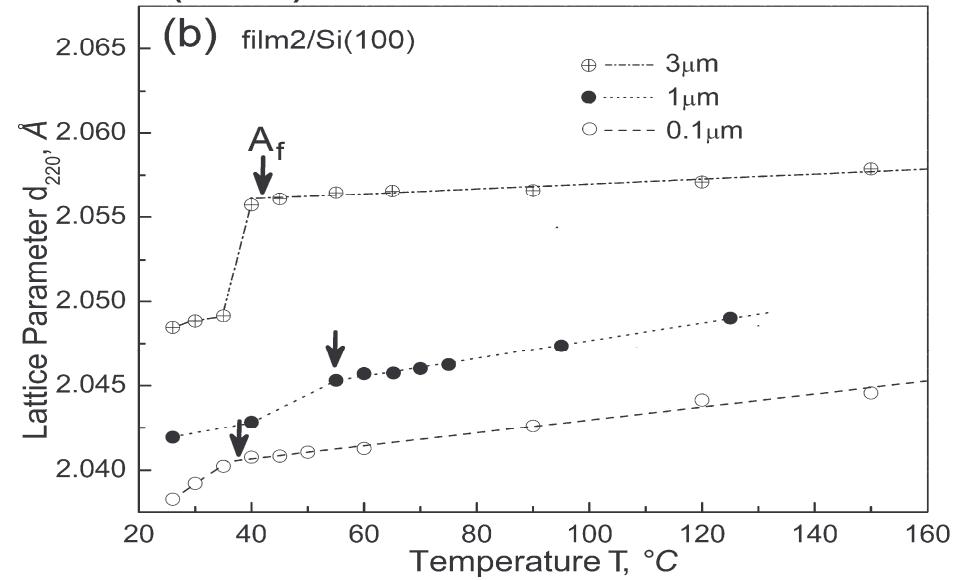
Ni52/Mo10, $d_f = 1 \mu\text{m}$



Ni52/Si(100)

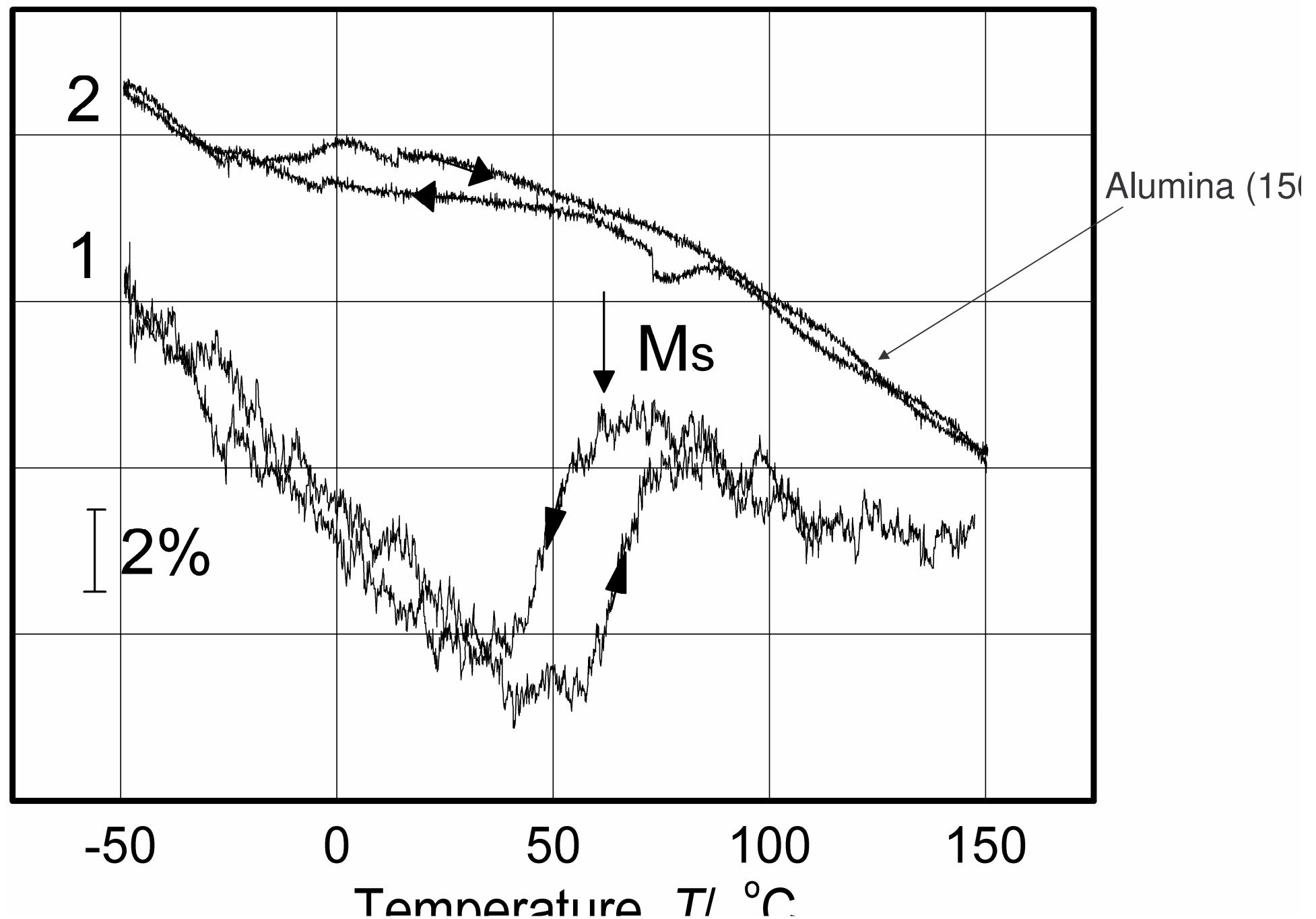


(b) film2/Si(100)



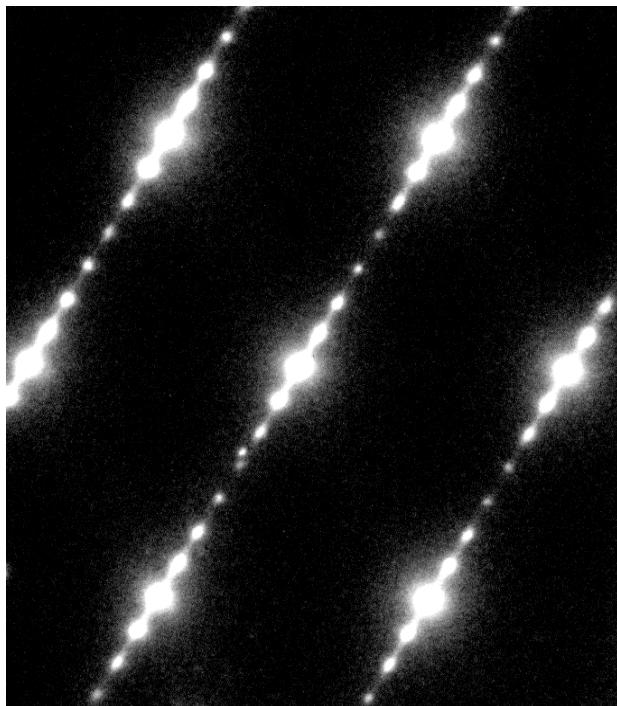
Transformation behavior studied by DMA

Ni49(0.6 μ m)/alumina(150 μ m)



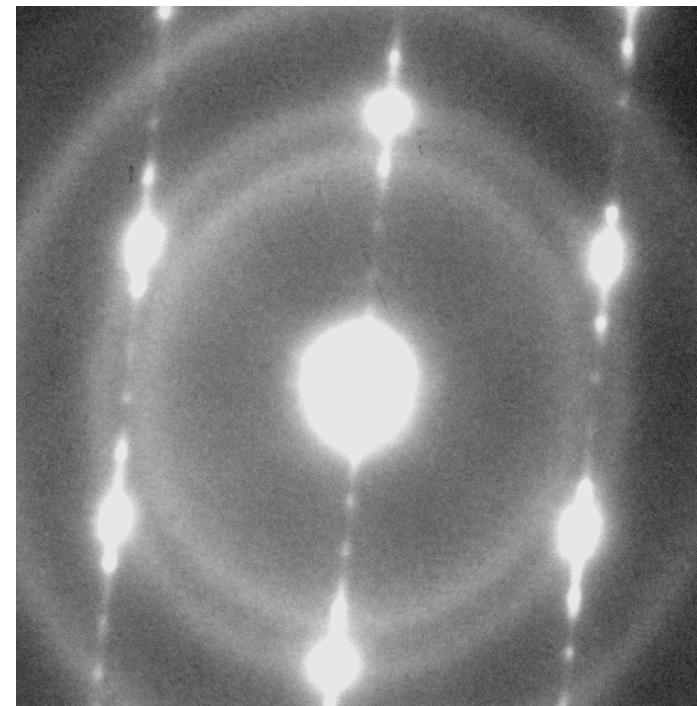
Structural Properties

Crystal structure by TEM of 5 μm free films



52 at.% Ni, $T_{\text{HT}}:1073 \text{ K}$ (3.6 ks)

14M-martensite

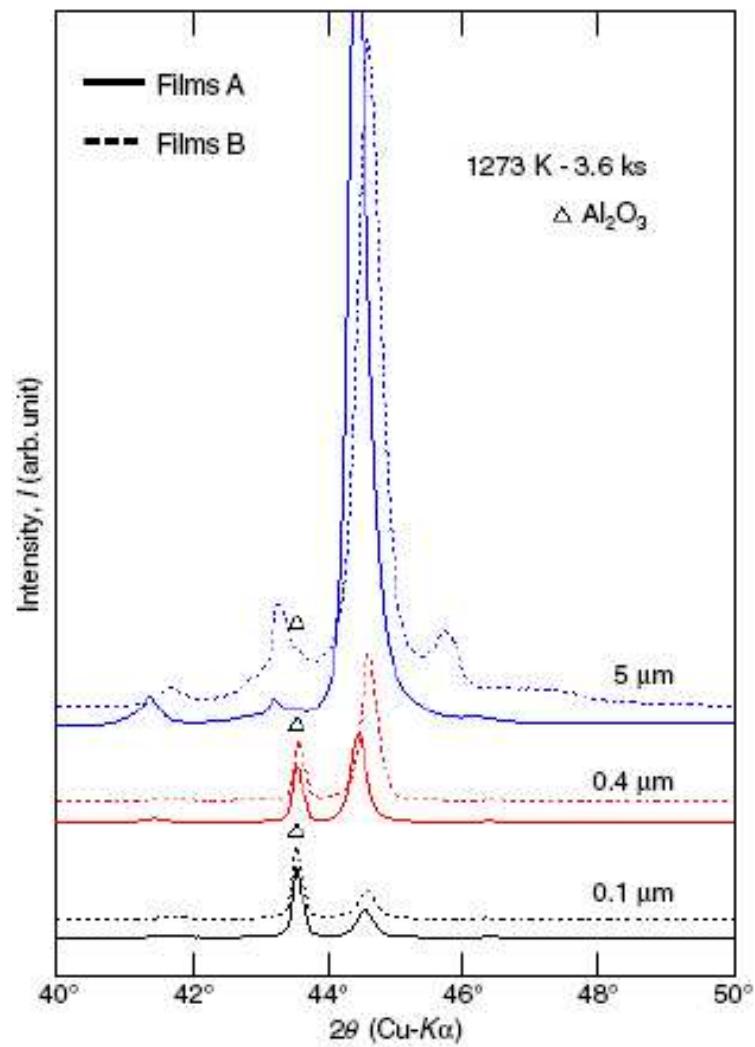


49.5 at.% Ni, $T_{\text{HT}}:1073 \text{ K}$ (3.6 ks)

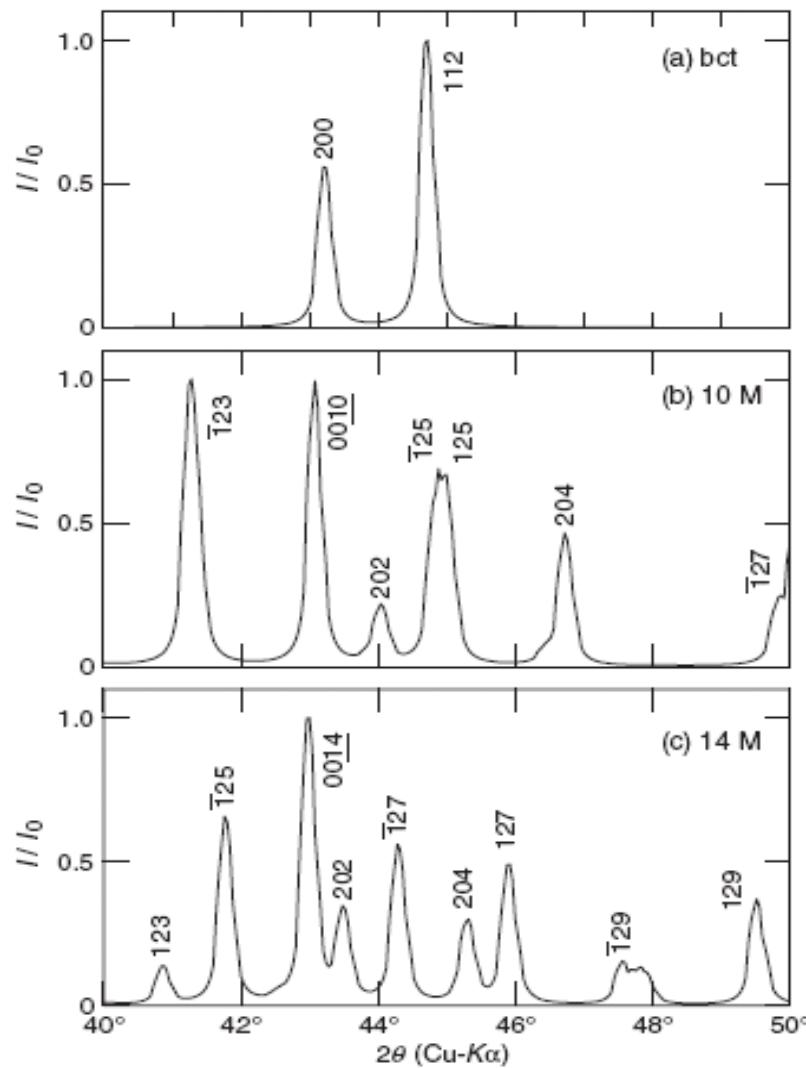
10M-martensite

XRPD

Ni49,52/alumina

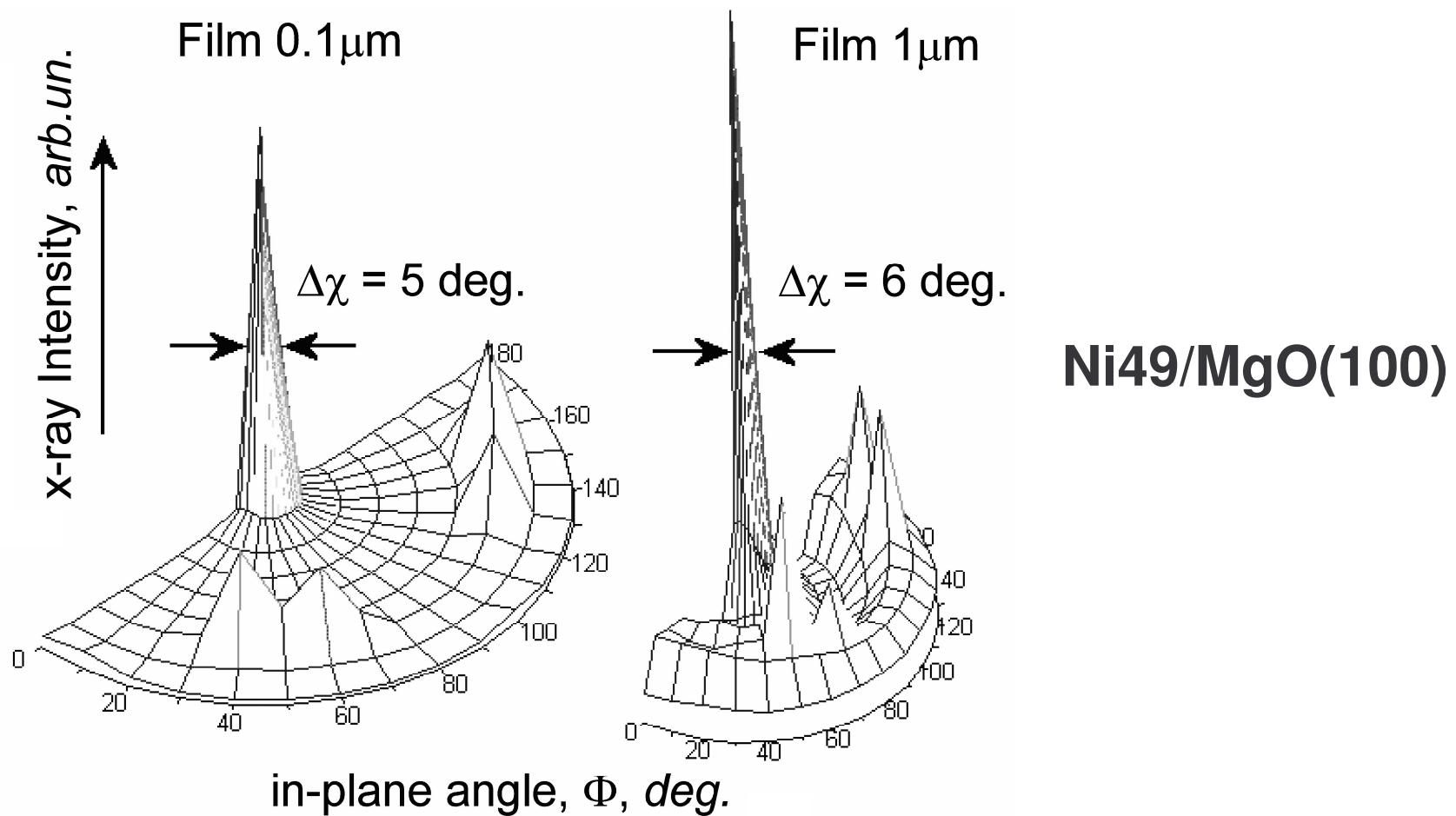


Calculated

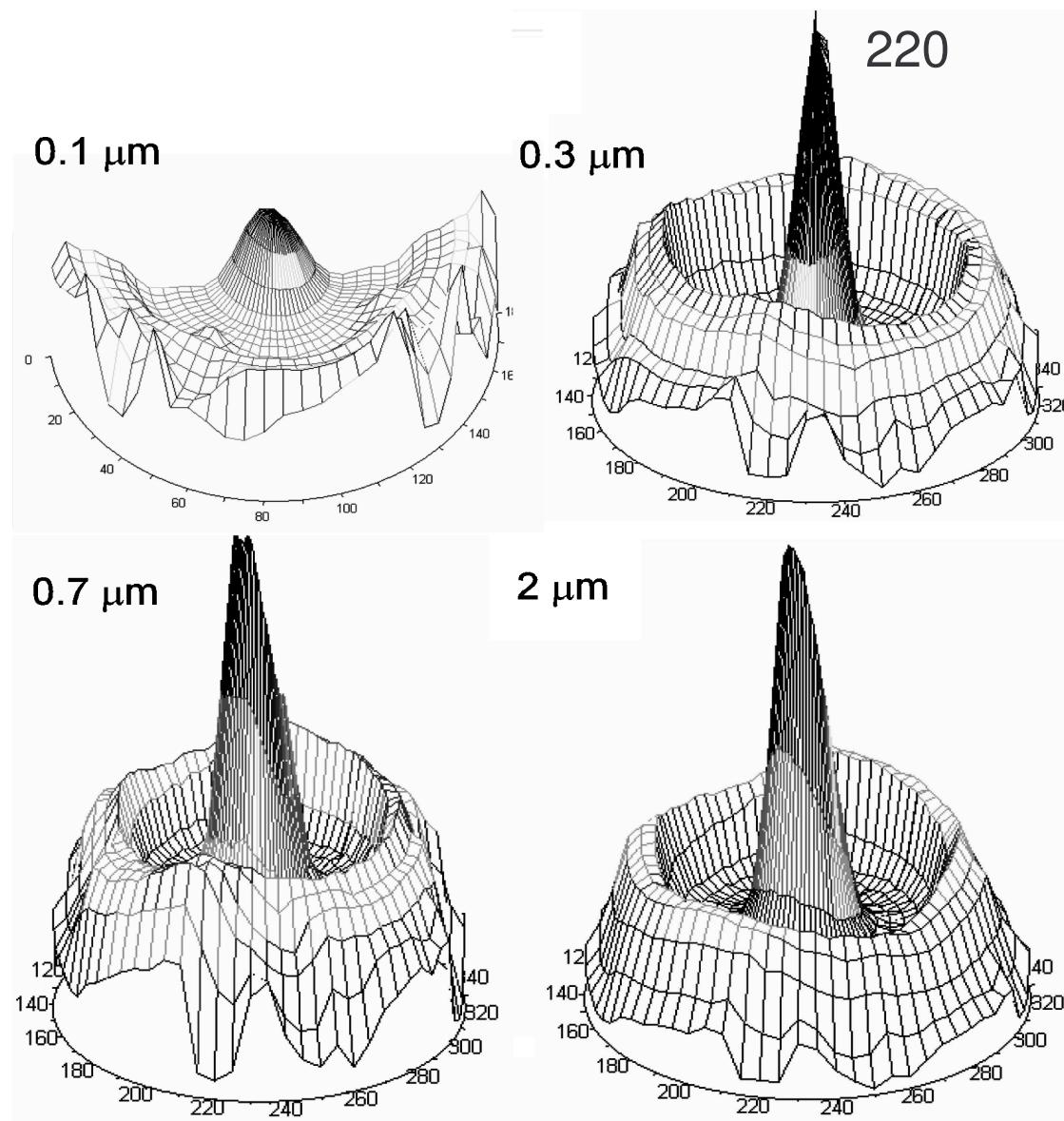


Chernenko et al., SMS 2005

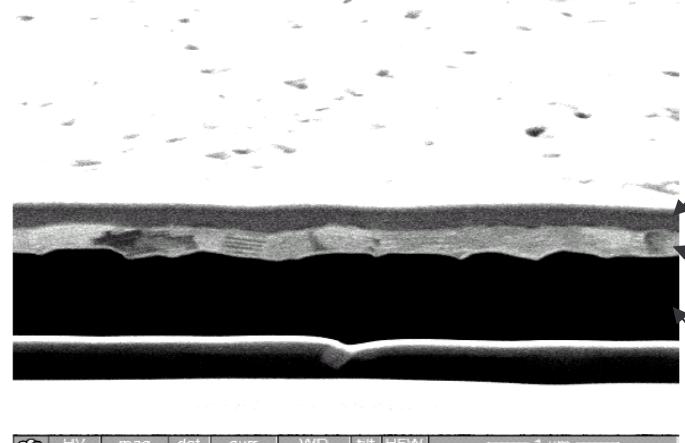
Texture studies by X-ray beamline of ANKA synchrotron source (FZK, Karlsruhe)



Fiber texture Ni49/Si(100)

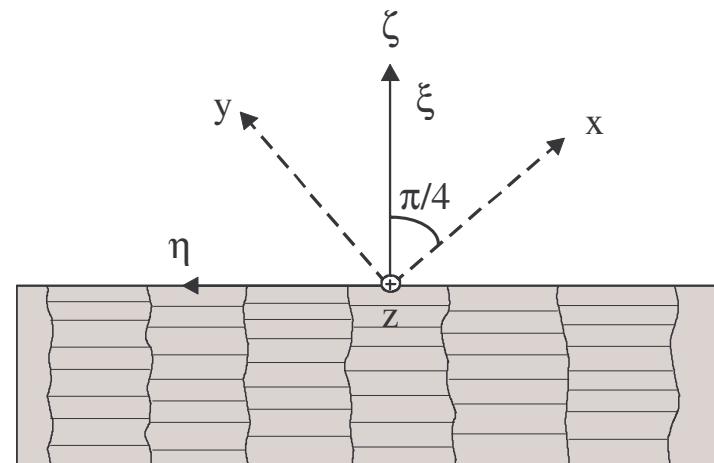
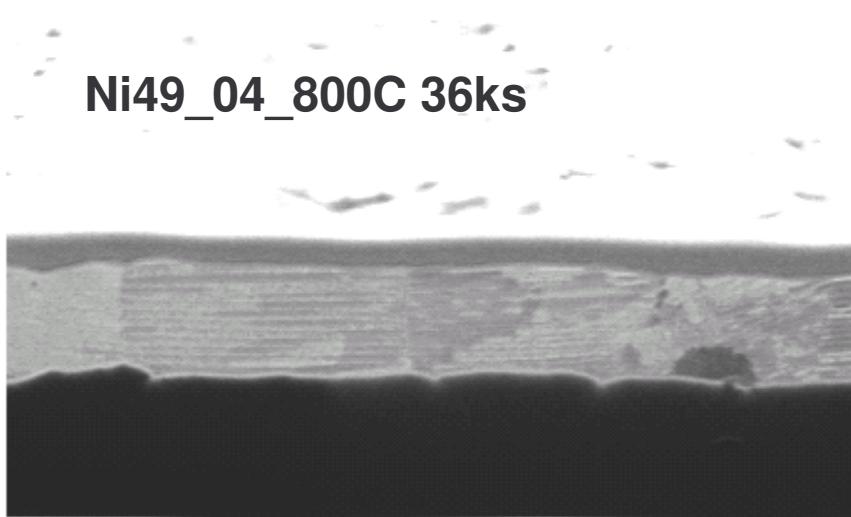


FIB images of cross-sectional areas of Ni49/Al₂O₃ film composites



Pt Model of columnar structure of twinned martensite

film
Al₂O₃

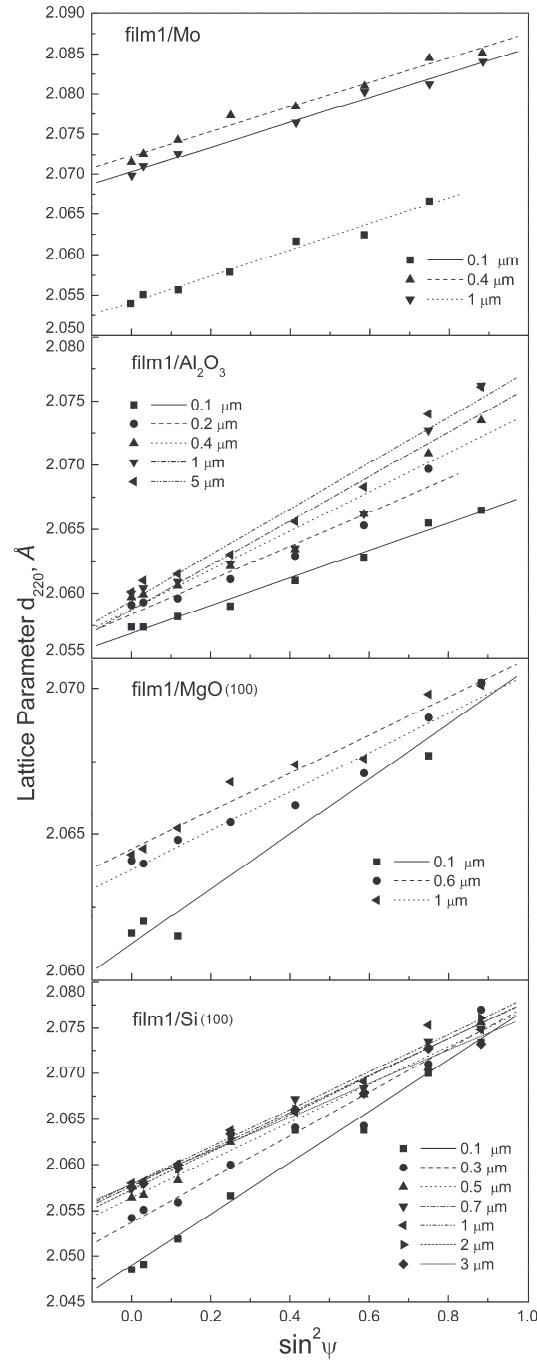


Ni49_1_800C 36ks

HV 30 kV mag 50000 x det SED curr 10.0 pA WD 19.8 mm bit 0.7 µm HFW 5 µm 1 µm

Chernenko et al., Mat.Trans.2006

Ni49/substrate

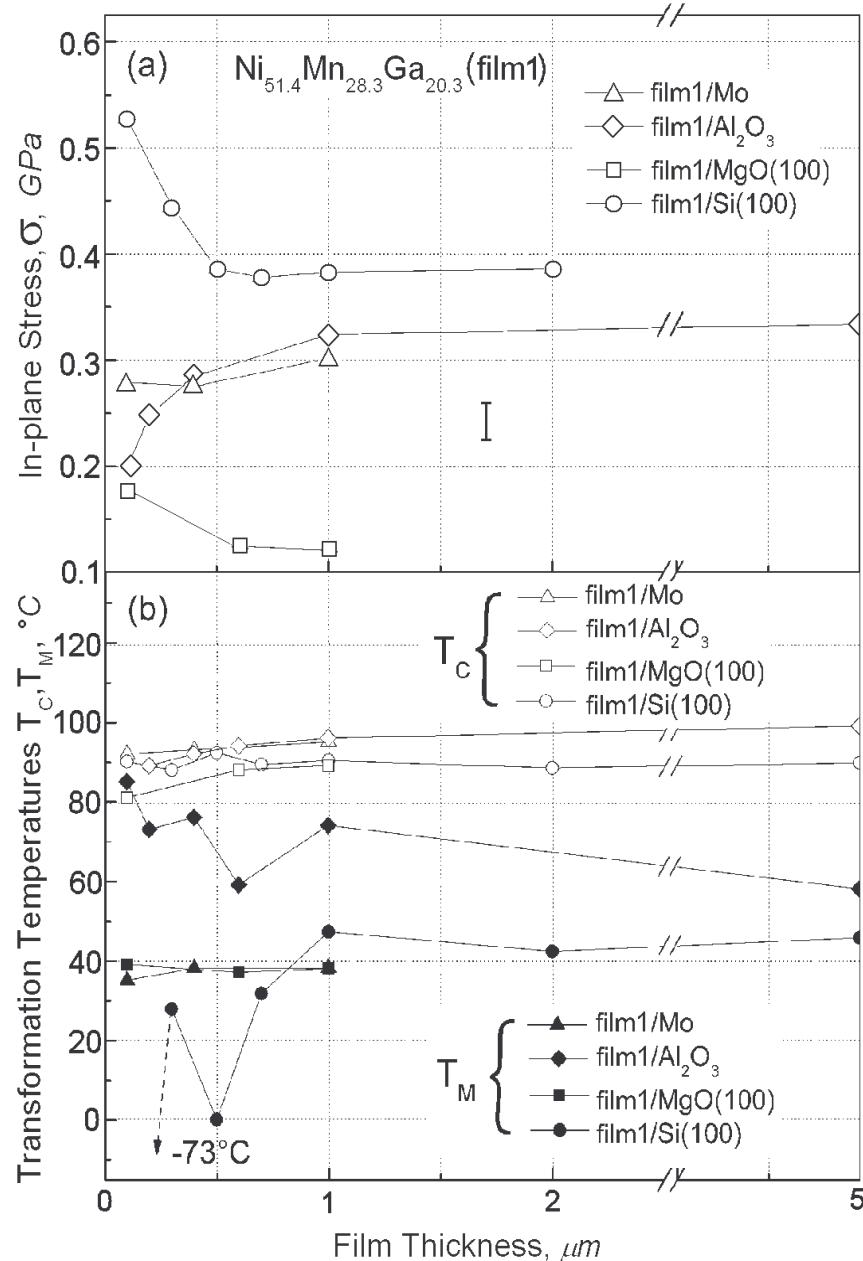


X-ray stress analysis

$$\sigma = \left(\frac{E}{1+\nu} \right) \frac{1}{d_o} \left(\frac{\partial d_{220}}{\partial \sin^2 \psi} \right)$$

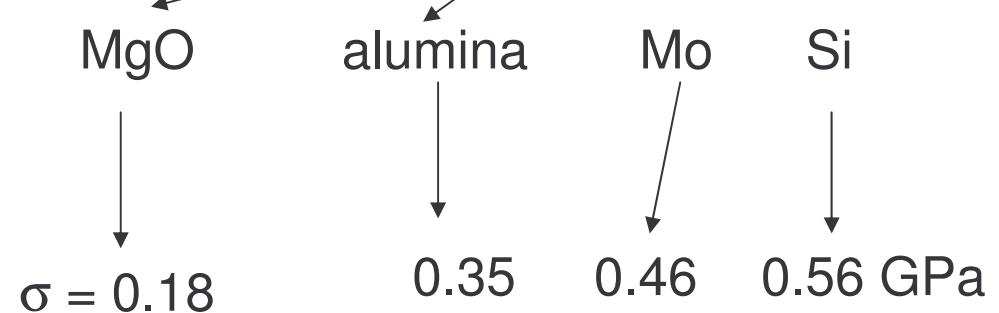
$$E = 50 \text{ GPa}, \nu = 0.3$$

Correlation: in-plane film stress and transformation temperatures of Ni49/substrate

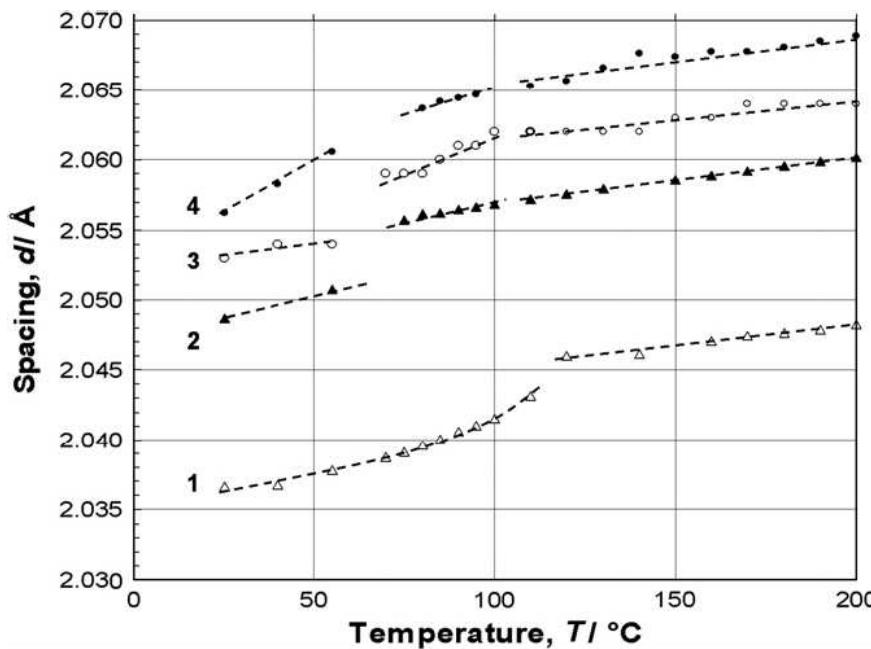


$$\sigma = E(\Delta\alpha\Delta T), \quad \Delta T = 650 \text{ K}$$

where $\alpha_f = 15, \alpha_s = 12.5; 7.5; 5$ and $3 \times 10^{-6} \text{ K}^{-1}$



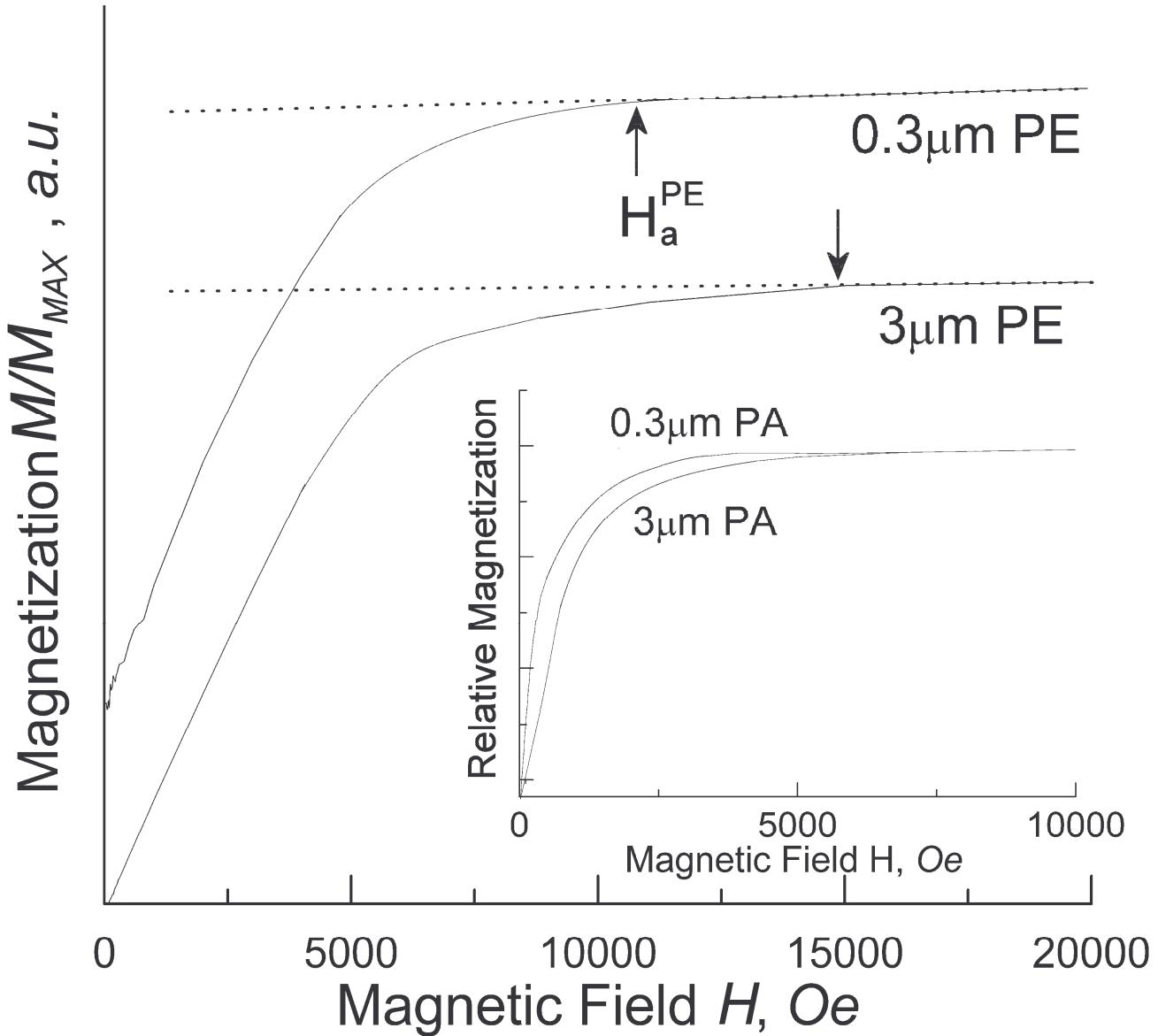
Thickness dependence of out-of-plane elastic modulus



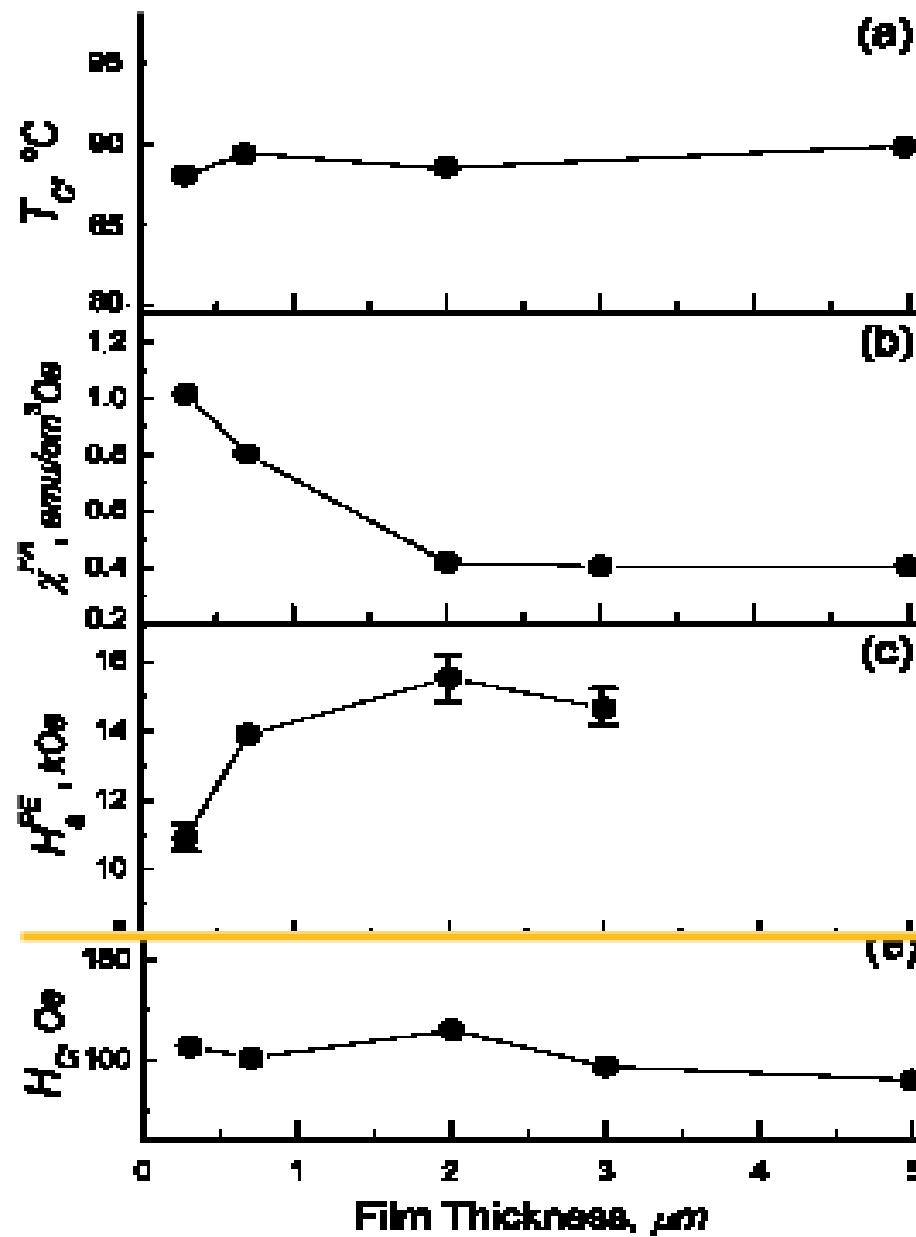
Transformation temperature A_f and d spacing measured at $200\text{ }^\circ\text{C}$ for $\text{Ni}_{51.4}\text{Mn}_{28.3}\text{Ga}_{20.3}$ films of different thickness h deposited on alumina ceramic substrate

h (μm)	A_f ($^\circ\text{C}$)	ΔT ($^\circ\text{C}$)	$\Delta\sigma$ (MPa)	d (nm)	$\Delta\varepsilon$	E (GPa)
5	65	0	0	0.206891	0	—
0.6	64	1	~0	0.206400	0.00237	—
1	78	13	130	0.206022	0.00426	31
0.1	114	49	490	0.204826	0.00998	49

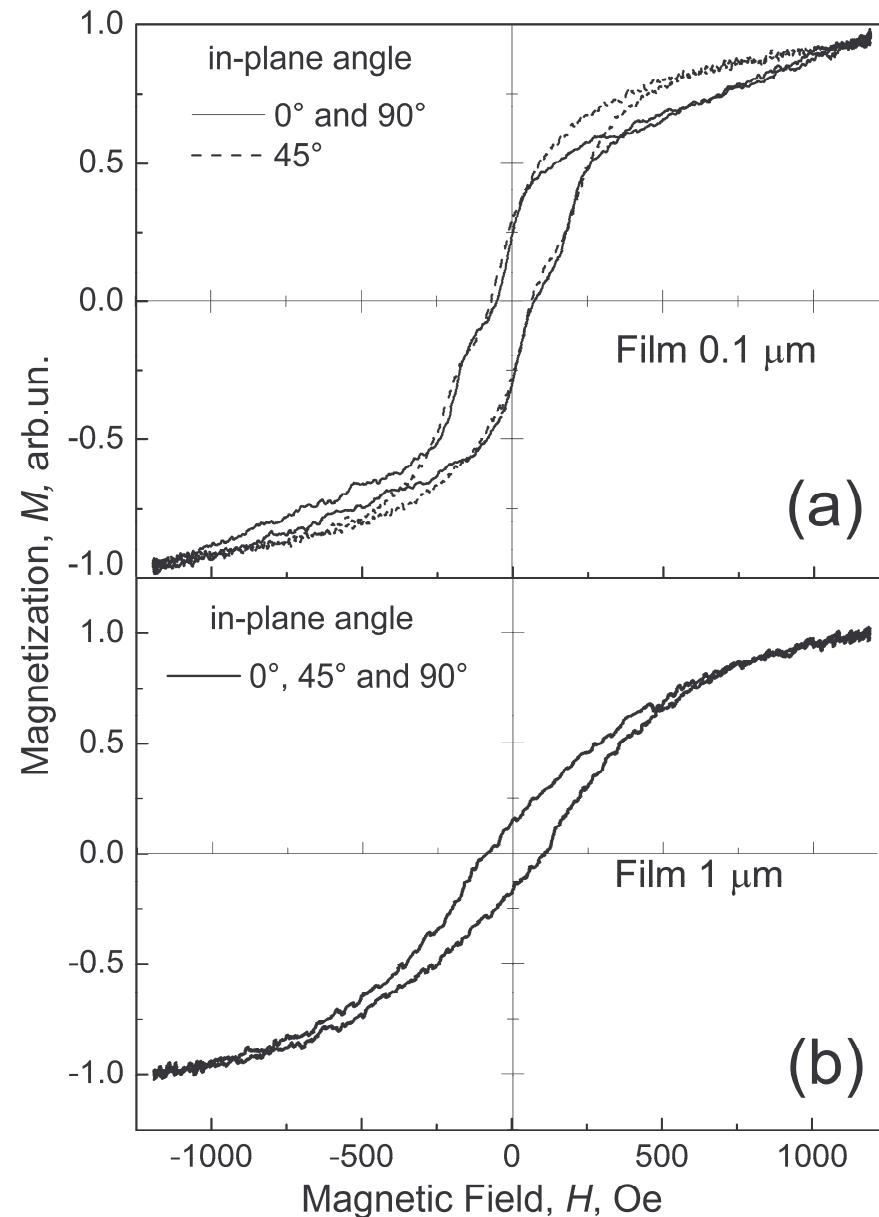
Magnetization of Ni49/Si(100) composite



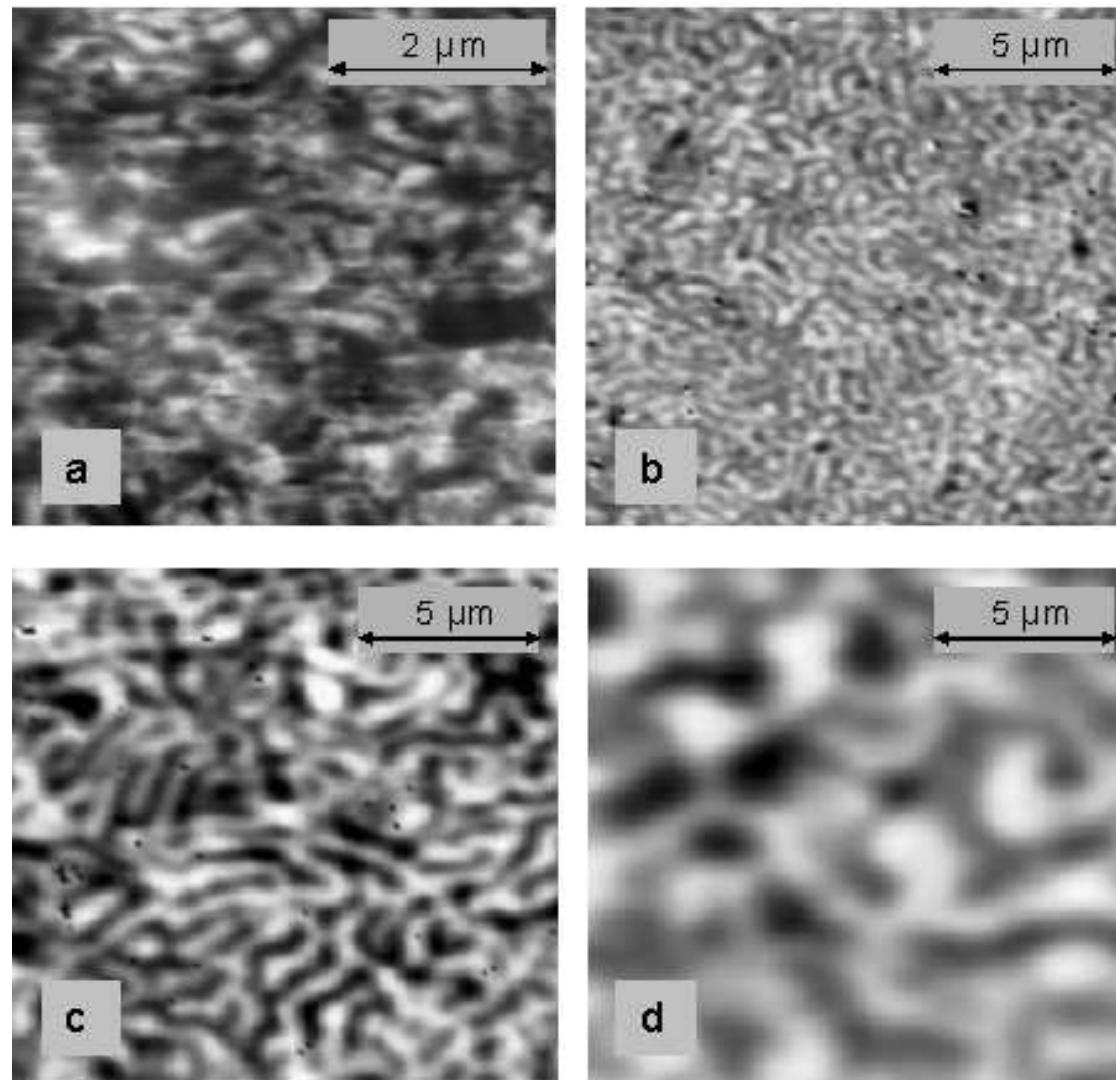
Thickness dependence of magnetic parameters for Ni49/Si



In-plane magnetic anisotropy for 0.1 μm NiMnGa/MgO film



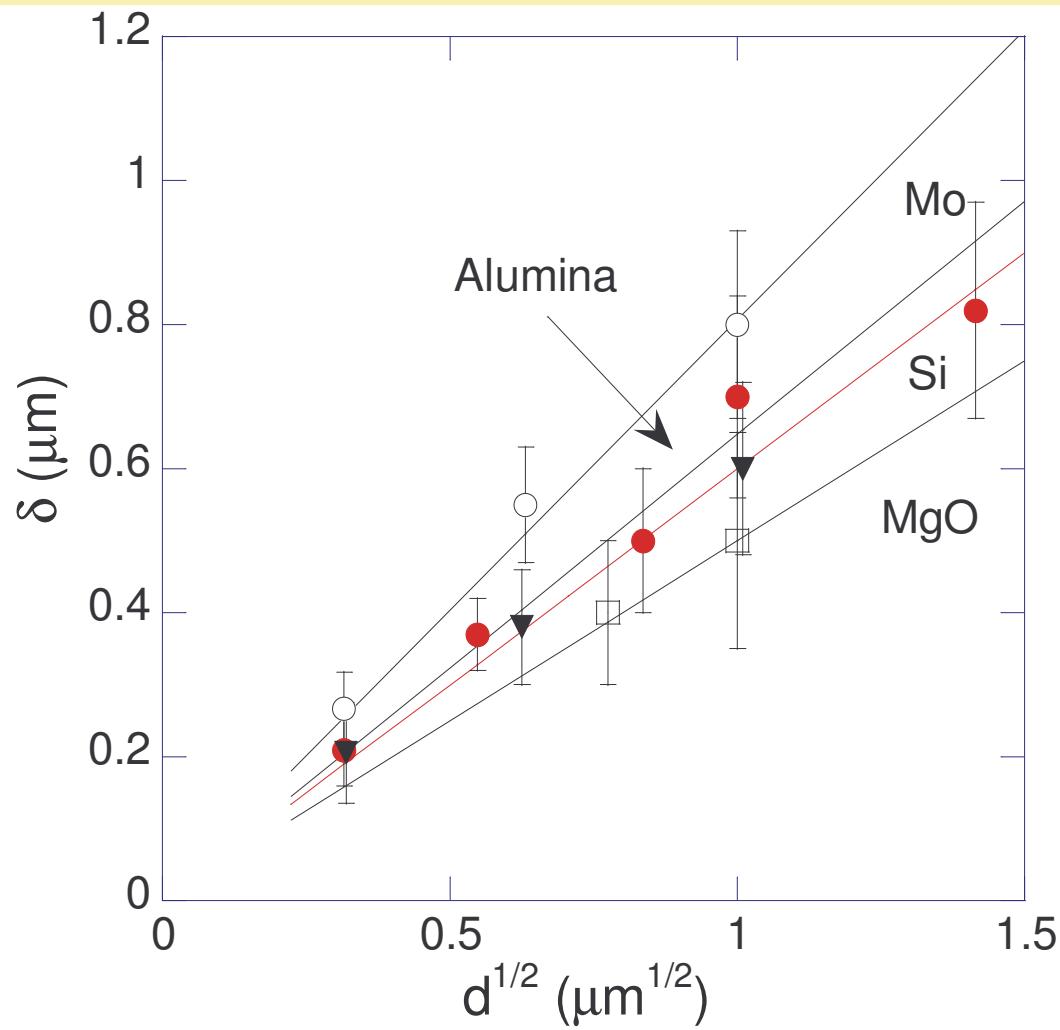
Magnetic Force Microscopy, H=0



$\text{Ni}_{49.5}\text{Mn}_{28.0}\text{Ga}_{22.5}/\text{Al}_2\text{O}_3$; $d = 0.1 \mu\text{m}$ (a),
0.4 μm (b); 1.0 μm (c); 5.0 μm (d)

Chernenko et al., JPCM 2005

Magnetic domain size: influence of substrate nature and microstructure



$$\delta = \sqrt{d\delta_{dw}}$$

$$\delta_{dw} = a\sqrt{E / A}$$

d<0.4 μm: Domain size becomes larger than film thickness

Analysis of χ vs angle between H and film normal and determination of :
 (a) angle between magnetization and film normal 67 deg
 (b) Q=2.9
 for Ni49/alumina 0.1 μm thick film

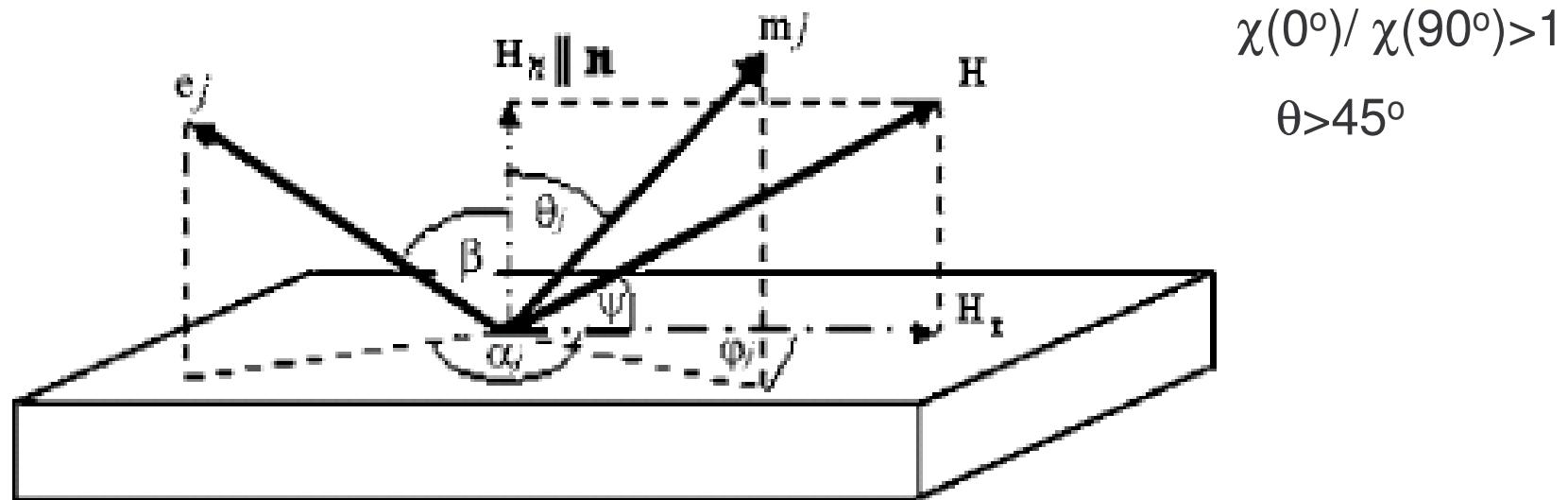
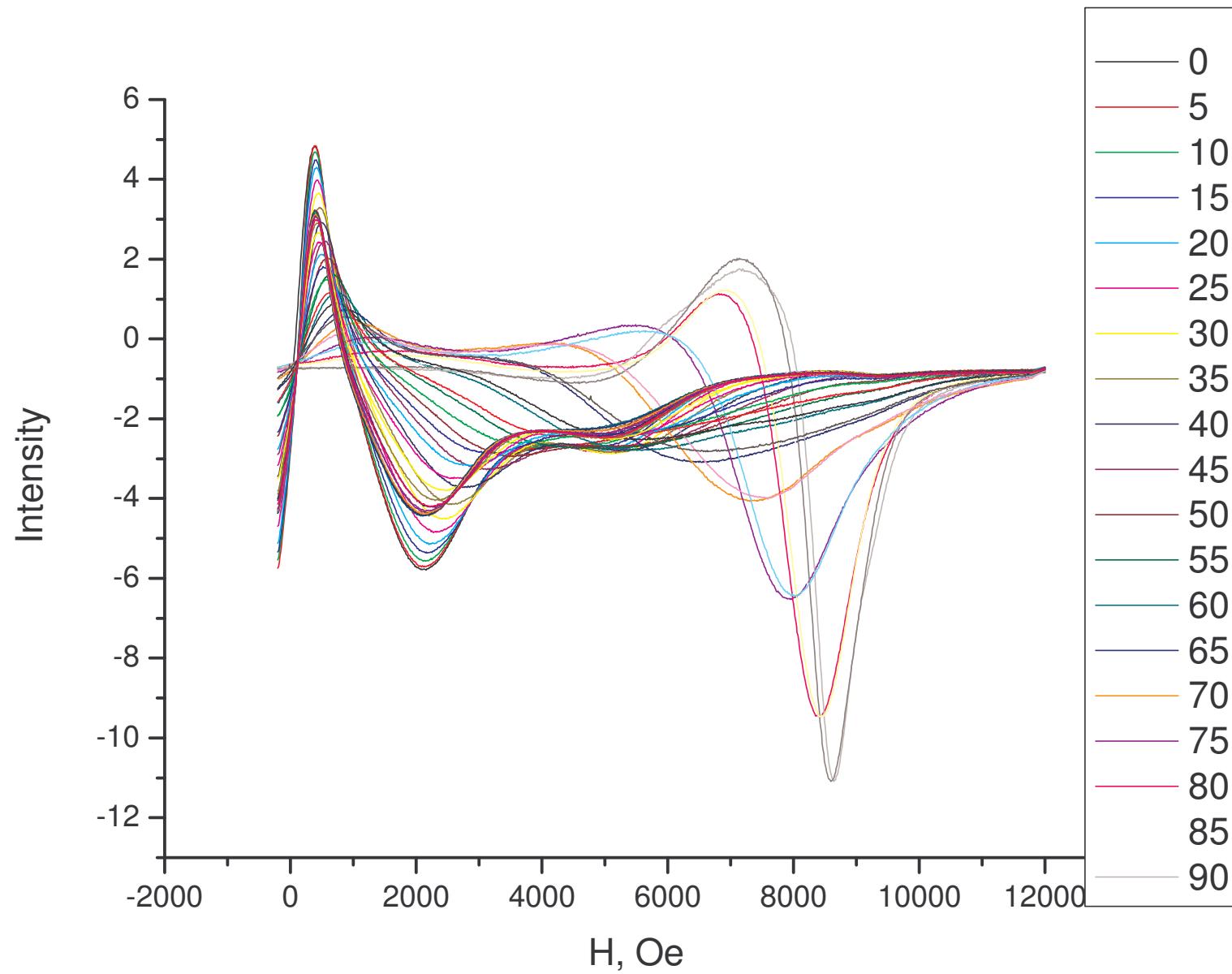


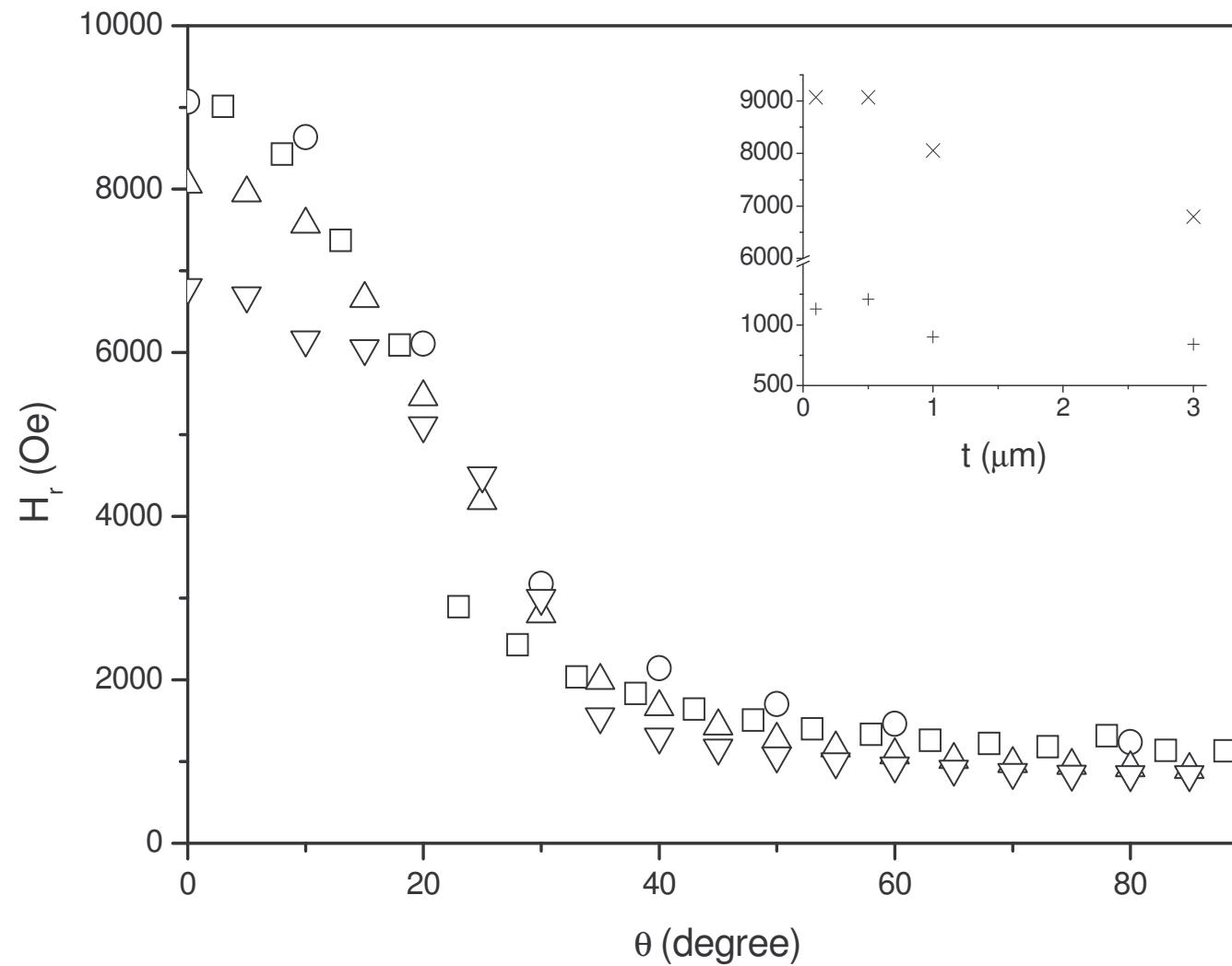
FIG. 3. The angles prescribing the directions of the external magnetic field ψ , easy-axis direction β , and magnetic vector position θ with respect to film normal or film plane.

FMR spectra for Ni49_1_Si

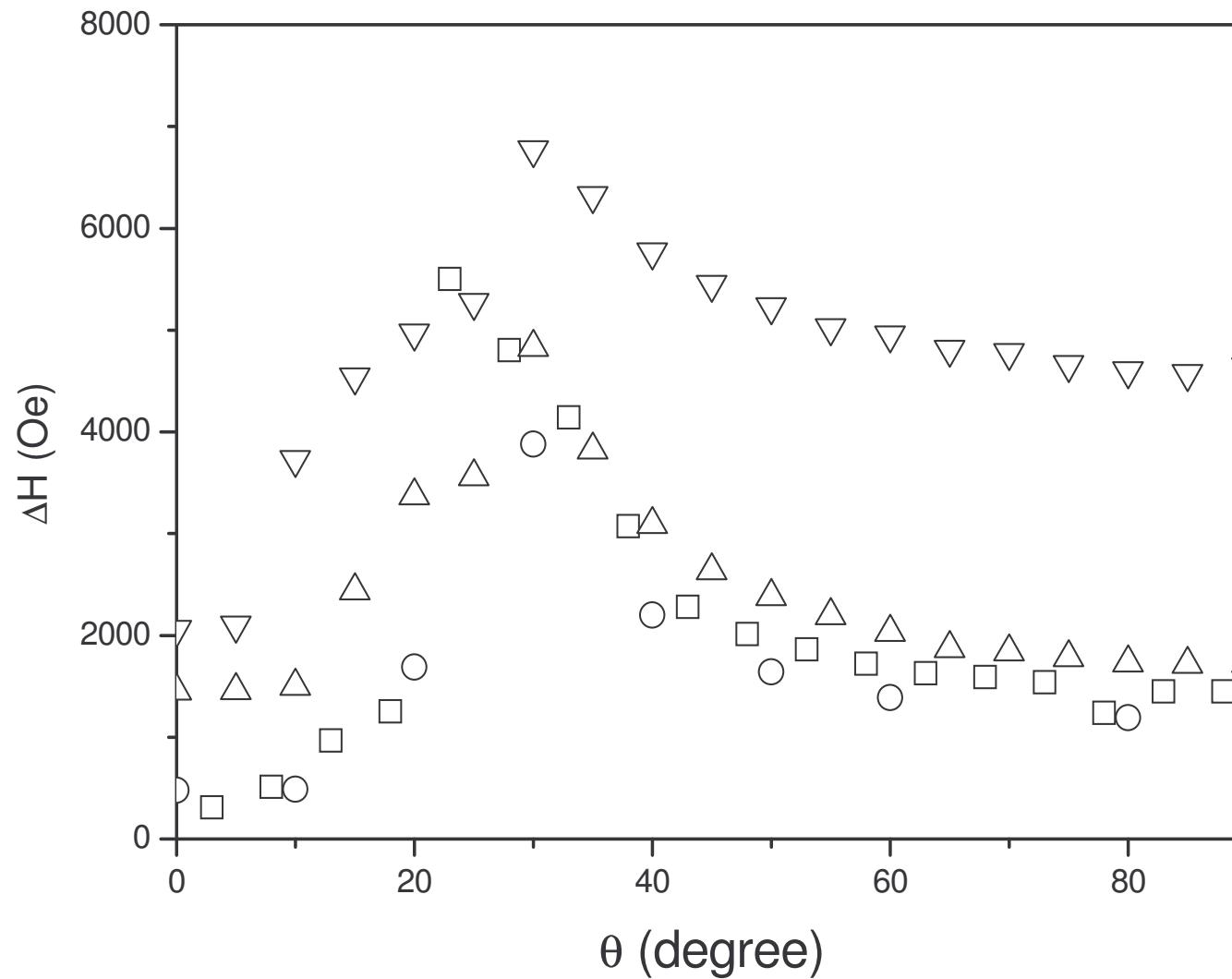


Resonance field vs angle between film normal and H

Inset: resonance field vs film thickness



Resonance linewidth vs angle



Modeling:

using $M_{sat}=500$ emu/cm³,

$g=2$

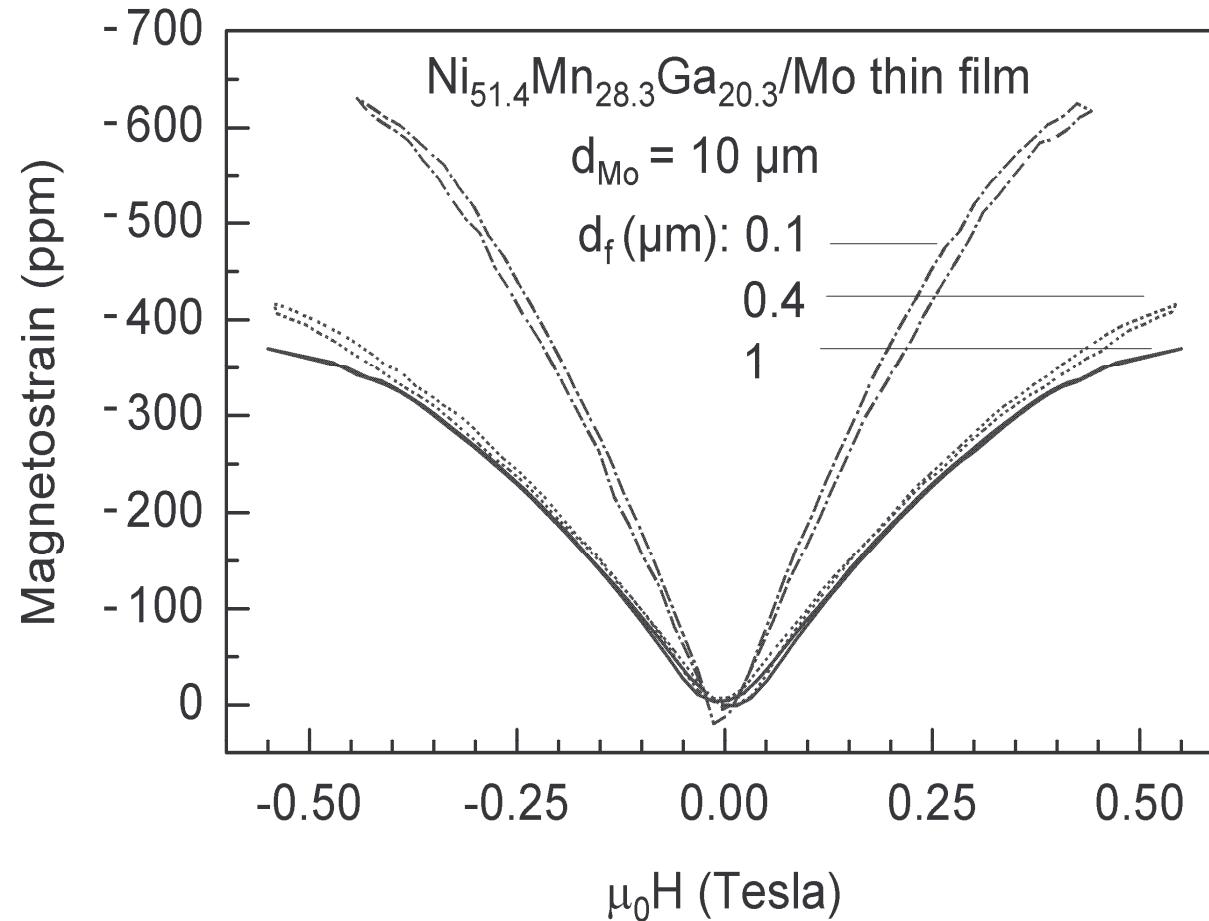
uniaxial anisotropy making an angle 45 degree with respect to film normal



values of uniaxial anisotropy constant for films Ni49/Si

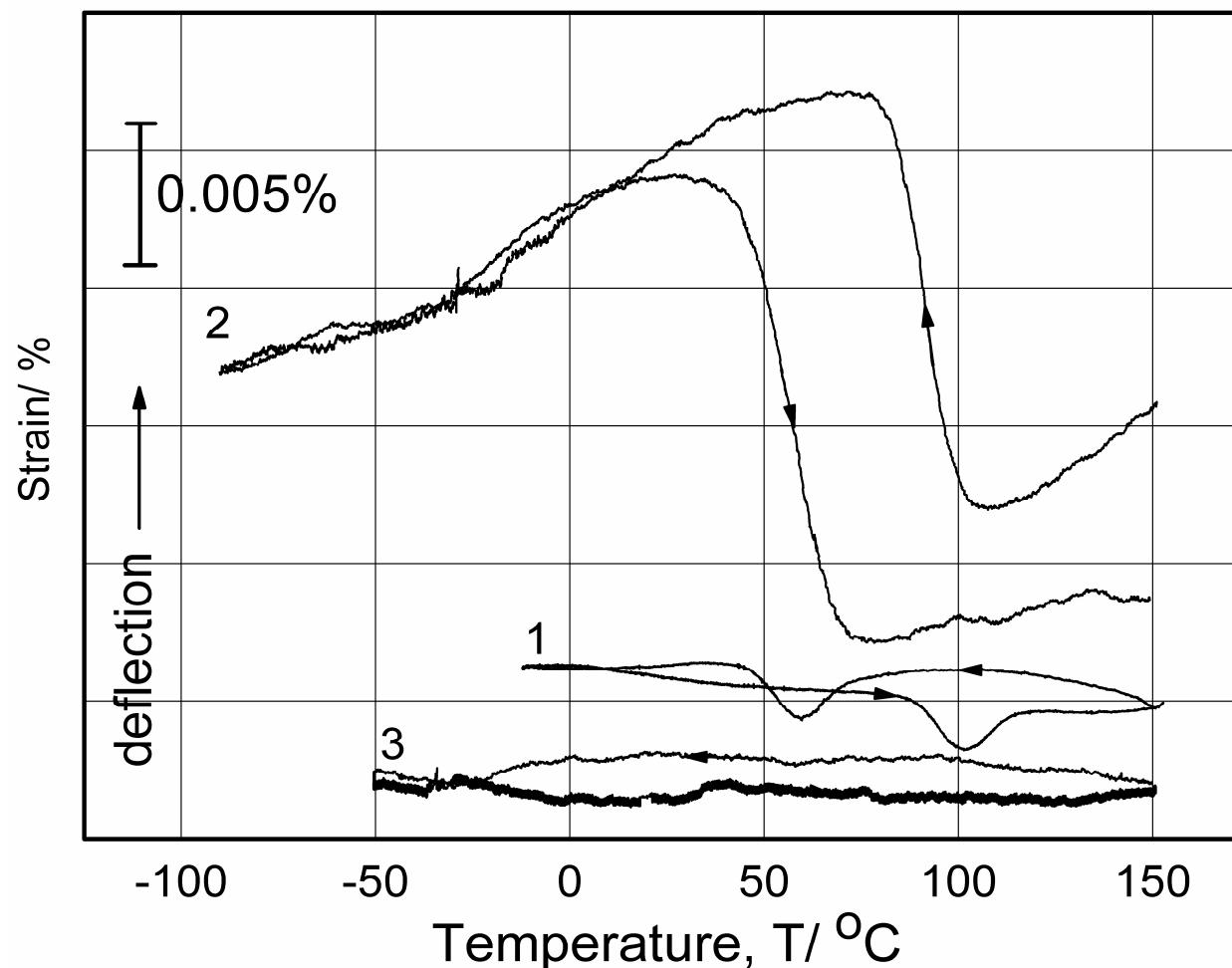
0.1 μm	0.5 μm	1 μm	3 μm
$2.8 \cdot 10^5$ erg/cm ³	$2.8 \cdot 10^5$ erg/cm ³	$4.2 \cdot 10^5$ erg/cm ³	$5.1 \cdot 10^5$ erg/cm ³

Magnetic activation of Ni49/Mo10



Thermally activated cantilever Ni49(5μm)/alumina(150μm) actuator

- 1 - no external load is applied;
- 2 - cantilever is loaded by the external stress of 150 MPa.
- 3 - alumina substrate cantilever loaded by the stress of 150 MPa.



Work in progress on FSMA's/substrate thin films composites

TEM and MFM imaging of films cross-section

X-rays grazing incidence studies

Deposition of films on heated substrates: epitaxial growth

Measurements of elastic modules as a function of film thickness

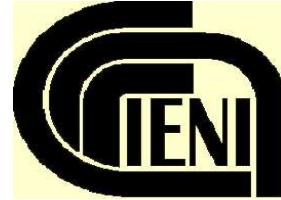
Modeling

Sensing and actuation properties



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Ferromagnetic Shape Memory Alloys: underlying physics and recent advances in thin film technology

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Milestones/Breakthroughs in MSMA history

Period	Country, Key Personnel	Conceptual approach
1984	UK Webster,Ziebeck et al.	Influence of Structural instability on magnetism of Ni ₂ MnGa Heusler compound
1989-1995	Ukraine, Spain Chernenko, Kokorin, Cesari et al.	Elaboration of Ni-Mn-Ga alloy system as novel ferromagnetic shape memory alloys
1996-1997	Finland, Ukraine, USA Ullakko, O'Handley et al.	Giant magnetostriction caused by twin boundary motion

Resulting in:

1998-present

Avalanche-like progress due to extension of experimental technique, modeling and developing other alloy systems

1993

**Ni₂MnGa AS A NEW FERROMAGNETIC ORDERED
SHAPE MEMORY ALLOY**

Vladimir A. Chernenko and Vladimir V. Kokorin

Institute of Metal Physics
Vernadsky str. 36
Kiev 252680 UKRAINE

ABSTRACT

Ni₂MnGa alloy was found to exhibit the thermoelastic martensitic transformation accompanying by anomalies of the electric, thermal, elastic and magnetic properties. The cubic to tetragonal cubic transformation was confirmed.

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MAGNETOPLASTIC DEFORMATION OF Dy CRYSTALS

1974

ABSTRACT

Several authors have reported that if single crystals of Dy or Tb are magnetized in fields of order 100 kOe parallel to a magnetic hard direction at 4.2 K, a permanent plastic deformation occurs. Neither the

1996

Large magnetic-field-induced strains in Ni₂MnGa single crystals

K. Ullakko, J. K. Huang, C. Kantner, and R. C. O'Handley

*Massachusetts Institute of Technology, Cambridge,
Massachusetts 02139*

V. V. Kokorin

Institute of Metal Physics, Kiev, Ukraine

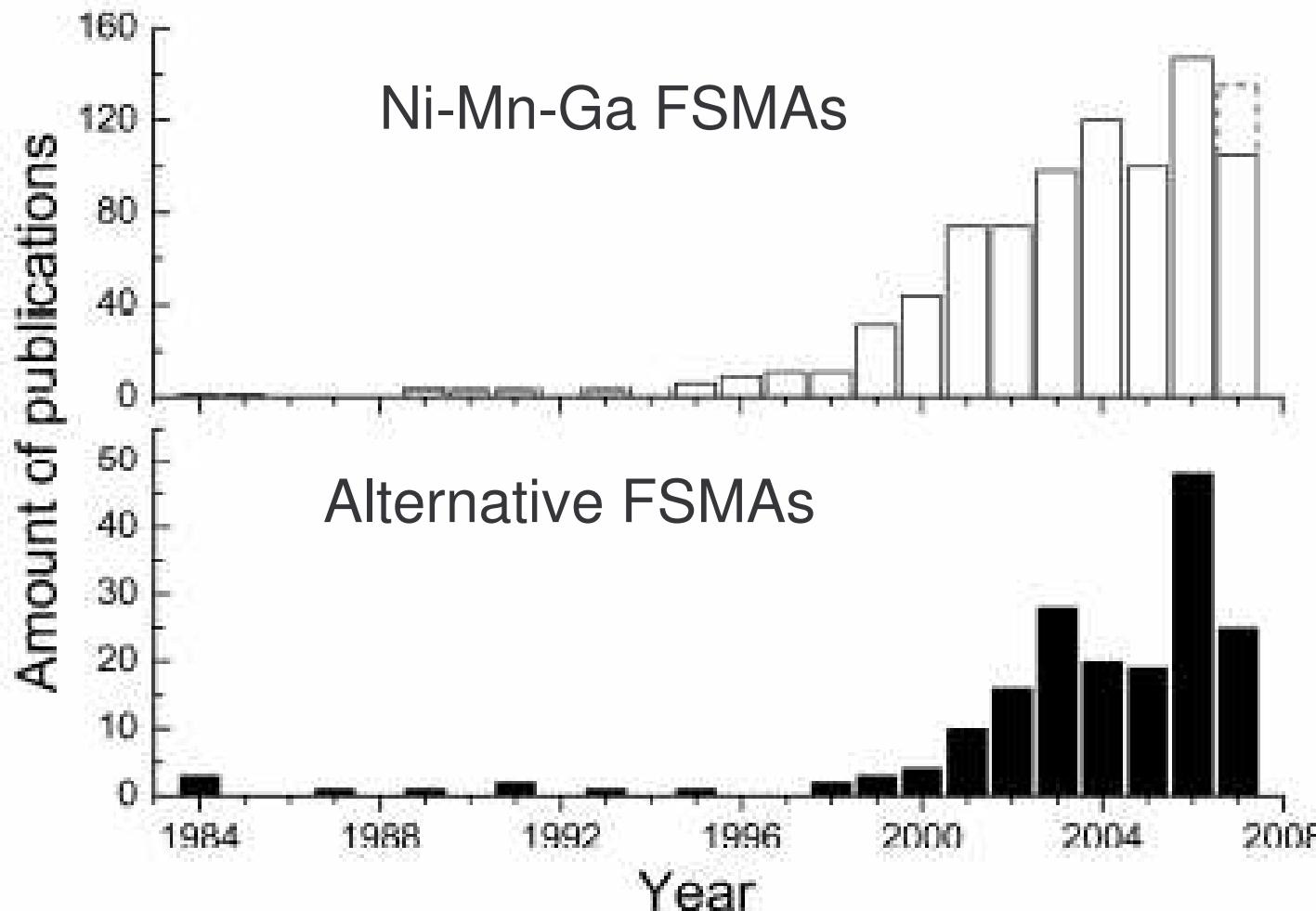
~Received 3 June 1996; accepted for publication 29 July 1996!

Strains of nearly 0.2% have been induced along @001# in unstressed crystals of Ni₂MnGa with magnetic fields of 8 kOe applied at 265 K. These stains are associated with the superelastic motion of twin boundaries in the martensitic phase that is stable below about 274 K. © 1996 American Institute of Physics. @S0003-6951~96!02439-4#

samples were spark-cut from the resulting button.

A major experimental problem was to devise a satisfactory method for holding the sample in the magnetic field. The crystal had to be free to deform, but also had to be held against large mechanical torques resulting from even slight misalignment. For example, if the sample is rotated 10° from its

FSMA science – vast interdisciplinary research field



Totally more than 1000 publications on FSMA

Materials science

Advances in Shape Memory Materials

Ferromagnetic Shape Memory Alloys

Editor V.A. Chernenko

Materials Science Forum Volumes 583
ISBN13: 978-0-87849-381-4, 302 pages, softcover,
Euro 167.00, USD 230.00

The thermoelastic martensitic transformations and shape memory materials have been in the focus of both solid state physics and materials science for a long time. The fundamental and applied aspects of this vast activity are extensively reviewed in the journal publications and conference proceedings. Several books and monographs on this subject, such as the most recent K. Otsuka and C.M. Wayman, eds. *Shape Memory Materials* (Cambridge University Press, Cambridge, 1998) and M. Kohl *Shape Memory Microactuators* (Springer, 2004), are also available in the literature.

Although structurally very similar to one typical for non-magnetic shape memory alloys, the thermoelastic



Advances in Shape Memory Materials
Ferromagnetic shape memory alloys

Edited by
V.A. Chernenko

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2-nd International Conference on Ferromagnetic Shape Memory Alloys

ICFSMSA 2009, July 1-3, Bilbao, Spain



Guggenheim museum

FSMA science field is affluent by concepts which help understanding FSME and related phenomena

Some of them are:

lattice instability

soft-mode behavior

electron concentration

ferromagnetic shape memory effect (FSME)
magnetic-field-induced twin boundary motion

magnetostress

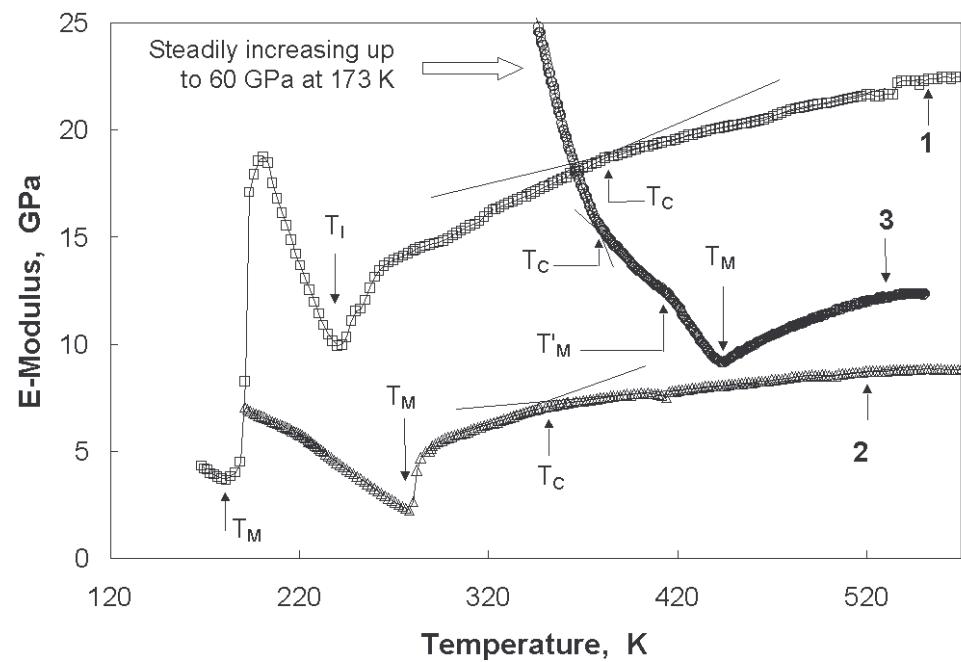
magnetomechanical mechanism

magnetic-field-induced superelasticity

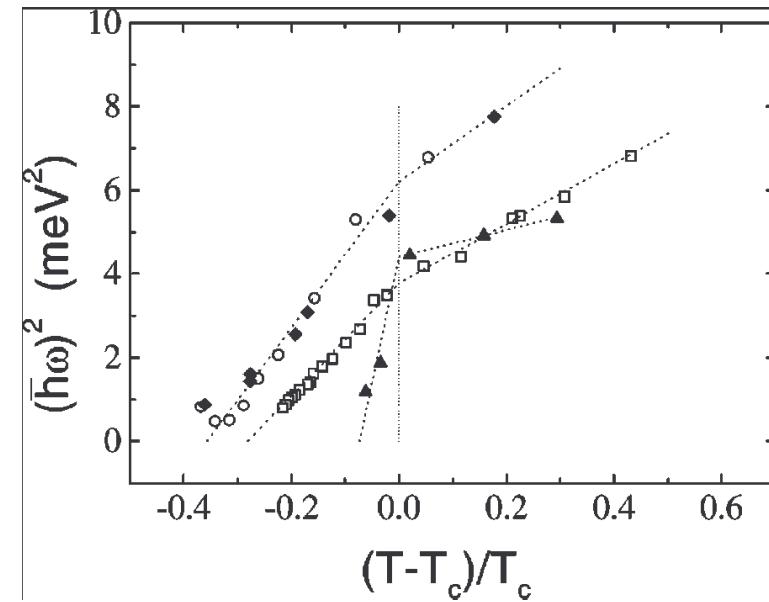
twinning-strain-induced change of magnetization

Modulus and phonon softening

TA2(1/3,1/3,0)

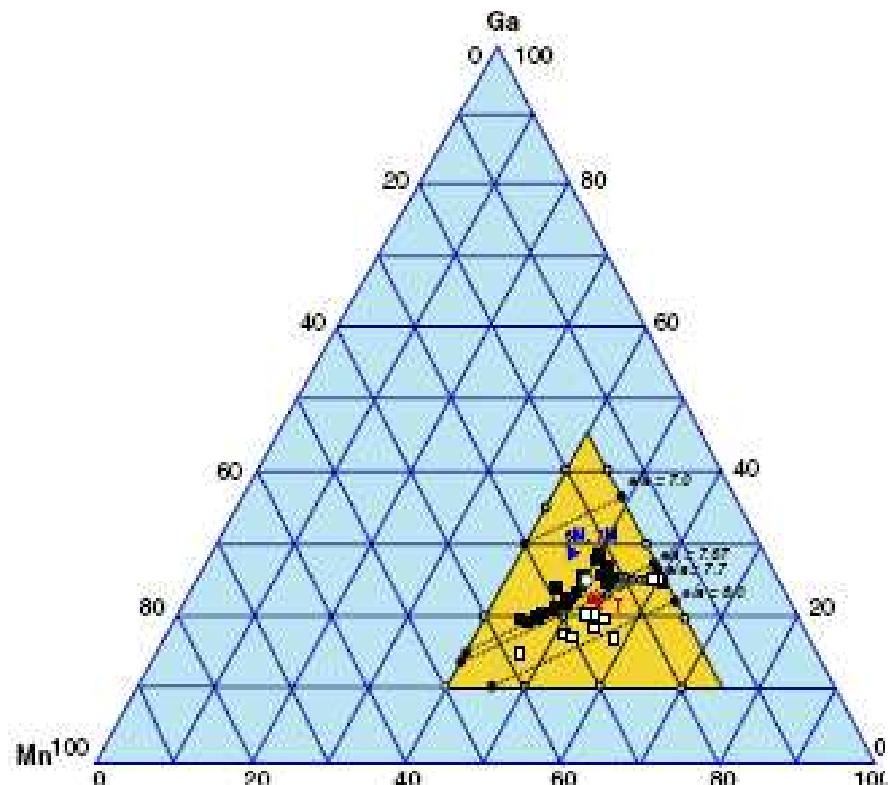


Chernenko et al. 2002

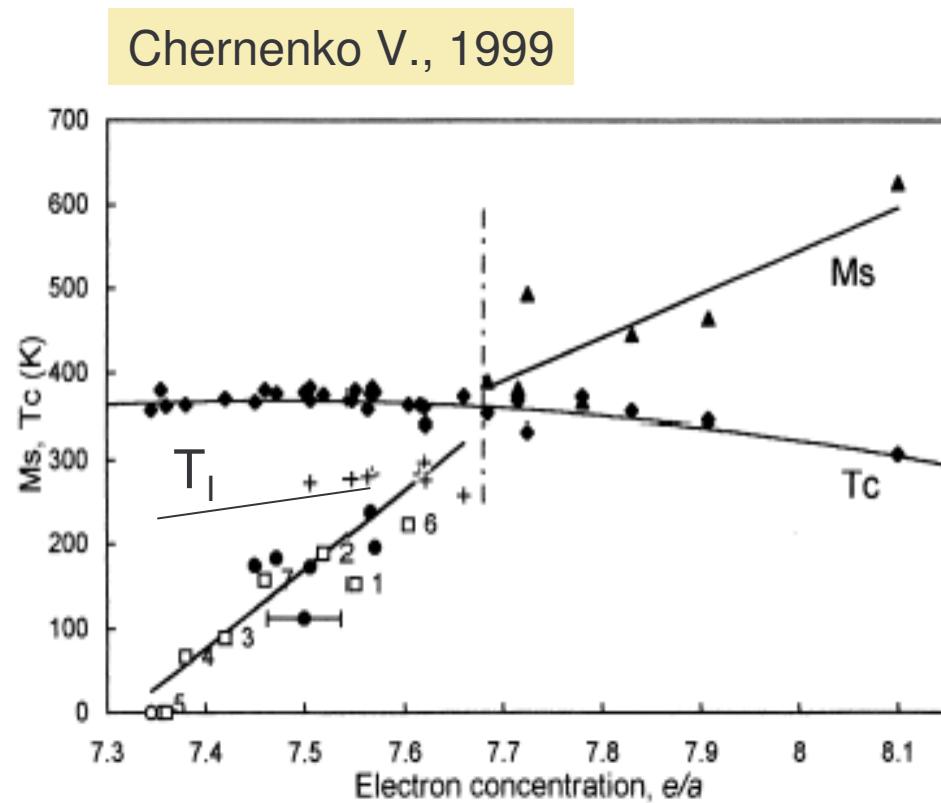


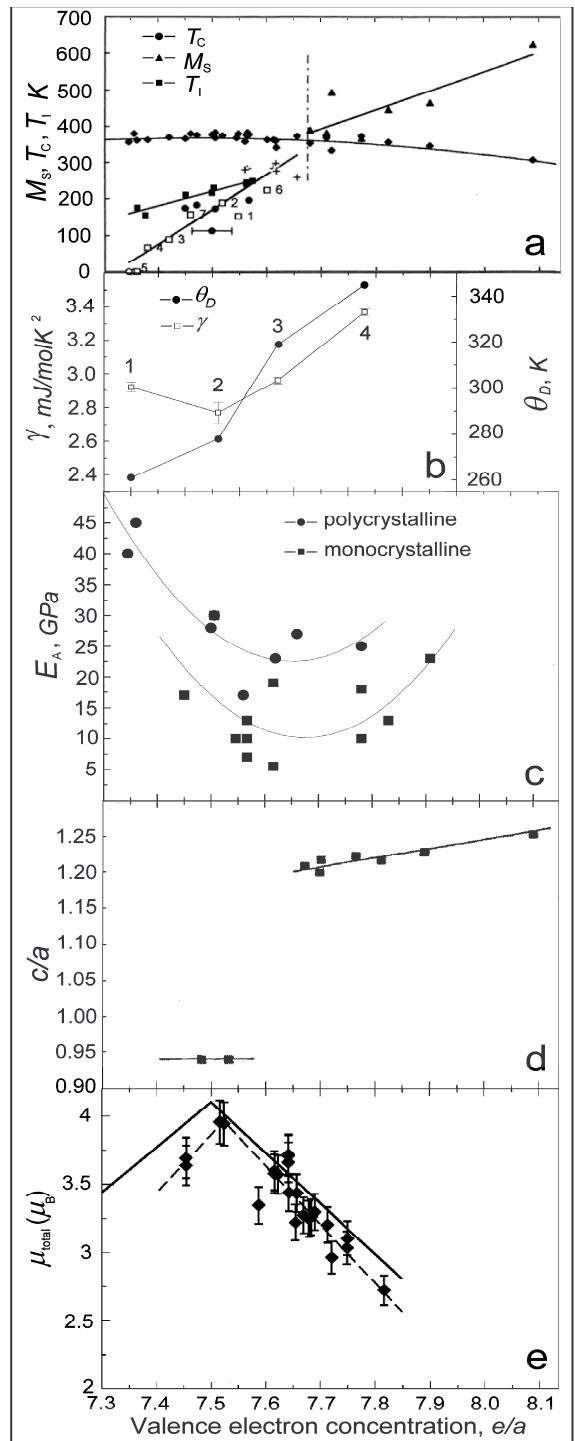
Manosa et al, 2001

Concentration dependence of transformations characteristics in Ni-Mn-Ga Heusler alloy system



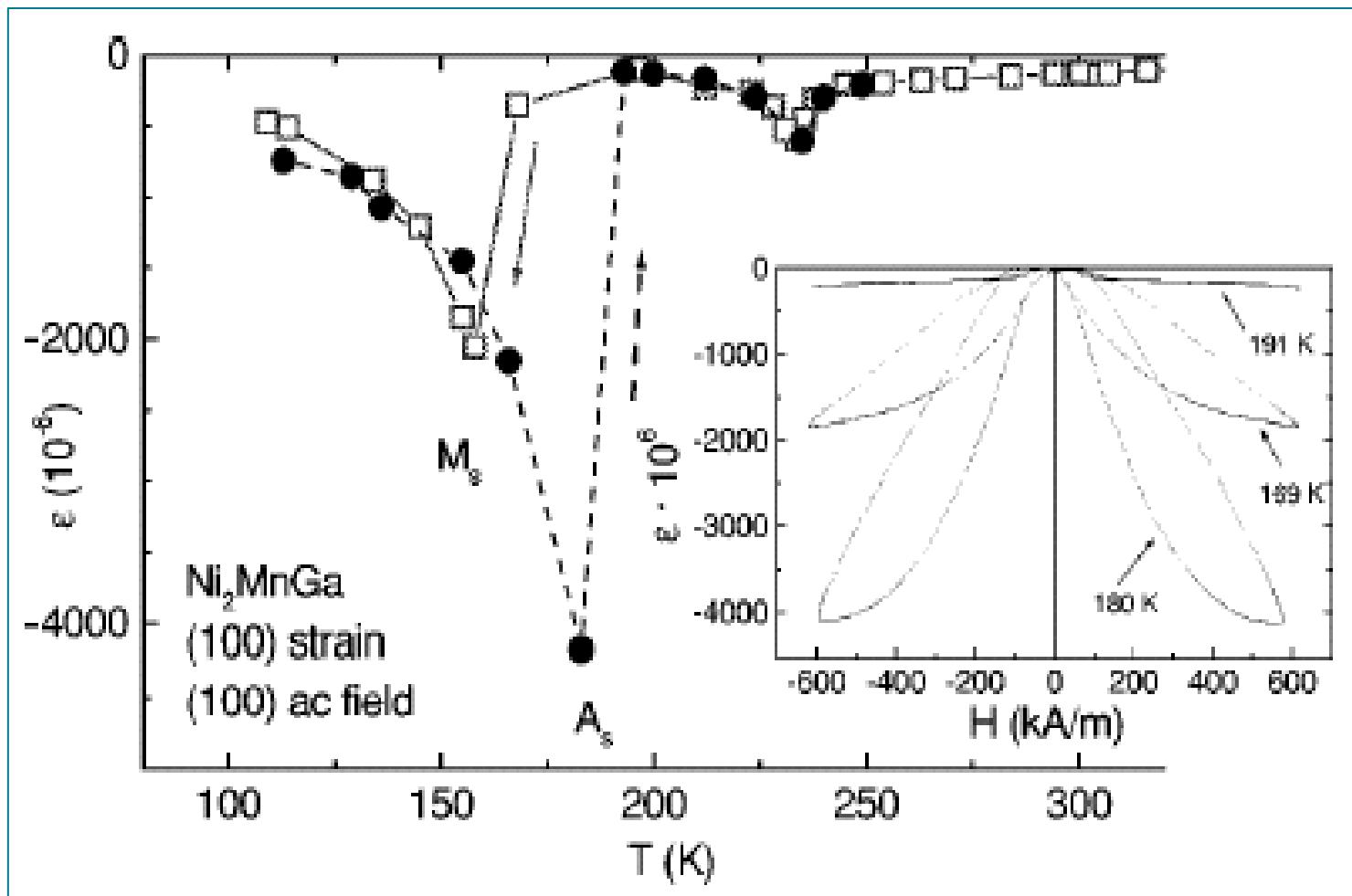
Chernenko V. et al., 1995



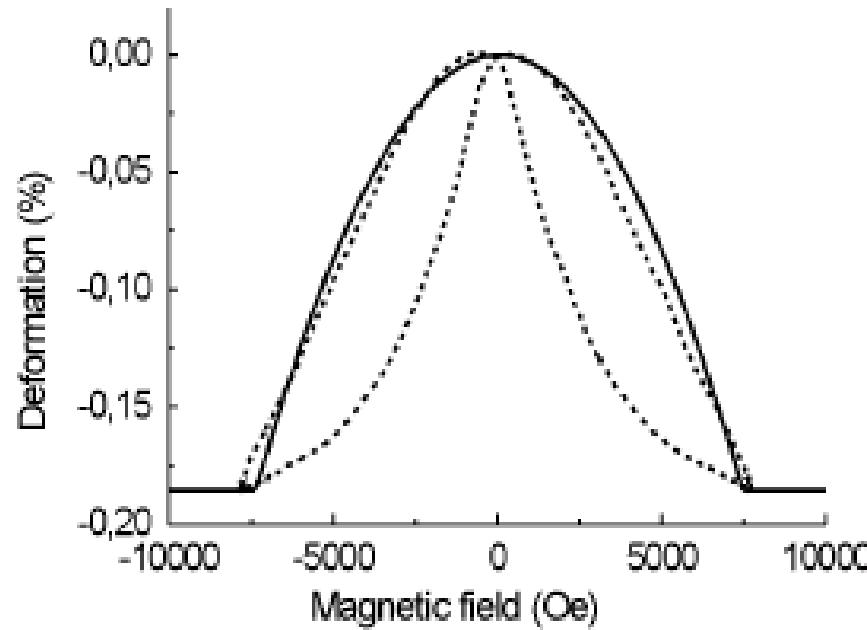


Effect of the valence electron concentration on the lattice, electronic and magnetic instabilities in Ni-Mn-Ga alloy system

Quasi-elastic magnetostain in polyvariant martensite



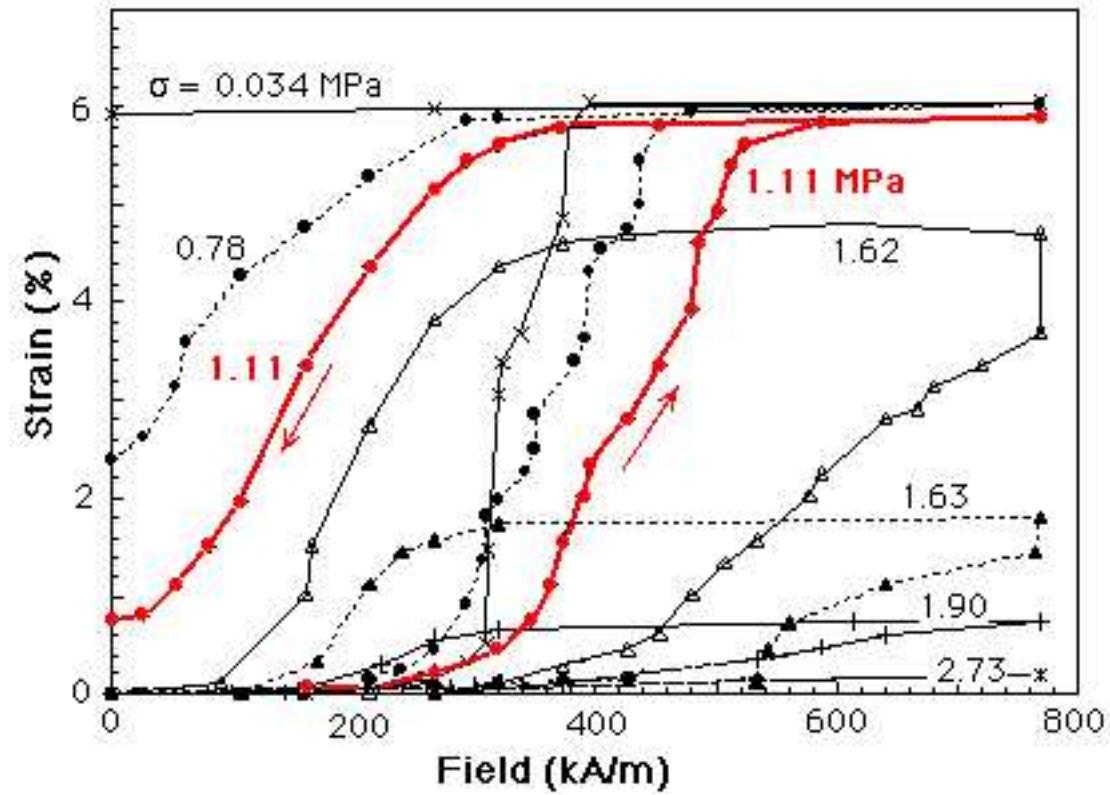
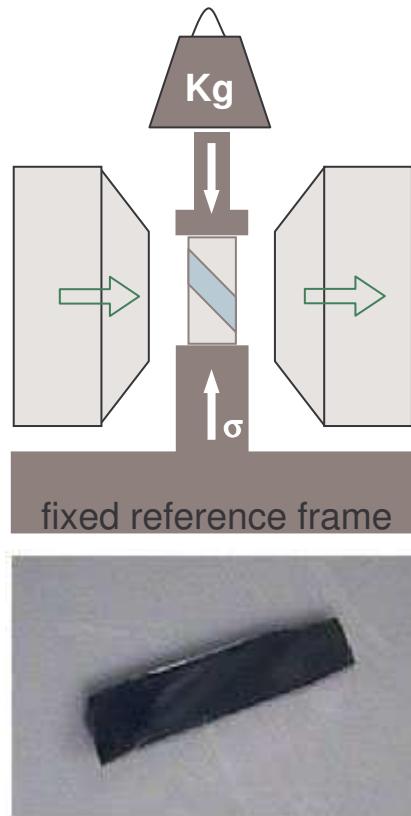
Quasi-elastic magnetostain: theory



[100] deformations induced by
field $H// [001]$ in Ni_2MnGa

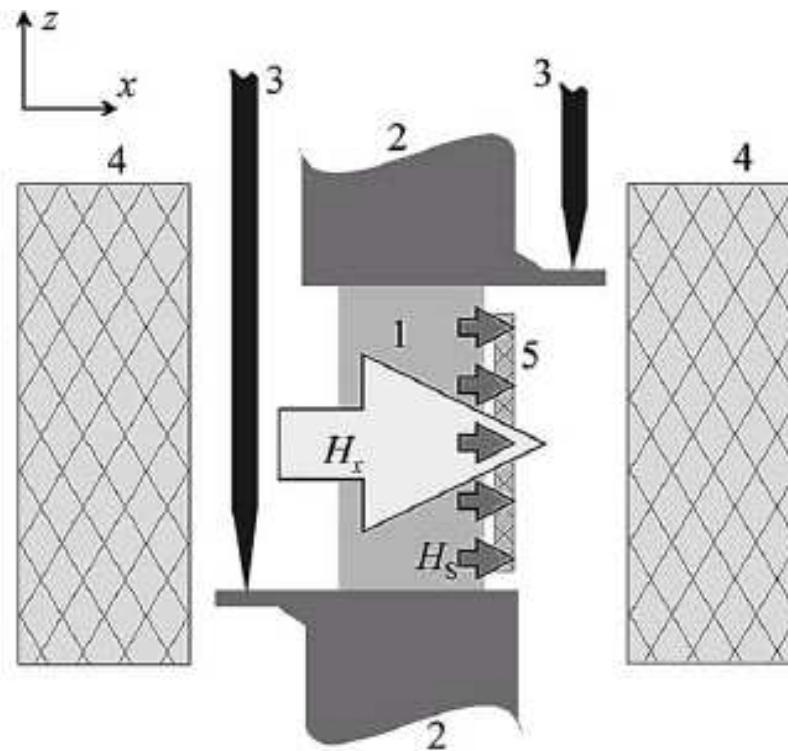
160 K, M=525 G, $\delta=-23$
 $(\partial \epsilon / \partial \sigma_3)^{-1} \approx 4.1 \text{ GPa}$

$$\langle \hat{\epsilon}(H, T) \rangle \approx \frac{2}{3} (\partial \hat{\epsilon} / \partial \sigma_3) 18 \delta M^2(T) \cos^2 \psi_{[001]}$$



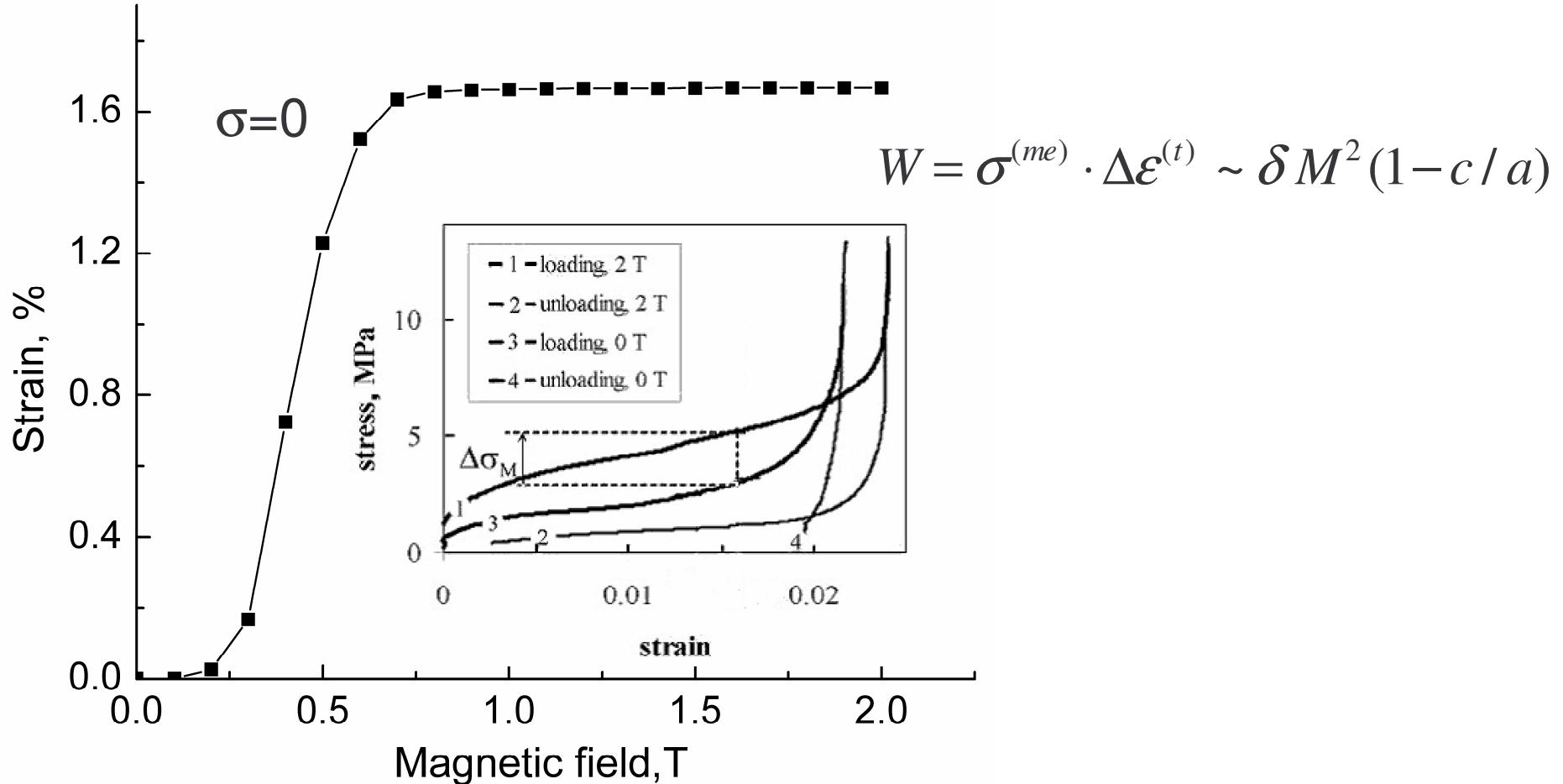
Murray 1999

Experiment with magneto-mechanical orthogonal loading



Muellner et al., 2002;2003

Strain vs H and stress-strain under H_{const}



Muellner et al., 2002;2003

MAGNETOELASTIC MODEL OF FERROMAGNETIC MARTENSITE

$$F = F_e + F_m + F_{me}. \quad (1)$$

In Landau-type approach, Eq. (1) is invariant under symmetry group PM cubic phase.

For cubic – tetragonal MT

$$F_e = \frac{3}{2}(C_{11} + 2C_{12})u_1^2 + \frac{1}{6}C'(u_2^2 + u_3^2) \quad (2)$$

$$u_1 = (\varepsilon_{xx} + \varepsilon_{yy} + \varepsilon_{zz})/3,$$

$$u_2 = \sqrt{3}(\varepsilon_{xx} - \varepsilon_{yy}), \quad u_3 = 2\varepsilon_{zz} - \varepsilon_{yy} - \varepsilon_{xx}$$

$$F_m = J y^2 / 2 + M^2 (\mathbf{m} \cdot \mathbf{D} \cdot \mathbf{m}) - \mathbf{m} \cdot \mathbf{H} \cdot \mathbf{M}, \quad (3)$$

$$y = M(T) / M(0) \quad \mathbf{m} = \mathbf{M}(T) / M(T)$$

$$F_{me} = -\delta_0 y^2(T) u_1 - \delta_1 \left[\sqrt{3} (m_x^2 - m_y^2) u_2 + (2m_z^2 - m_y^2 - m_x^2) u_3 \right], \quad (4)$$

Lvov, Chernenko et al
1998 - 2004

Magnetostress concept and mechanical equivalence

follows explicitly from magnetoelastic model of martensite

$$F = \frac{1}{2} C' \varepsilon_{ii}^2 + \frac{1}{2} A M_i^2 - \delta M_i^2 \varepsilon_{ii} - M_i H_i$$

$$A \sim \delta(1 - c/a)$$

$$\sigma_{ik} = (\partial F / \partial \varepsilon_{ik})_T \quad \begin{array}{c} \longrightarrow \\ \downarrow \end{array} \quad \sigma_{ii} = C' \varepsilon_{ii} + \delta M_i^2$$

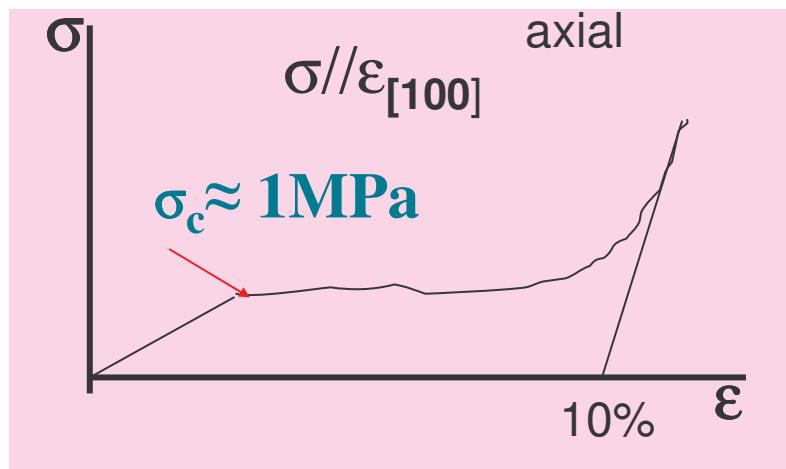
generalized Hooke's law for stressed ferromagnetic crystal

Conclusion:

magnetoelastic coupling is responsible for magnetostress

Equivalence principle (*JMMM, 1999; EPJ-B, 2002*)

magnetoelastic stress state = mechanical stress state

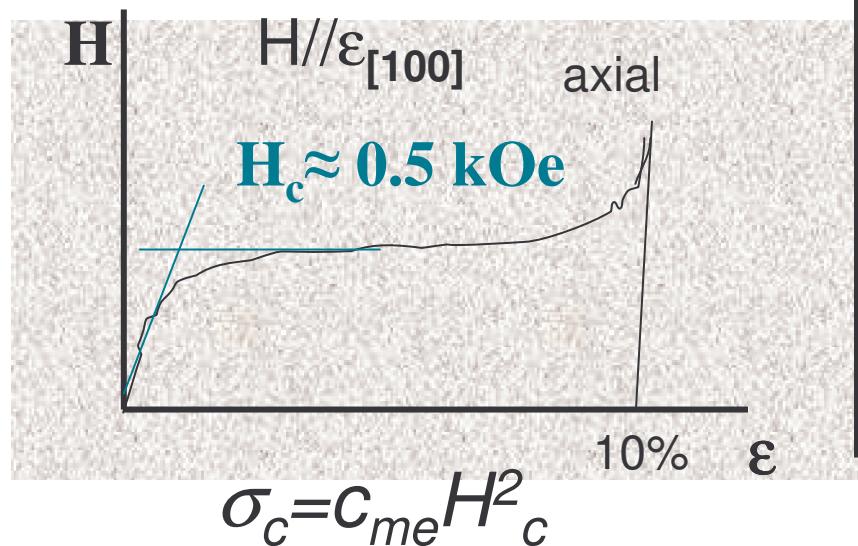


(Quasi-)Elastic regime:

$$\sigma^{\text{mech}}, \sigma^{\text{mag}} = \mathbf{S} \cdot \boldsymbol{\epsilon}$$

where $\epsilon \ll \epsilon^{(\text{MT})}$

ϵ reversible after force removal

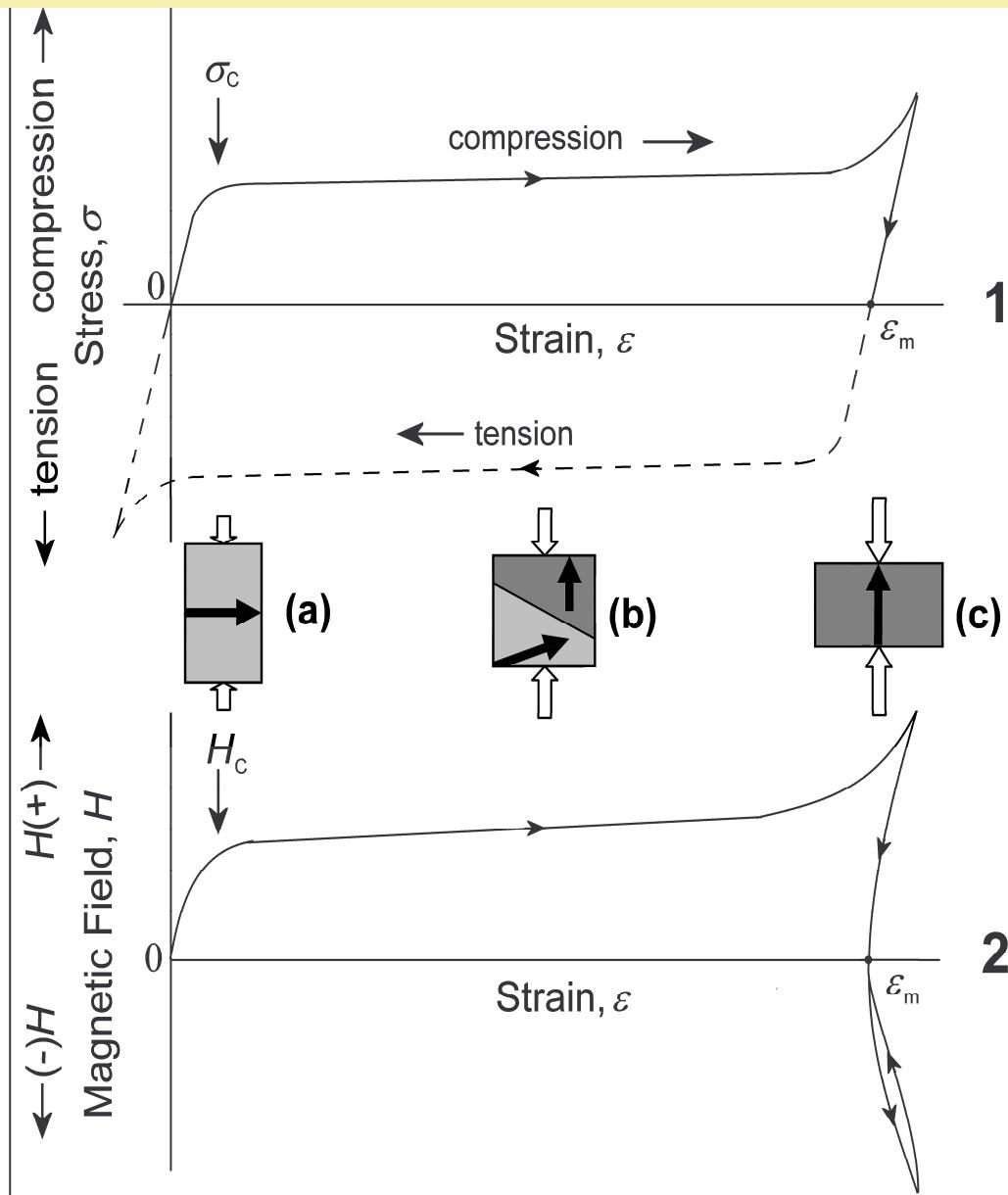


Anelastic (flow) regime:

$$\epsilon \sim \epsilon^{(\text{MT})} \sim f(\alpha)$$

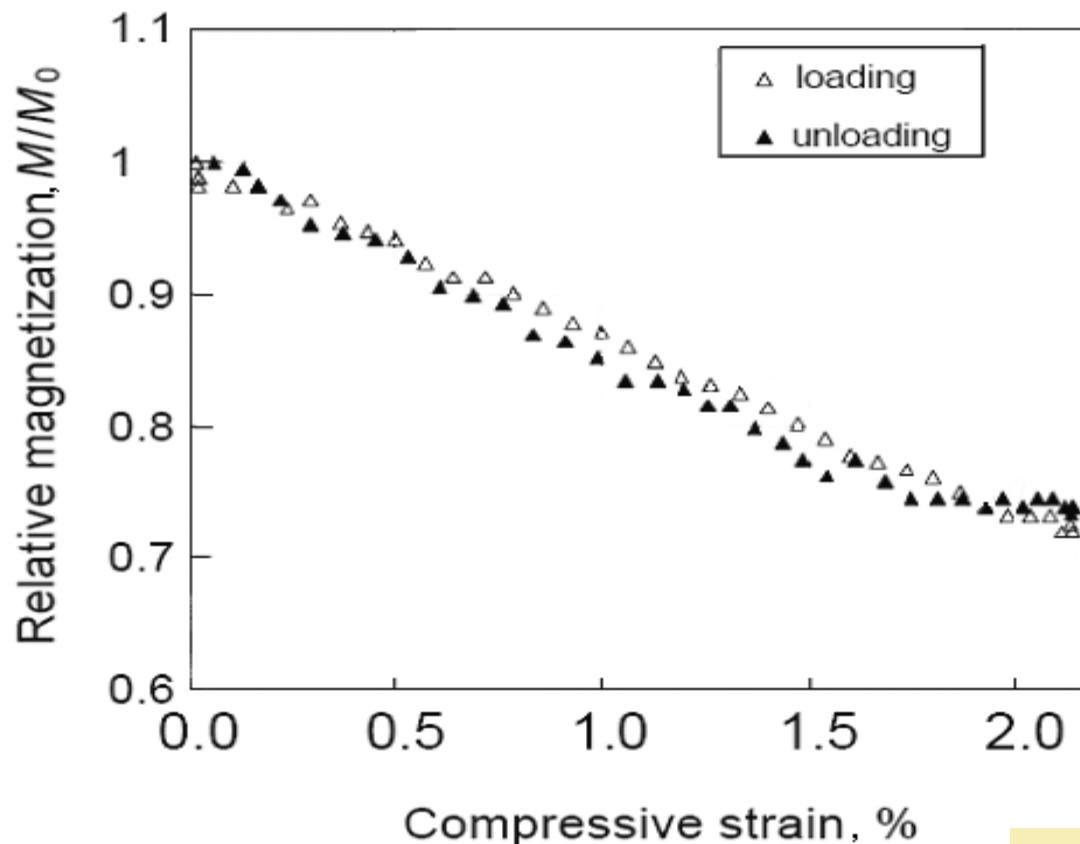
ϵ reversible when force direction is changed

Difference in mechanical and magnetic actuation



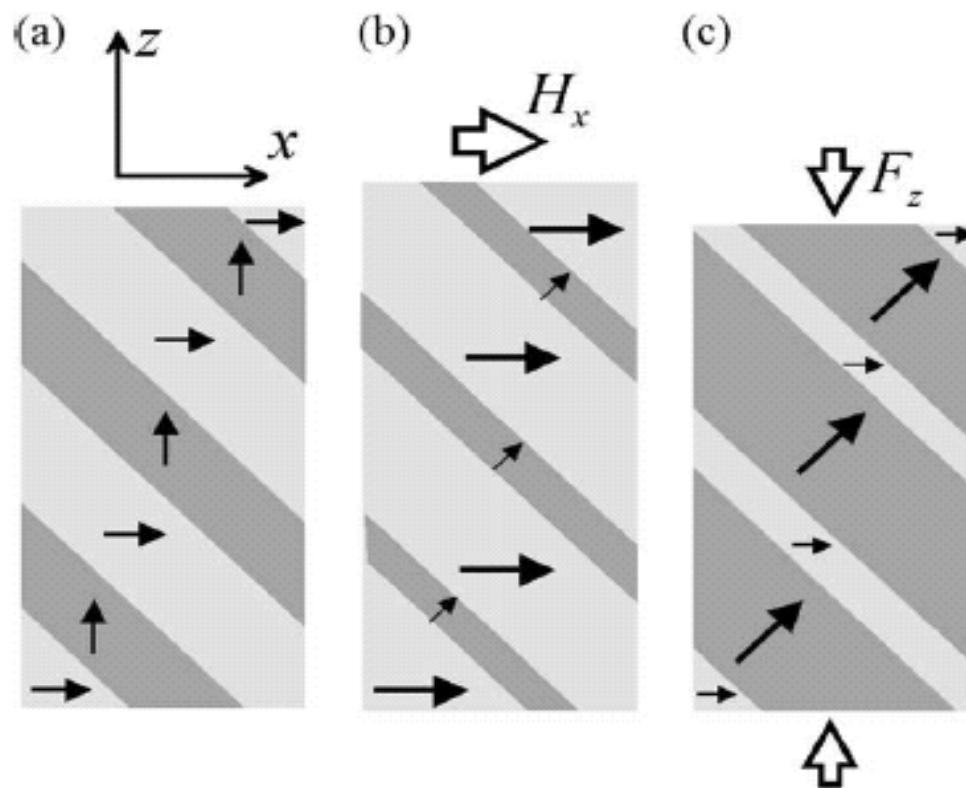
Reverse to giant MFIS effect

Magnetization change due to twinning strain under orthogonal loading

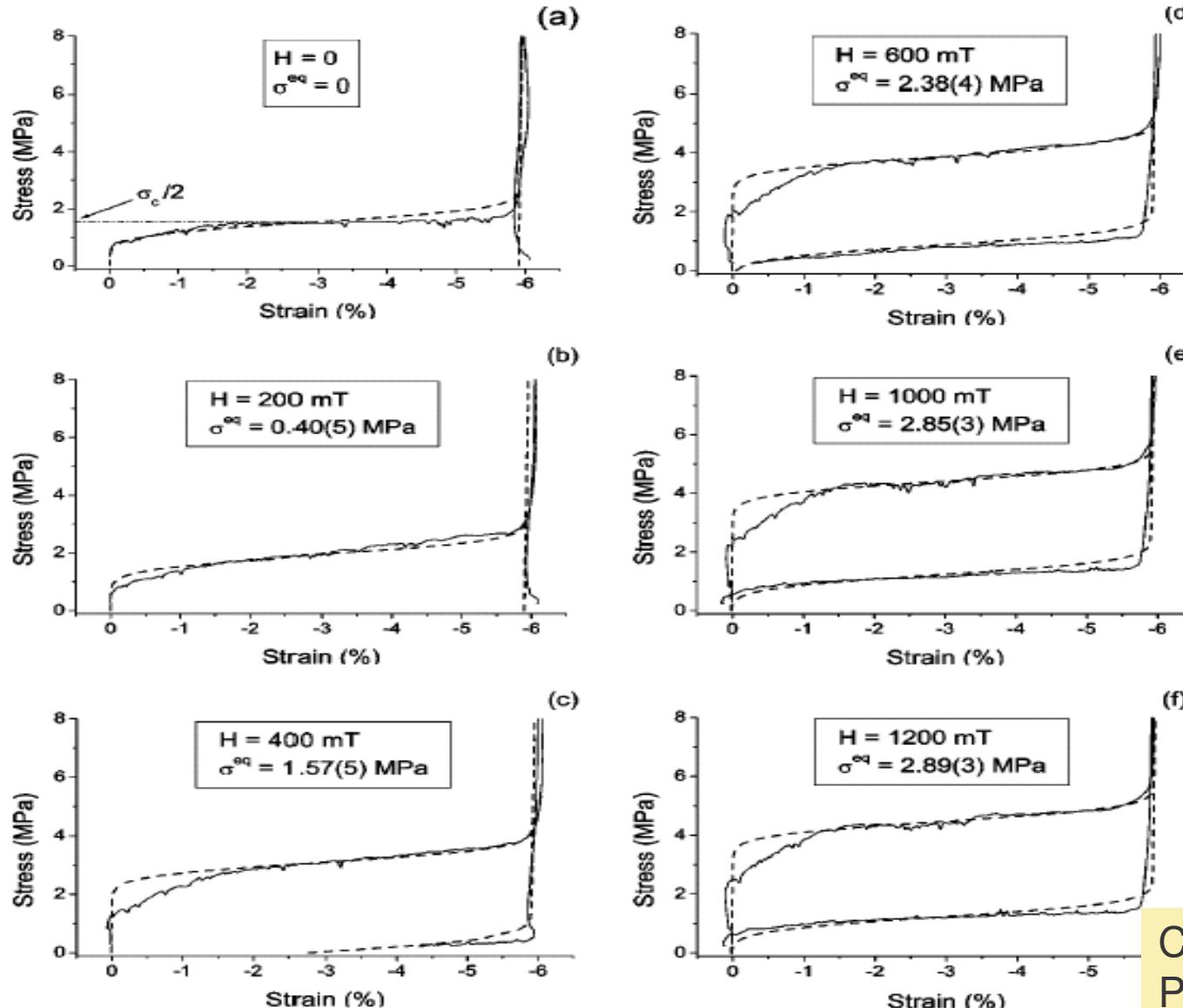


Muellner et al., 2003

Schematic of magnetization process through deformation in ferromagnetic martensite



Magnetomechanics of single variant $\text{Ni}_{52.0}\text{Mn}_{24.4}\text{Ga}_{23.6}$ marteniste



$\epsilon[010]$

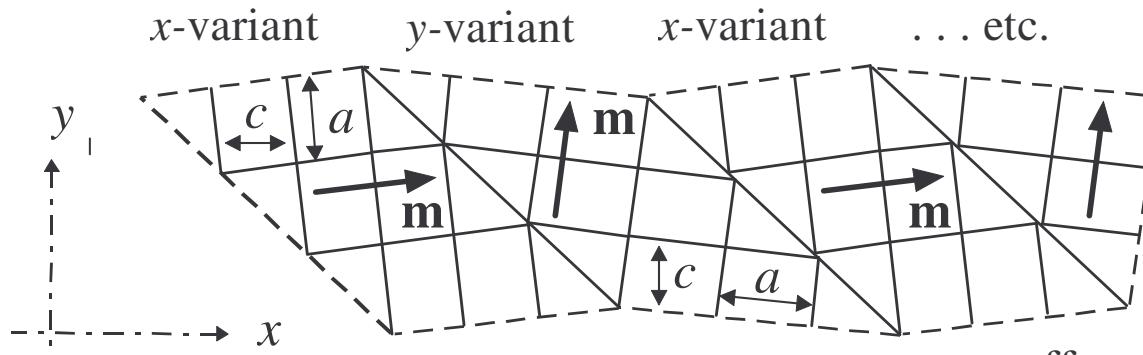
$H/[100]$

tetragonal
martensite
 $c/a < 1$

$T_{\text{exp}} = 293 \text{ K}$

Chernenko,Lvov et al
PRB, 2004

Magnetic field-induced superelasticity of Ni-Mn-Ga marteniste: modeling



- (i) reorientation of variants is caused by total stress which is sum of mechanical and field-induced stresses, tensile if positive, compressive if negative, it is acting in x direction in each twin variant.
- (ii) two-stage twin boundary rearrangement
- (iii) critical stress σ_n^{eff} is different for twin variants
- (iv) statistical distribution of critical stresses around certain σ_c

$$p_n = \frac{1}{Z} \exp \left\{ - \frac{(|\sigma_n^{\text{eff}}| - |\sigma_c|)^2}{2\sigma_0^2} \right\},$$

$$\varepsilon_{yy}^{(xy)} = -(S^{-1}/2)\sigma^{\text{eff}}(\sigma_{yy}, H) + (c/a - 1)[\alpha_y(\sigma_{yy}, H) - \alpha_y(0, 0)].$$

For each stress value, N of disappeared twins (boundaries jumps) is found from condition $|\sigma^{\text{eff}}| \geq |\sigma_n^{\text{eff}}|$. which gives rise with (ii) to:

$$N(\sigma^{\text{eff}}) = N_0 \sum p_n \theta(|\sigma^{\text{eff}}| - |\sigma_n^{\text{eff}}|),$$

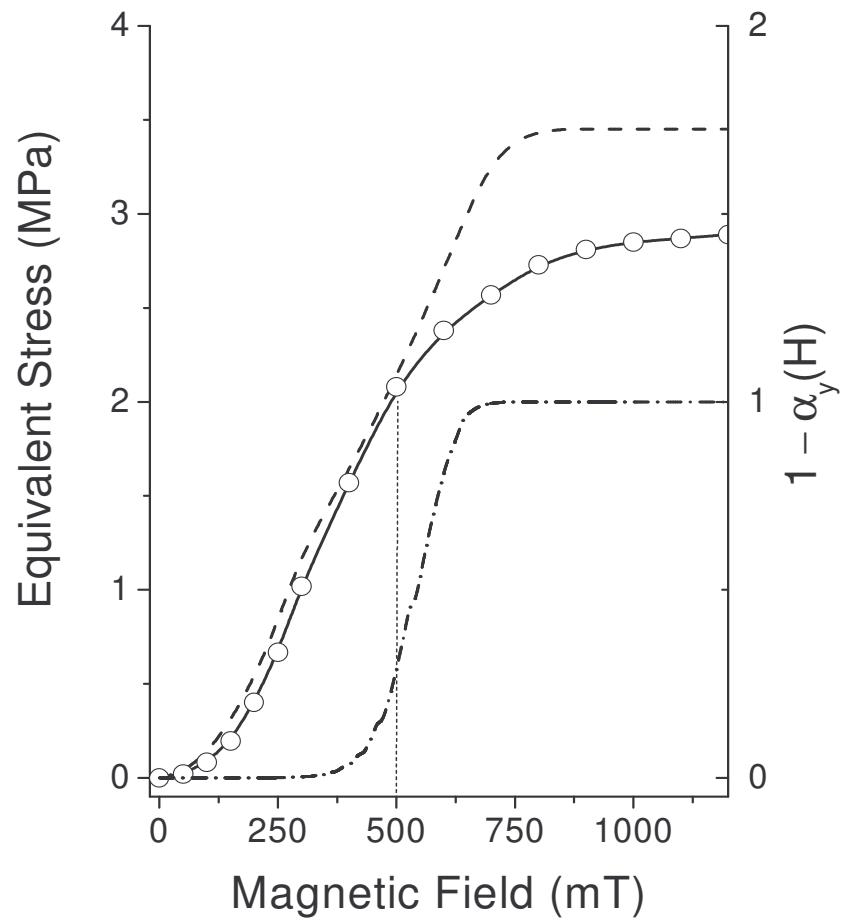
According to (i): $\sigma^{\text{eff}} = \sigma^{\text{eff}}(\sigma_{ik}^n, H_j)$, $N = N(\sigma_{ik}, H_j)$.

During compression along y ($H \parallel x$)

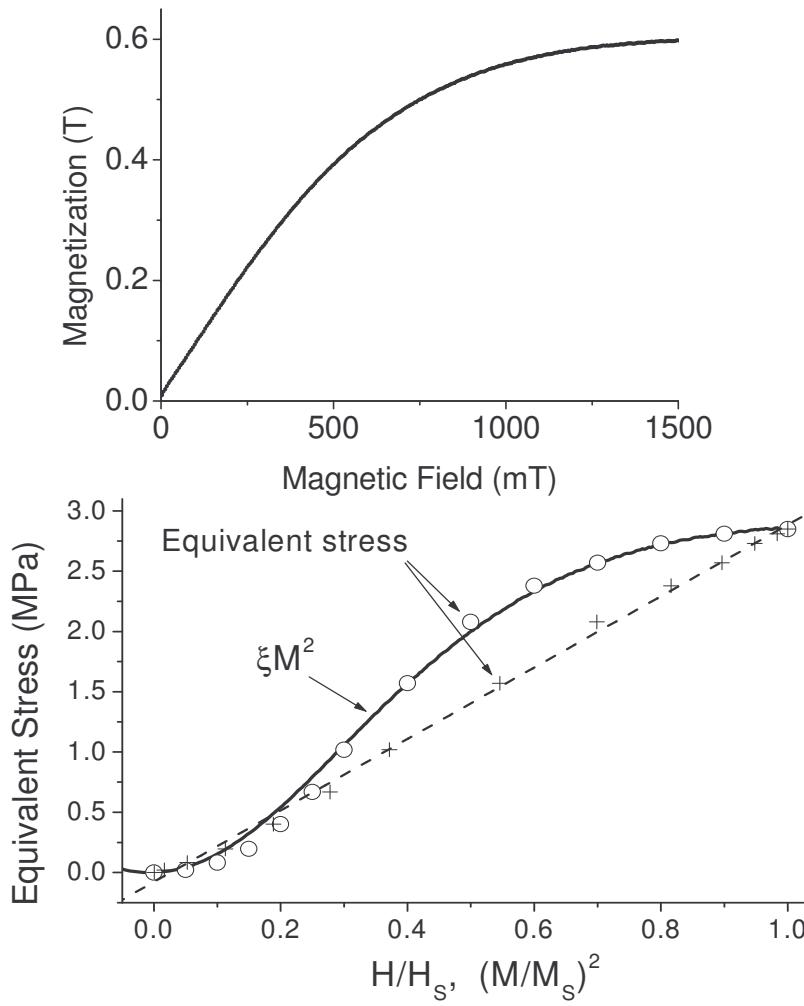
$$\alpha_y(\sigma_{yy}, H) = \begin{cases} \alpha_y(0, H)[1 - N(\sigma_{yy}, H)/N_0], & \sigma^{\text{eff}} < 0, \\ \alpha_y(0, H) + [1 - \alpha_y(0, H)]N(\sigma_{yy}, H)/N_0, & \sigma^{\text{eff}} \geq 0. \end{cases}$$

During unloading with $H \parallel x$

$$\alpha_y(\sigma_{yy}, H) = \begin{cases} \alpha_y(\sigma_{yy}^{\max}, H)[1 - N(\sigma_{yy}, H)], & \sigma^{\text{eff}} < 0, \\ \alpha_y(\sigma_{yy}^{\max}, H), & \sigma^{\text{eff}} > 0, \end{cases}$$



Experiment: open circles
 Fitting: solid line
 Thermoelastic model: dash
 $\sigma_2(H_s)/2\sqrt{3}=345 \text{ MPa}$
 x -variant fraction: dash-dot



Magnetoelastic origin of magneto-strain in martensite !

Interdependence between magnetic properties and lattice parameters of Ni-Mn-Ga martensite

Research line followed scheme:

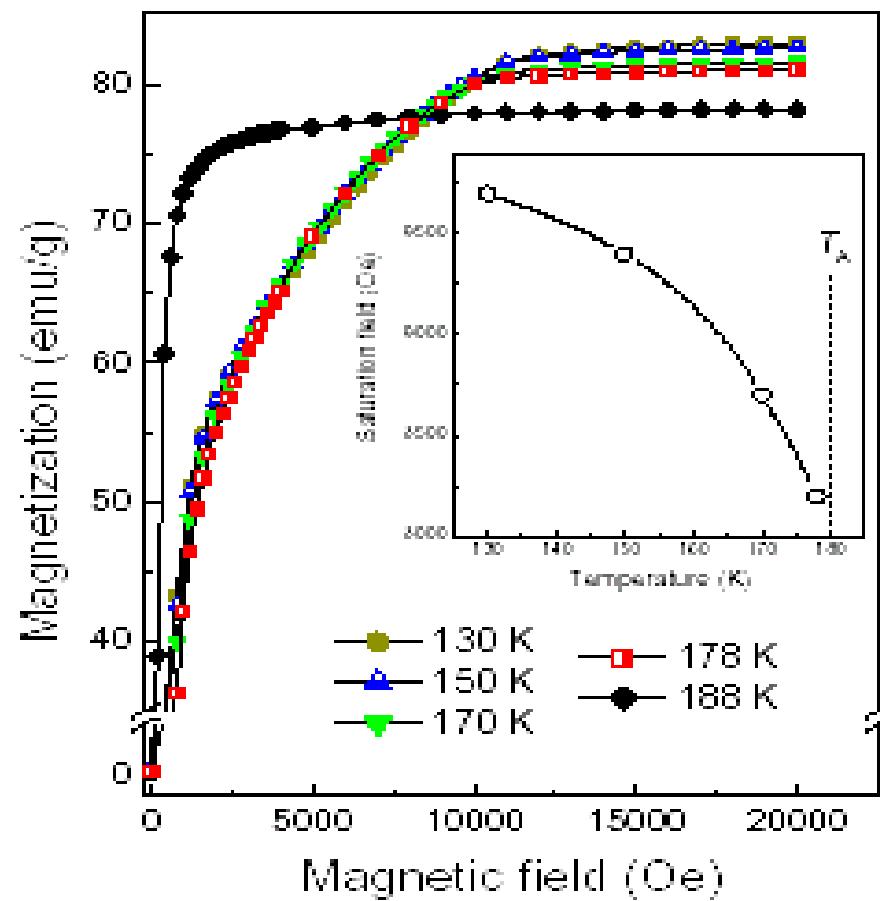
- (i) measurements of $H_s(T)$
- (ii) determination of $c/a(T)$ according to our phenomenological theory:

$$c/a = 1 - \frac{[(H_s/M) + |D_1 - D_2|]}{12|\delta|}.$$

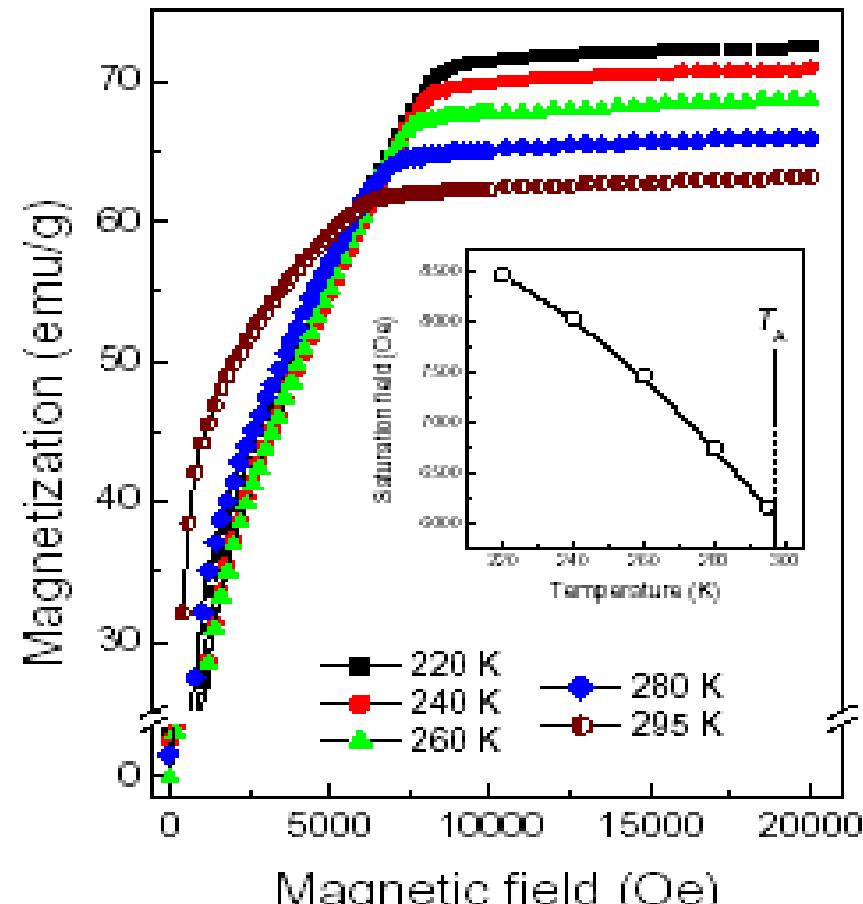
- (iii) comparison with $c/a(T)$ function obtained by X-rays measurements.

$\text{Ni}_{49.5}\text{Mn}_{25.4}\text{Ga}_{25.1}$ (A1) : $T_M = 170\text{ K}$, $T_A = 180\text{ K}$, $T_C = 381\text{ K}$

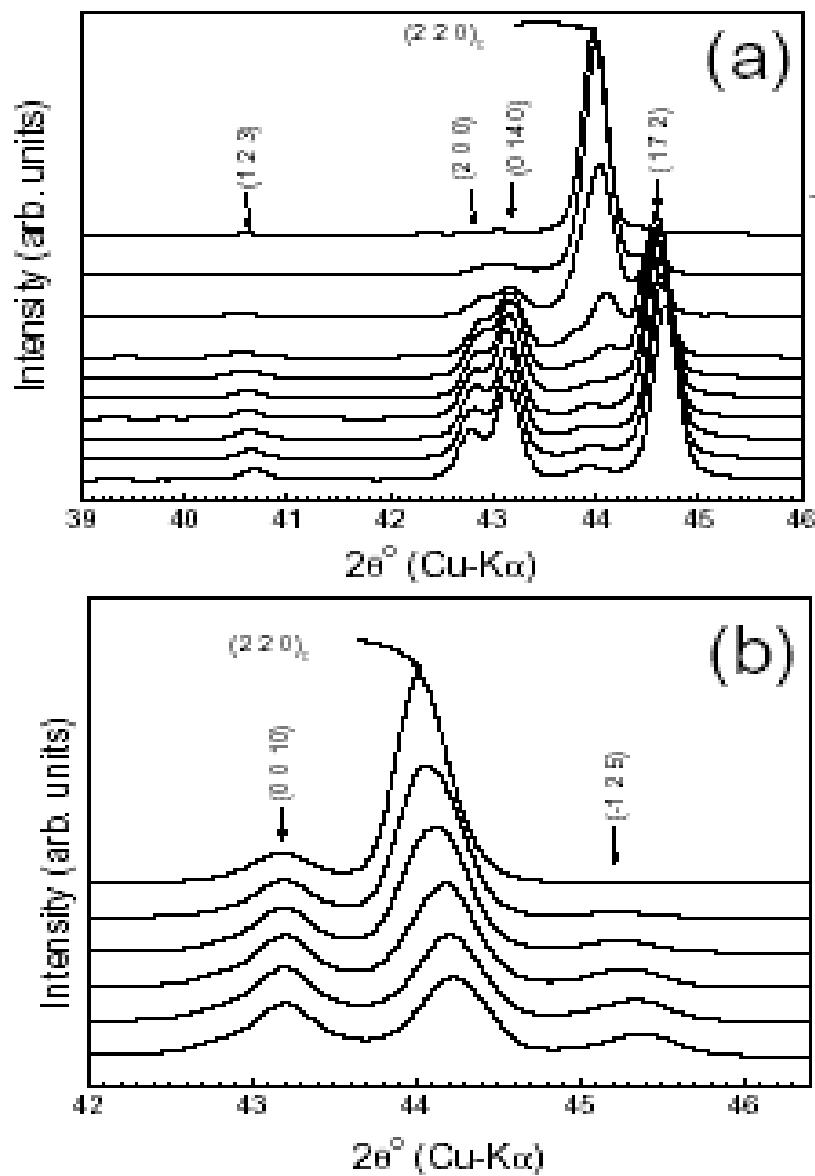
$\text{Ni}_{52.6}\text{Mn}_{23.6}\text{Ga}_{23.8}$ (A2) : $T_M = 284\text{ K}$, $T_A = 297\text{ K}$, $T_C = 363\text{ K}$



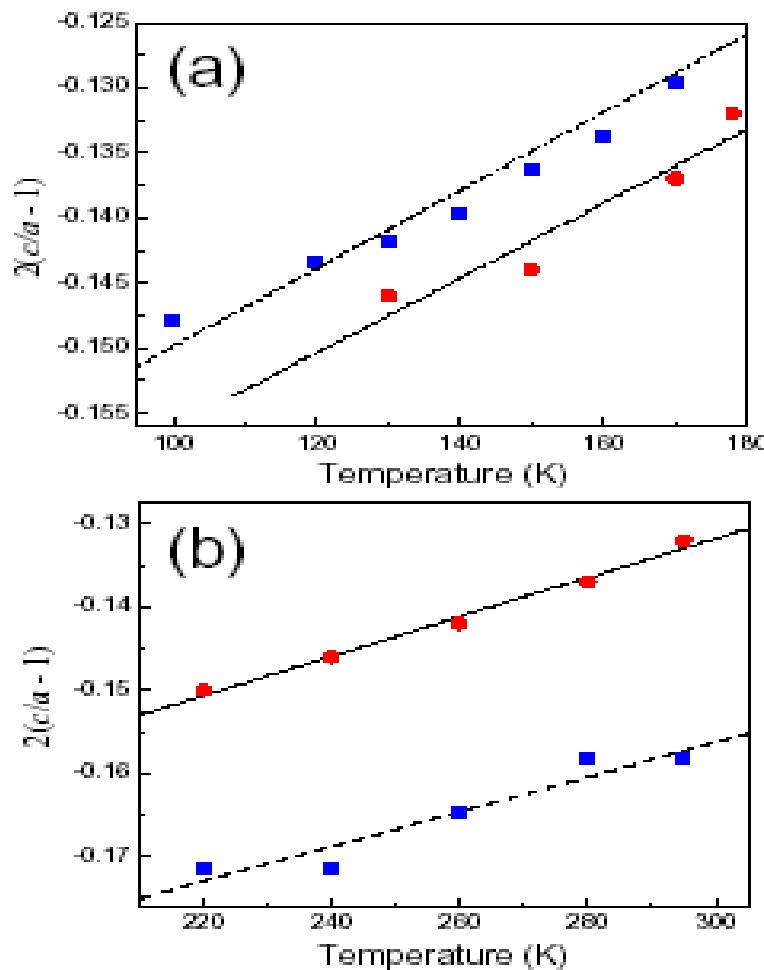
A1



A2



A1



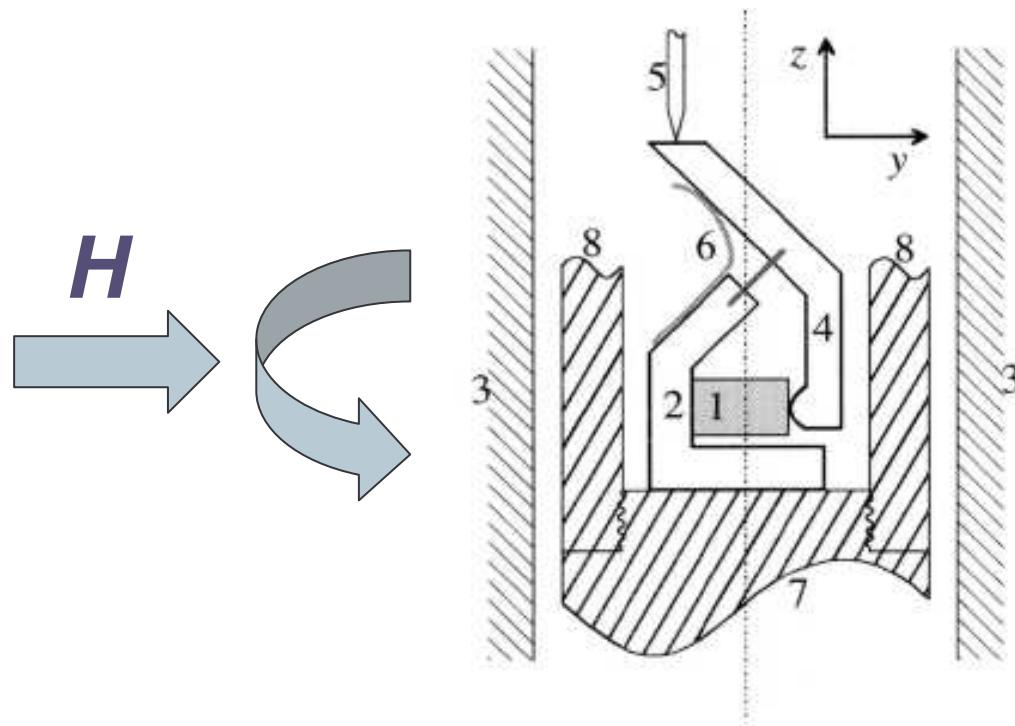
A2

From X-rays (squares)

From M(T) (circles)

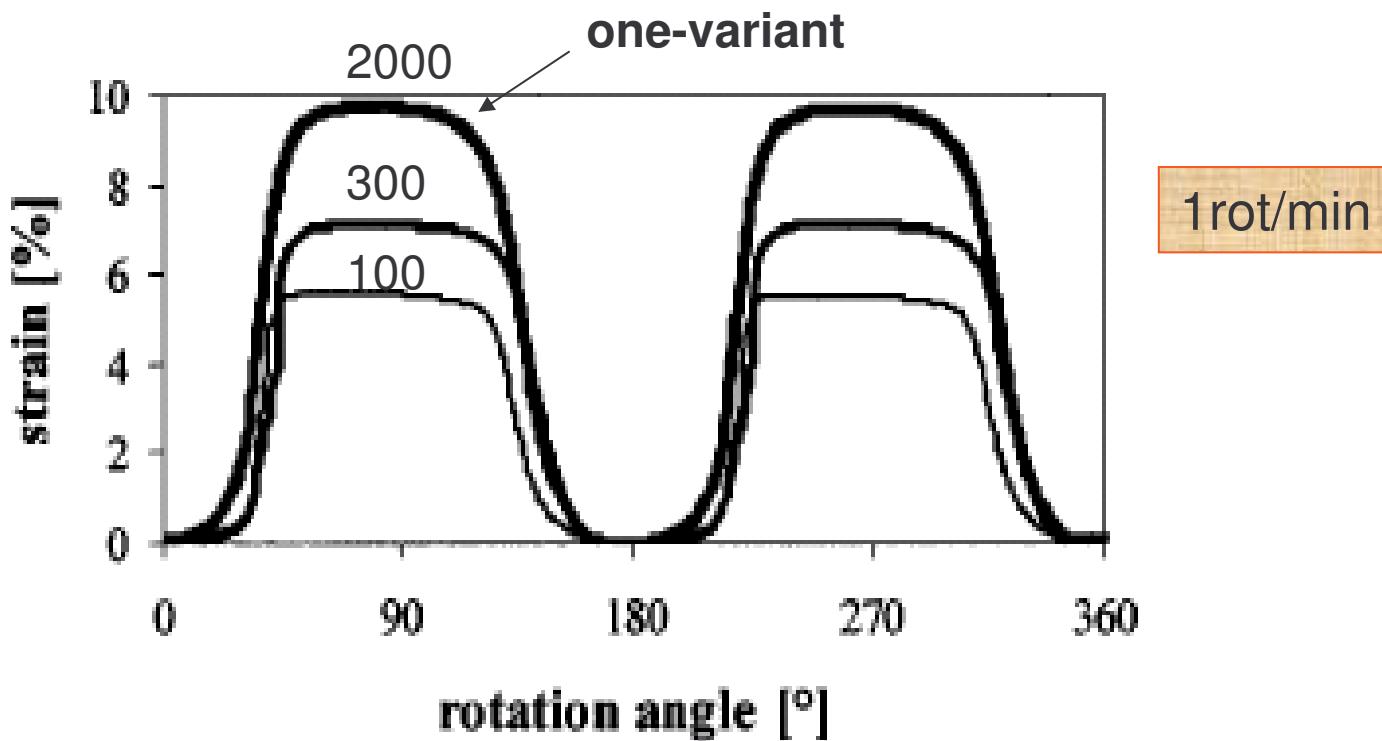
So, validity of Eq. c/a vs H_s is confirmed!

Magnetostress loading in rotation configuration



Muellner et al., 2002;2004

Periodic field-induced strain in Ni-Mn-Ga 14M martensite in rotating field of 10 kOe



Muellner et al., 2004

Ferromagnetic martensite

Magnetisation curves (**quasielastic regime of magnetostain**

- (i) Reversible displacements of 180 deg. domain walls
- (ii) Reversible rotation of magnetic vectors of variants
- (iii) $\beta = \text{const}$

$$M(H,T) = M(T)[\beta\Delta(H,T) + (1-\beta)\cos\psi(H,T)]$$

$H/[100] < H_s$

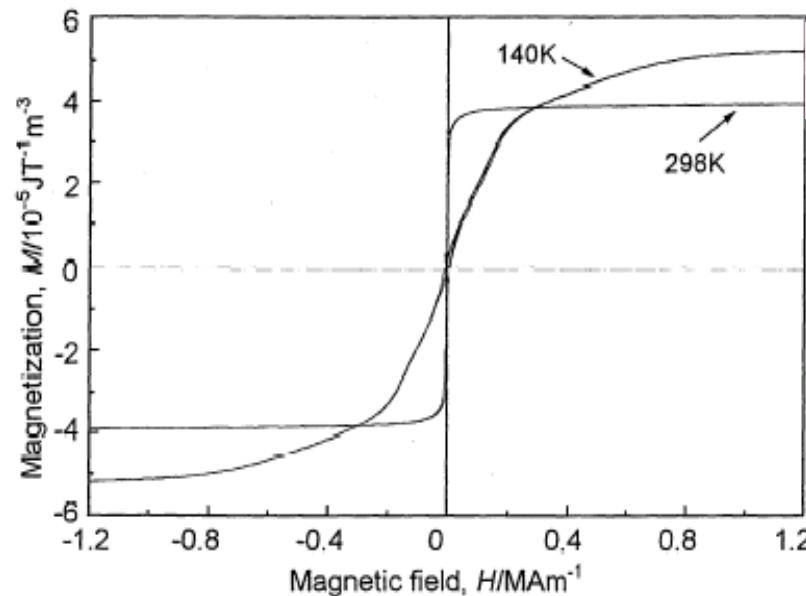
$$M(T) = M(0) \tanh[(T_c/T)M(T)/M(0)]$$

$$\Delta(H,T) = H / D_2 M(H,T)$$

for $H < H_c$

$$\Delta(H,T) = 1$$

for $H > H_c$

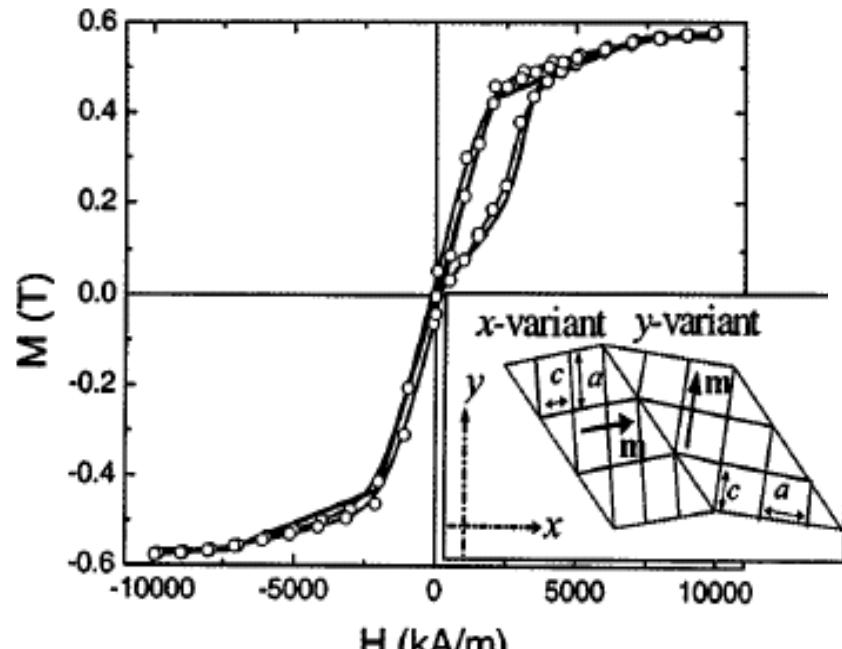


Lvov, Chernenko et al., 2003

Ferromagnetic martensite

Magnetisation curves (twin rearrangement)

(iii) $\alpha = f(H)$



For xy variant, $H/[010]$

$$M_{\text{rot}} = M_S [(H/H_S) \theta(H_S - H) + \theta(H - H_S)],$$

$$M_{\text{dis}} = M_S [(H/H_c) \theta(H_c - H) + \theta(H - H_c)],$$

$$M_{\text{rel}} = \alpha_y(H) M_{\text{dis}} + [1 - \alpha_y(H)] M_{\text{rot}},$$

$$\alpha_y(\sigma) = \begin{cases} \alpha_y(0)[1 - N(\sigma)/N_0], & \sigma < 0 \\ \alpha_y(0) + [1 - \alpha_y(0)]N(\sigma)/N_0, & \sigma \geq 0 \end{cases}$$

$$\alpha_y(\sigma) = \begin{cases} \alpha_y(\sigma_{\max})[1 - N(\sigma)/N_0], & \sigma < 0 \\ \alpha_y(\sigma_{\max}), & \sigma \geq 0 \end{cases}$$

$T_n = \sigma(H_n)$

$$\sigma(H) = c_{me} H^2,$$

$$\langle M \rangle = \frac{1}{3}(M^{xy} + M^{zy} + M^{zx})$$

where $c_{me} = -12\delta M_S^2 / H_S^2$

$M^{xy} = M^{zy} = M_{\text{rel}}$, $M^{zx} = M_{\text{rot}}$

H increasing

H decrease

$M^{xy} = M^{zy} = M_{\text{dis}}$

Conclusions

Magnetoelastic model satisfactorily describes magnetomechanics in quasi-elastic regime

While

In case of twin rearrangements, statistical approach should be additionally involved to account for peculiar magnetomechanics and magnetization phenomena in ferromagnetic thermoelastic martensites

Acknowledgments

- L'vov V.A., Taras Shevchenko University, Kiev, Ukraine
- Cesari E., UIB, Palma de Mallorca, Spain
- Mullner P., Boise University, Boise, USA
- Besseghini S., CNR-IENI, Lecco, Italy
- Kostorz G.K., ETH, Zuerich, Switzerland
- Pasquale M., IEN, Torino, Italy
- Takagi T., Tohoku University, Sendai, Japan